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A $\mathcal{U}_q(\mathfrak{sl}(2))$ -representation with no quantum symmetric algebra

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Algebra. — *A $U_q(\mathfrak{sl}(2))$ -representation with no quantum symmetric algebra.* Nota (*)
di OLIVIA ROSSI-DORIA, presentata dal Corrisp. C. De Concini.

ABSTRACT. — We show by explicit calculations in the particular case of the 4-dimensional irreducible representation of $U_q(\mathfrak{sl}(2))$ that it is not always possible to generalize to the quantum case the notion of symmetric algebra of a Lie algebra representation.

KEY WORDS: Quantized enveloping algebra; Representation; Symmetric algebra.

RIASSUNTO. — *Una rappresentazione di $U_q(\mathfrak{sl}(2))$ che non ha algebra simmetrica quantica.* Si dimostra, mediante calcoli espliciti per la rappresentazione irriducibile di dimensione 4 di $U_q(\mathfrak{sl}(2))$, che non è sempre possibile generalizzare al caso quantico la nozione di algebra simmetrica di una rappresentazione di un'algebra di Lie.

1. PRELIMINARIES

The quantized enveloping algebra $U_q(\mathfrak{g})$ was introduced by Drinfeld and Jimbo in [1, 2] and [3]. Refer to these papers for the general definition and to Rosso [6] for the representation theory of finite dimensional $U_q(\mathfrak{g})$ -modules. As for the quantized enveloping algebra $U_q(\mathfrak{sl}(2))$, recall that it is the Hopf algebra over $\mathbb{C}(q)$ generated by $E, F, K^{\pm 1}$ with the following relations

$$(1.1) \quad \begin{aligned} KEK^{-1} &= q^2 E \\ KFK^{-1} &= q^{-2} F \\ [E, F] &= \frac{K - K^{-1}}{q - q^{-1}} \end{aligned}$$

with comultiplication Δ , antipode S and counit ε given by

$$(1.2) \quad \begin{aligned} \Delta(E) &= E \otimes 1 + K \otimes E \\ \Delta(F) &= F \otimes K^{-1} + 1 \otimes F \\ \Delta(K) &= K \otimes K \end{aligned}$$

$$(1.3) \quad \begin{aligned} S(E) &= -K^{-1} E \\ S(F) &= -FK \\ S(K) &= K^{-1} \end{aligned}$$

$$(1.4) \quad \begin{aligned} \varepsilon(E) &= \varepsilon(F) = 0 \\ \varepsilon(K) &= 1. \end{aligned}$$

(*) Pervenuta in forma definitiva all'Accademia il 23 ottobre 1998.

The results about finite dimensional representation of $\mathcal{U}_q(\mathfrak{sl}(2))$ are completely analogous to those for the enveloping algebra $\mathcal{U}(\mathfrak{sl}(2))$. In particular one has the following

THEOREM 1. (1) *Let V be an irreducible representation of $\mathcal{U}_q(\mathfrak{sl}(2))$. Then the action of both E and F on V is nilpotent, $\dim(\ker E) = \dim(\ker F) = 1$. For any vector $v \in \ker E$, $Kv = \varepsilon q^n v$ with $\varepsilon = \pm 1$ and n the non negative integer such that $\dim(V) = n + 1$. The pair (ε, n) is called the highest weight of V .*

(2) *Given a pair (ε, n) , with ε as above and n a non negative integer, there exists, up to isomorphism, a unique irreducible representation $V_{(\varepsilon, n)}$ of highest weight (ε, n) .*

For the proof see [6] or for example [4, Theorem VI.3.5]. In what follows for the representations of $\mathcal{U}_q(\mathfrak{sl}(2))$ let V_n for $V_{(1, n)}$.

2. THE QUANTUM SYMMETRIC ALGEBRA

We want to give a definition of a quantum analogue of the symmetric algebra of a Lie algebra representation.

Let's start with the classical case. For a representation V of $\mathcal{U}(\mathfrak{g})$ the symmetric algebra $S(V)$ is the quotient of the tensor algebra $T(V)$ by the ideal generated by the antisymmetric component $\Lambda^2(V)$ of the tensor product representation $V \otimes V$. This is exactly the polynomial algebra in the classes x_1, \dots, x_n of the basis vectors v_1, \dots, v_n for V . Since $\Lambda^2(V)$ is a submodule of $V \otimes V$ the Hopf algebra action of $\mathcal{U}(\mathfrak{g})$ on the tensor algebra transfers to the quotient algebra $S(V)$. There is a natural way to characterize the elements of $\Lambda^2(V)$. Since $\mathcal{U}(\mathfrak{g})$ is cocommutative the involutive automorphism σ of $V \otimes V$, given by the switch of the two tensor factors, commutes with the action of $\mathcal{U}(\mathfrak{g})$ on $V \otimes V$. Therefore $V \otimes V$ decomposes as a representation of $\mathcal{U}(\mathfrak{g})$ in two subrepresentation of eigenvalues ± 1 for σ . The antisymmetric tensors are the eigenvectors of eigenvalue -1 , while the symmetric tensors are the eigenvectors of eigenvalue $+1$.

Now for a representation V of the quantum group $\mathcal{U}_q(\mathfrak{g})$ associated to the Lie algebra \mathfrak{g} we define

DEFINITION 2. *The quantum symmetric algebra $S_q(V)$ of V is a \mathbb{Z}^+ -graded algebra over $\mathbb{C}(q)$ with a graded $\mathbb{C}[q, q^{-1}]$ -subalgebra $\overline{S_q(V)}$ such that*

- (1) *the degree-one part of $S_q(V)$ is V ;*
- (2) *$\overline{S_q(V)}$ is free as a graded $\mathbb{C}[q, q^{-1}]$ -module and in each degree has the same dimension of the polynomial algebra in $\dim V$ variables;*
- (3) *the natural map $\overline{S_q(V)} \otimes_{\mathbb{C}[q, q^{-1}]} \mathbb{C}(q) \rightarrow S_q(V)$ is an isomorphism;*
- (4) *$\overline{S_q(V)}/(q-1)$ is the classical symmetric algebra of the $\mathcal{U}(\mathfrak{g})$ -representation corresponding to V , i.e. the polynomial algebra over \mathbb{C} in $\dim V$ variables;*
- (5) *there is an Hopf algebra action of $\mathcal{U}_q(\mathfrak{g})$ on the algebra $S_q(V)$ such that an integral form of $\mathcal{U}_q(\mathfrak{g})$ acts on $\overline{S_q(V)}$ inducing the usual action of $\mathcal{U}(\mathfrak{g})$ on the symmetric algebra $\overline{S_q(V)}/(q-1)$.*

In order to satisfy condition (5) we have to consider the quotient of the tensor algebra $T(V)$ modulo the ideal generated by the $U_q(\mathfrak{g})$ -submodule of $V \otimes V$ corresponding to the antisymmetric component in the classical case.

The composition \check{R} of the R -matrix with the switch σ of the two tensor factors plays in the quantum case the role of σ and commutes with the action of $U_q(\mathfrak{g})$ on $V \otimes V$, which therefore decomposes into irreducible submodules that are eigenspaces for \check{R} . More precisely

PROPOSITION 2 [5, (1.38)]. *Let $V = U^\lambda$ be the irreducible module with highest weight λ for $U_q(\mathfrak{g})$. Then \check{R} satisfies the equation*

$$\prod_{\nu} \left(\check{R} \pm q^{\frac{1}{2}(\nu, \nu + 2\rho) - (\lambda, \lambda + 2\rho)} \right) = 0$$

where ν ranges over the irreducible summands U^ν of $V \otimes V$ and ρ is half the sum of the positive roots. The sign is positive for the U^ν corresponding to the antisymmetric summands and negative for the U^ν corresponding to the symmetric summands.

As a consequence, for an irreducible representation V , we need to take the quotient of the tensor algebra $T(V)$ by the ideal generated by the eigenvectors for \check{R} of eigenvalues with negative sign, although it is not necessarily true that this verifies all the requested properties for $S_q(V)$. On the contrary we shall give in the next section a counterexample to the existence of the quantum symmetric algebra.

3. THE COUNTEREXAMPLE

Let's consider the irreducible $n + 1$ -dimensional representation of $U_q(\mathfrak{sl}(2))$. There is a quantum version of the classical Clebsch-Gordan formula. It gives the decomposition $V_n \otimes V_n \simeq V_{2n} \oplus V_{2n-2} \oplus \cdots \oplus V_0$ into irreducible summands. It is easy to check from the formula for a highest weight vector of V_{2n-2p} (cf. [4, Chapter VIII])

$$(3.1) \quad v^{(2n-2p)} = \sum_{i=0}^p (-1)^i \frac{[p]! [n-p+i]! [n-i]!}{[p-i]! [i]! [n]! [n-p]!} q^{i(n-i+1)} v_i \otimes v_{p-i},$$

where $v_i = F^i v_0$ and v_0 is the highest weight vector for V_n , that the V_{2n-2p} with odd p are the ones with negative eigenvalues for \check{R} . In the case of V_3 the tensor product $V_3 \otimes V_3$ decomposes into irreducible summands $V_6 \oplus V_4 \oplus V_2 \oplus V_0$, and V_4 and V_0 generate the ideal of relations for the quantum symmetric algebra, if it exists. A basis of V_4 is given by the 2-tensors

$$(3.2) \quad \begin{aligned} & v_0 \otimes v_1 - q^3 v_1 \otimes v_0 \\ & v_0 \otimes v_2 + (q^{-1} - q^3) v_1 \otimes v_1 - v_2 \otimes v_0 \\ & v_0 \otimes v_3 + (q^{-1} + q - q^3) v_1 \otimes v_2 + (q^{-2} - 1 - q^2) v_2 \otimes v_1 - q^{-3} v_3 \otimes v_0 \\ & (q + q^{-1}) v_1 \otimes v_3 + (q^{-2} - q^4) v_2 \otimes v_2 - (q + q^{-1}) v_3 \otimes v_1 \\ & (q^2 + q^{-2}) v_2 \otimes v_3 - (q + q^5) v_3 \otimes v_2 \end{aligned}$$

and V_0 has generator given by $v_0 \otimes v_3 - q^3 v_1 \otimes v_2 + q^4 v_2 \otimes v_1 - q^3 v_3 \otimes v_0$. It follows that the quotient algebra is generated on $\mathbb{C}(q)$ by x_0, x_1, x_2, x_3 with the following commutation relations:

$$\begin{aligned}
 (3.3) \quad & x_1 x_0 = q^{-3} x_0 x_1 \\
 & x_2 x_0 = x_0 x_2 + (q^{-1} - q^3) x_1^2 \\
 & x_2 x_1 = \frac{1 + q^2 + q^4 - q^6}{q^3(q^2 + 1)} x_1 x_2 + \frac{q^3 - q^{-3}}{q^3(q^2 + 1)} x_0 x_3 \\
 & x_3 x_0 = \frac{q^4 + q^2 + 1 - q^{-2}}{q^3(q^2 + 1)} x_0 x_3 + \frac{1 - q^6}{q^2(q^2 + 1)} x_1 x_2 \\
 & x_3 x_1 = x_1 x_3 + \frac{1 - q^6}{q(q^2 + 1)} x_2^2 \\
 & x_3 x_2 = q^{-3} x_2 x_3.
 \end{aligned}$$

Now it is not difficult to see that this algebra does not satisfy the property (2) of the definition of the quantum symmetric algebra for V_3 . Let's take the monomial $x_2 x_1 x_0$ and write it in terms of ordered monomials. There are two ways of doing it and they give different results. We obtain

$$(3.4) \quad x_2 x_1 x_0 = \frac{(q^6 - 1)}{q^9(q^2 + 1)} x_0^2 x_3 + \frac{(1 + q^2 + q^4 - q^6)}{q^6(q^2 + 1)} x_0 x_1 x_2 + \frac{(1 - q^4)}{q^4} x_1^3,$$

if we commute first x_1 with x_0 , and

$$\begin{aligned}
 (3.5) \quad x_2 x_1 x_0 &= \frac{(q^6 - 1)(q^4 + q^2 + 1 - q^{-2})}{q^9(q^2 + 1)^2} x_0^2 x_3 + \\
 &+ \frac{-1 + q^2 + 2q^4 + 4q^6 - q^{10} - q^{12}}{q^8(q^2 + 1)^2} x_0 x_1 x_2 + \frac{(1 + q^2 + q^4 - q^6)(1 - q^4)}{q^4(q^2 + 1)} x_1^3
 \end{aligned}$$

if we start commuting x_2 with x_1 .

Then we find the following linear dependence relation on $\mathbb{C}(q)$ between ordered monomials of degree 3

$$(3.6) \quad \frac{(q^6 - 1)^2}{q^{11}(q^2 + 1)^2} x_0^2 x_3 - \frac{(q^6 - 1)^2}{q^8(q^2 + 1)^2} x_0 x_1 x_2 + (q^2 - 1)^2 x_1^3 = 0.$$

REMARK. Clearly for $q = 1$ this is an empty relation.

As a consequence of relation (3.6) the dimension of the degree-three component is less than the dimension of the degree-three component of the polynomial algebra in four variables. We have then finally proved the following

THEOREM 3. *For $\mathcal{U}_q(\mathfrak{g}) = \mathcal{U}_q(\mathfrak{sl}(2))$ and $V = V_3$ the quantum symmetric algebra $S_q(V)$ does not exist.*

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