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Ground States of Nonlinear Schrödinger Equations with potentials vanishing at infinity

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Analisi matematica. — Ground States of Nonlinear Schrödinger Equations with potentials vanishing at infinity. Nota di ANTONIO AMBROSETTI, VERONICA FELLI E AN-DREA MALCHIODI, presentata (*) dal Socio A. Ambrosetti.

ABSTRACT. — In this preliminary *Note* we outline the results of the forthcoming paper [2] dealing with a class on nonlinear Schrödinger equations with potentials vanishing at infinity. Working in weighted Sobolev spaces, the existence of a ground state is proved. Furthermore, the behaviour of such a solution, as the Planck constant tends to zero (semiclassical limit), is studied proving that it concentrates at a point.

KEY WORDS: Nonlinear Schrödinger equations; Weighted Sobolev spaces; Critical point theory.

RIASSUNTO. — Stati fondamentali per equazioni di Schrödinger nonlineari con potenziali che si annullano all'infinito. In questa Nota preliminare presentiamo i risultati del lavoro [2] dove studiamo una classe di equazioni di Schrödinger nonlineari con potenziali che tendono a zero all'infinito. Lavorando in spazi di Sobolev con peso, dimostriamo l'esistenza di una soluzione fondamentale. Di tale soluzione è anche studiato il comportamento quando la costante di Planck tende a zero (limite semiclassico) dimostrando che essa si concentra in un punto.

1. INTRODUCTION

We consider, for $N \ge 3$, the stationary Nonlinear Schrödinger Equations

(1)
$$\begin{cases} -\varepsilon^2 \Delta v + V(x)v = K(x)v^p, & x \in \mathbb{R}^N, \\ v \in W^{1,2}(\mathbb{R}^N), & v(x) > 0, & \lim_{|x| \to \infty} v(x) = 0, \end{cases}$$

where 1 . We address here two problems: (*i* $) the existence, for <math>\varepsilon > 0$ fixed, of a solution v_{ε} of (1) with minimal energy (ground state); (*ii*) the behavior (concentration) of v_{ε} as $\varepsilon \to 0$. The main novelty with respect to most of the (broad) literature dealing with (1) is that we assume that the potentials *V* and *K* decay to zero as $|x| \to \infty$. Precisely, we suppose

(V) $V: \mathbb{R}^N \to \mathbb{R}$ is smooth and $\exists \alpha, a_1, a_2 > 0$ such that

$$\frac{a_1}{1+|x|^{\alpha}} \le V(x) \le a_2,$$

and, respectively

(K) $K: \mathbb{R}^N \to \mathbb{R}$ is smooth and $\exists \beta, a_3 > 0$:

$$0 < K(x) \le \frac{a_3}{1 + |x|^{\beta}}$$

If $V \sim |x|^{-\alpha}$ as $|x| \to \infty$, with $\alpha > 0$, the spectrum of the linear operator $-\varDelta + V$ is

(*) Nella seduta del 12 marzo 2004.

 $[0, +\infty)$, and this prevents the use of perturbation methods as in [3, 4]. On the other hand, we cannot even apply critical point theory working in $W^{1,2}(\mathbb{R}^N)$. For these reasons, we consider the weighted space $L_{\mathcal{K}}^{q}$ of measurable $u : \mathbb{R}^N \to \mathbb{R}$ such that

$$|u|_{q,K} = \left[\int_{\mathbb{R}^N} K(x) |u(x)|^q dx\right]^{\frac{1}{q}} < \infty,$$

as well as the weighted Sobolev spaces $\mathcal{H}_{\varepsilon}$ defined by setting

$$\mathcal{H}_{\varepsilon} = \left\{ u \in \mathcal{O}^{1,\,2}(\mathbb{R}^{N}) : \int_{\mathbb{R}^{N}} \left[\varepsilon^{2} \left| \nabla u(x) \right|^{2} + V(x) u^{2}(x) \right] \, dx < \infty \right\}.$$

 $\mathcal{H}_{\varepsilon}$ is a Hilbert space with scalar product

$$(u \mid v)_{\varepsilon} = \int_{\mathbb{R}^{N}} \left[\varepsilon^{2} \nabla u(x) \cdot \nabla v(x) + V(x) u(x) v(x) \right] dx$$

and norm $||u||_{\varepsilon}^2 = (u | u)_{\varepsilon}$. Let

$$\sigma = \sigma_{N, \alpha, \beta} = \begin{cases} \frac{N+2}{N-2} - \frac{4\beta}{\alpha(N-2)}, & \text{if } 0 < \beta < \alpha\\ 1 & \text{otherwise.} \end{cases}$$

The above weighted spaces have been introduced in [6], where the following result is proved.

THEOREM 1. Let $N \ge 3$ and suppose that (V), (K) hold with $\alpha \in (0, 2]$ and $\beta > 0$, respectively. Then for all $\varepsilon > 0$, $\mathcal{H}_{\varepsilon} \hookrightarrow L_{K}^{p+1}$ provided

$$\sigma \leq p \leq \frac{N+2}{N-2}$$

Furthermore, the embedding of $\mathcal{H}_{\varepsilon}$ into L_{K}^{q} is compact provided

(2)
$$\sigma$$

2. AN EXISTENCE RESULT

The preceding Theorem implies that the functional I_{ε}

$$I_{\varepsilon}(u) = \frac{1}{2} \int_{\mathbb{R}^{N}} \left[\varepsilon^{2} \left| \nabla u(x) \right|^{2} + V(x) u^{2}(x) \right] dx - \frac{1}{p+1} \int_{\mathbb{R}^{N}} K(x) \left| u(x) \right|^{p+1} dx, \quad u \in \mathcal{H}_{\varepsilon},$$

is well defined and $I_{\varepsilon} \in C^{1}(\mathcal{H}_{\varepsilon}, \mathbb{R})$. It is easy to check that I_{ε} has the Mountain Pass geometry. Furthermore, the Palais-Smale condition is satisfied if (2) holds, since in such a case $\mathcal{H}_{\varepsilon}$ is compactly embedded into L_{K}^{σ} . This immediately implies:

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LEMMA 2. Let (V), (K) hold with $0 < \alpha \le 2$, $\beta > 0$, respectively, and suppose that p satisfies (2). Then

$$b_{\varepsilon} = \inf_{u \in \mathcal{H}_{\varepsilon} \setminus \{0\}} \max_{t \ge 0} I_{\varepsilon}(tu)$$

is a critical level of I_{ε} and carries a critical point $v_{\varepsilon} \in \mathcal{H}_{\varepsilon}$ of I_{ε} .

Here v_{ε} is a critical point in the sense that $(I_{\varepsilon}'(v_{\varepsilon}) | u) = 0$ for all $u \in \mathcal{H}_{\varepsilon}$. By local elliptic regularity, it follows that v_{ε} is for all $\varepsilon > 0$ a positive (classical) solution of the equation

(3)
$$-\varepsilon^2 \Delta v + V(x)v = K(x)v^p, \quad x \in \mathbb{R}^N.$$

REMARKS. *a*) Lemma 2 also follows from [7, Thm. 3.1] combined with Theorem 1. Let us point out that the case in which $p = \sigma$ or $p = \frac{N+2}{N-2}$ is also studied in [7, Thm. 3.2], under some further restriction on *V* and *K*. Lemma 2 is also somewhat related to the results of [5].

b) If V is abounded away from zero and infinity», namely $0 < \inf_{\mathbb{R}^N} V \leq \sup_{\mathbb{R}^N} V < +\infty$, we can directly work in $W^{1,2}(\mathbb{R}^N)$. In such a case we recover the compactness assuming that (K) holds. If $V \sim |x|^{-\alpha}$ as $|x| \to \infty$, with $\alpha \in (0, 2]$, while K is bounded away from zero and infinity, one could show that $b_{\varepsilon} = 0$ and hence there are no Mountain Pass solution. The same remark holds if $p < \sigma$. On the other hand, solutions with higher energy might exist. For example, if both V and K are bounded away from zero and infinity, the existence of a solution to (1) is proved *e.g.* in [3] provided ε is *sufficiently small* and $V^{\theta}K^{-2/(p-1)}$, with $\theta = \frac{p+1}{p-1} - \frac{N}{2}$, has a «stable» stationary point, like *e.g.* a maximum or a minimum.

Next, we will show that the Mountain-Pass solution $v_{\varepsilon} \in \mathcal{H}_{\varepsilon}$ is indeed a ground state, provided $0 < \alpha < 2$. In order to prove this fact, some sharp decay estimates are carried out, leading to the following integral estimate (in the lemma below, we have highlighted the dependence on ε because it will be used in the next section, dealing with the concentration phenomenon).

LEMMA 3. Let (V), (K) hold with $0 < \alpha < 2$, $\beta > 0$, respectively, and suppose that p satisfies (2). Moreover, let v_{ε} be solutions of (3) and suppose there exists $\Gamma > 0$ such that

(4)
$$\|v_{\varepsilon}\|_{\mathcal{H}_{\varepsilon}} \leq \Gamma \varepsilon^{n}$$

Then there exist $R_{\Gamma} > 0$, and constants $C_{1,2} > 0$, depending only on Γ , such that, for all $R \ge R_{\Gamma}$ there holds

(5)
$$\int_{|x|>R} [\varepsilon^2 |\nabla v_{\varepsilon}|^2 + V(x) v_{\varepsilon}^2] dx \leq C_1 \varepsilon^N \exp\left\{-C_2 \varepsilon^{-1} R^{\frac{2-\alpha}{2}}\right\}.$$

The preceding lemma is used to show that $v_{\varepsilon} \in L^{2}(\mathbb{R}^{N})$. Roughly, if $y \in \mathbb{R}^{N}$ with $|y| \gg 1$, one gets:

$$\int_{|x-y|<1} v_{\varepsilon}^2 dx = \int_{|x-y|<1} V(x) v_{\varepsilon}^2 \cdot \frac{1}{V(x)} dx \leq c_1 |y|^{\alpha} \int_{|x-y|<1} V(x) v_{\varepsilon}^2 dx.$$

Then (5) implies

x

$$\int_{-y|<1} v_{\varepsilon}^2 dx \le c_2 |y|^a \exp\{-c_3 |y|^{1-a/2}\},\$$

and from this it follows that $v_{\varepsilon} \in L^{2}(\mathbb{R}^{N})$. Then one gets that $v_{\varepsilon} \in W^{1,2}(\mathbb{R}^{N})$ as well as that $\lim_{|x|\to\infty} v(x) = 0$, proving our main existence result:

THEOREM 4. Let (V), (K) hold with $0 < \alpha < 2$, $\beta > 0$, respectively, and suppose that p satisfies (2). Then the Mountain-Pass solution v_{ε} found in Lemma 2 is such that $v_{\varepsilon} \in W^{1,2}(\mathbb{R}^N)$, $v_{\varepsilon} \in C^2(\mathbb{R}^N)$, $v_{\varepsilon}(x) > 0$ and $\lim_{|x| \to \infty} v_{\varepsilon}(x) = 0$ and thus is a ground state of (1).

3. Concentration as $\varepsilon \rightarrow 0$

Concerning the behavior of the Mountain-Pass solution v_{ε} , our main result is Theorem 5 below. Let

$$\mathcal{C}(x) := [V(x)]^{\theta} [K(x)]^{-2/(p-1)}, \qquad \theta = \frac{p+1}{p-1} - \frac{N}{2}.$$

The auxiliary potential \mathcal{A} has been previously introduced dealing with potentials V and K bounded away from zero and infinity, see *e.g.* [3], where it is proved that concentration occurs at the «stable» stationary points of \mathcal{A} . Moreover, let us point out that \mathcal{A} has a global minimum since $\lim_{|x| \to \infty} \mathcal{A}(x) = +\infty$ provided (2) holds.

THEOREM 5. Let the same assumptions as in Theorem 4 hold. Then, the preceding ground state v_{ε} concentrates at a global minimum x^* of \mathfrak{A} . More precisely, v_{ε} has a unique maximum x_{ε} such that $x_{\varepsilon} \rightarrow x^*$ as $\varepsilon \rightarrow 0$, and

$$v_{\varepsilon}(x) \sim U^*\left(\frac{x-x_{\varepsilon}}{\varepsilon}\right), \quad as \ \varepsilon \to 0,$$

where U* is the unique positive radial solution of

$$-\Delta U^* + V(x^*) U^* = K(x^*)(U^*)^p.$$

The proof of Theorem 5 is based upon the following Lemmas. The first one provides an uniform bound on the Mountain-Pass critical level of I_{ε} .

LEMMA 6. There exists $\Gamma > 0$ such that $b_{\varepsilon} \leq \Gamma \varepsilon^{N}$, for all $\varepsilon > 0$ small.

To see this, we introduce the functional J_{ε} : $W^{1,2}(\mathbb{R}^N) \mapsto \mathbb{R}$, by setting

$$J_{\varepsilon}(u) = \frac{1}{2} \int_{\mathbb{R}^{N}} \left[\varepsilon^{2} |\nabla u|^{2} + a_{2} u^{2} \right] dx - \frac{1}{p+1} \int_{\mathbb{R}^{N}} K(x) |u|^{p+1} dx$$

where, according to assumption (V), sup $V \leq a_2$. Since $W^{1,2}(\mathbb{R}^N) \subset \mathcal{H}_{\varepsilon}$, we infer

(6)
$$b_{\varepsilon} \leq \tilde{b}_{\varepsilon} := \inf_{u \in W^{1,2}(\mathbb{R}^N) \setminus \{0\}} \max_{t \geq 0} J_{\varepsilon}(tu).$$

Furthermore, letting

$$\mathcal{J}_{\varepsilon}(u) = \frac{1}{2} \int_{\mathbb{R}^N} \left[\left| \nabla u \right|^2 + a_2 u^2 \right] dx - \frac{1}{p+1} \int_{\mathbb{R}^N} K(\varepsilon x) \left| u \right|^{p+1} dx.$$

One finds that $u_{\varepsilon}(x)$ is a critical point of $\mathcal{J}_{\varepsilon}$ iff $\tilde{u}_{\varepsilon}(x) := u_{\varepsilon}(x/\varepsilon)$ is a critical point of J_{ε} . Using the perturbation method introduced in [1], we can look for critical points of $\mathcal{J}_{\varepsilon}$ near those of the *unperturbed* functional $\mathcal{J}_0 \equiv \mathcal{J}_{\varepsilon=0}$. Up to translation, we can assume that $K(0) = \max K$. Then $\mathcal{J}_{\varepsilon}$ has, for $\varepsilon > 0$ small, a critical point u_{ε} such that $u_{\varepsilon} \to U$ as $\varepsilon \to 0$, where U is the unique positive radial solution of

$$-\varDelta U + a_2 U = K(0) U^p, \qquad U \in W^{1,2}(\mathbb{R}^N).$$

Furthermore, *U* is a Mountain Pass critical point of \mathcal{J}_0 and the same holds true for u_{ε} . The preceding information, together with $\mathcal{J}_{\varepsilon}(u_{\varepsilon}) \to \mathcal{J}_0(U)$ as $\varepsilon \to 0$ and the equality $J_{\varepsilon}(\tilde{u}_{\varepsilon}) = \varepsilon^N \mathcal{J}_{\varepsilon}(u_{\varepsilon})$, readily imply that $\tilde{b}_{\varepsilon} \leq \Gamma \varepsilon^N$, and the lemma follows using (6).

Lemma 6 implies that (4) holds. Then Lemma 3 applies and this allows us to prove a pointwise uniform exponential decay for the solutions v_{ε}

LEMMA 7. There exist C_1 , $C_2 > 0$ and d > 0, depending only on Γ , p, N, α and β , such that

(7)
$$|v_{\varepsilon}(x)| \leq C_1 |x|^d \varepsilon^{-d} \exp\left\{-C_2 \varepsilon^{-1} |x|^{\frac{2-\alpha}{2}}\right\}; \quad for |x| \gg 1.$$

Let x_{ε} denote a point of maximum for v_{ε} . Thus, $\Delta v_{\varepsilon}(x_{\varepsilon}) \leq 0$ whence $V(x_{\varepsilon})K^{-1}(x_{\varepsilon}) \leq v_{\varepsilon}^{p-1}(x_{\varepsilon})$. Then, using (7), we deduce that

$$V(x_{\varepsilon})K^{-1}(x_{\varepsilon}) \leq C_1 |x_{\varepsilon}|^{d(p-1)} \varepsilon^{-d(p-1)} \exp\left\{-C_2 \varepsilon^{-1} |x_{\varepsilon}|^{\frac{2-\alpha}{2}}\right\}.$$

Since from (V) - (K) we know that $V(x)K^{-1}(x) \sim |x|^{\beta - \alpha}$ as $|x| \to \infty$, there exists a constant C' > 0 such that

(8) $|x_{\varepsilon}| \leq C', \quad \forall \varepsilon \sim 0.$

Finally one also shows that there exists a constant C'' > 0

(9) $\|v_{\varepsilon}\|_{L^{\infty}} \ge C''.$

Equations (8), (9) and the preceding Lemmas allow us to carry over standard arguments which lead to prove Theorem 5.

The complete proofs of the results sketched in the present *Note*, are contained in the forthcoming paper [2].

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