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**Periodic solutions to a non-linear parametric
differential equation of the third order**

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Equazioni differenziali. — *Periodic solutions to a non-linear parametric differential equation of the third order.* Nota (*) di JAN ANDRES e JAN VORÁČEK, presentata dal Corrisp. R. CONTI.

RIASSUNTO. — Si dimostra un teorema di esistenza di soluzioni periodiche dell'equazione differenziale ordinaria del terzo ordine $x''' + a(t, x, x', x'')x'' + b(t, x, x', x'')x' + h(x) = e(t, x, x', x'')$ con le funzioni a, b e periodiche in t di periodo ω .

The equation to be discussed is

$$(1) \quad x''' + a(t, x, x', x'')x'' + b(t, x, x', x'')x' + h(x) = e(t, x, x', x'')$$

with continuous functions a, b, h, e . Furthermore we assume the ω -periodicity in the variable t of a, b, e .

We proceed by the well-known Leray-Schauder fixed point technique [see e.g. [1]].

The equation (1) results for $p := 1$ from the following family of differential equations depending on the parameter p :

$$(2_p) \quad x''' + p[a(t, x, x', x'')x'' + b(t, x, x', x'')x' + (h(x) - cx)] + cx = pe(t, x, x', x'')$$

with a suitable constant $c \neq 0$.

For the existence of ω -periodic solutions of (1) the following conditions are sufficient:

- (i) All the ω -periodic solutions $x(t)$ of (2_p) and $x'(t), x''(t)$ are for $0 \leq p \leq 1$ a priori bounded by a constant independent of p .

(*) Pervenuta all'Accademia il 21 settembre 1984.

(ii) The linear equation

$$x''' + cx = 0$$

[resulting from (2_p) for $p := 0$] has no ω -periodic solution different from identical zero.

It is clear that the last condition (ii) will be satisfied for any $c \neq 0$.

Putting

$$\omega_1 := \frac{\omega}{2\pi}$$

we have for every ω -periodic function $f(t)$ with square integrable f' and f'' :

$$(3) \quad \int_t^{t+\omega} f'^2(s) \, ds \leq \omega_1^2 \int_t^{t+\omega} f''^2(s) \, ds \quad \text{for all } t$$

[cf. e.g. [2]].

Substituting into $a(t, x, x', x'')$ for x a fixed solution $x(t)$ of (1), we obtain the composed function $a(t, x(t), x'(t), x''(t))$ of the variable t . This function we write simply $a_x(t)$; in the same sense we use the symbols $b_x(t)$, $h_x(t)$, $e_x(t)$.

LEMMA 1. *If the following inequalities*

$$(4) \quad |a(t, x, y, z)| \leq A,$$

$$(5) \quad |b(t, x, y, z)| \leq B.$$

$$(6) \quad |e(t, x, y, z)| \leq E$$

hold for all t, x, y, z , then for any

$$(7) \quad \theta := 1 - \omega_1(A + \omega_1 B) > 0$$

every ω -periodic solution $x(t)$ of (2_p) satisfies the inequality

$$(8) \quad \int_t^{t+\omega} x''^2(s) \, ds \leq \omega \left[\frac{E\omega_1}{\theta} \right]^2 := D_2^2$$

for all t .

Proof. Substituting a fixed $x(t)$ into (2_p) , multiplying the obtained identity by $x'(t)$ and integrating we come to

$$\int_t^{t+\omega} x''^2(s) ds = p \int_t^{t+\omega} [a_x(s) x''(s) x'(s) + b_x(s) x'^2(s) - e_x(s) x'(s)] ds$$

and hence by (4), (5), (6)

$$\int_t^{t+\omega} x''^2(s) ds \leq \int_t^{t+\omega} [A |x''(s) x'(s)| + B x'^2(s) + E |x'(s)|] ds.$$

Thereof by (3) and the Schwarz inequality it results

$$\theta^2 \int_t^{t+\omega} x''^2(s) ds \leq (E \omega_1)^2 \omega$$

for all t , i.e. (8).

COROLLARY. *If all the assumptions of Lemma 1 are fulfilled, we get from (8) and (3) also*

$$(9) \quad \int_t^{t+\omega} x'^2(s) ds \leq \omega_1^2 D_2^2 := D_1^2 \quad \text{for all } t.$$

Since for some $t \leq t_1 \leq t + \omega$ it must be

$$x'(t_1) = 0$$

and consequently

$$x'(s) = \int_{t_1}^s x''(u) du,$$

we obtain from the Schwarz inequality that

$$(10) \quad |x'(s)| \leq \sqrt{\omega} D_2 := D'$$

holds for every real s .

LEMMA 2. *If all the assumptions of Lemma 1 are fulfilled and if moreover there exist such real numbers $c \neq 0$, $h > 0$ that*

$$(11) \quad xh(x) \operatorname{sgn} c \leq |c| x^2 \quad \text{for every } |x| \geq h$$

holds, then every ω -periodic solution $x(t)$ of (2_p) satisfies

$$(12) \quad |x(t)| < R + |D'| \omega := D$$

with

$$(13) \quad R := \max \left[h, \frac{A |D_2| + B |D_1|}{c \sqrt{\omega}} + \frac{E}{c} \right]$$

and henceforth also

$$(14) \quad \int_t^{t+\omega} x^2(s) ds < \omega D^2 := D_0^2, \quad \text{for every } t.$$

Proof. We substitute again $x(t)$ into (2_p) . Multiplying the resulting identity by $x(t)$ and integrating, we obtain for every t :

$$\begin{aligned} p \int_t^{t+\omega} [a_x(s) x''(s) + b_x(s) x'(s) - e_x(s)] x(s) ds &= (p-1) \int_t^{t+\omega} c x^2(s) ds - \\ &- p \int_t^{t+\omega} x(s) h_x(s) ds \end{aligned}$$

and further [cf. (11) and (4)-(6)]

$$\begin{aligned} (15) \quad &\int_t^{t+\omega} [(1-p) |c| x^2(s) + p x(s) h_x(s) \operatorname{sgn} c] ds \leq \\ &\leq \int_t^{t+\omega} [A |x''(s) x(s)| + B |x'(s) x(s)| + E |x(s)|] ds. \end{aligned}$$

If on $\langle t, t+\omega \rangle$ the inequality $|x(s)| > R$ holds, we have

$$c^2 R^2 \omega < c^2 \int_t^{t+\omega} x^2(s) ds \leq (A |D_2| + B |D_1| + E \sqrt{\omega})^2$$

and this contradicts to (13). Thus on $\langle t, t+\omega \rangle$ there must exist a point t_1 with $|x(t_1)| \leq R$ and the lemma is proved.

LEMMA 3. *If all the assumptions of the foregoing lemma are valid, then denoting*

$$H := \max_{|x| \leq D} |h(x)|$$

we have for an arbitrary ω -periodic solution $x(t)$ of (2_p)

$$(16) \quad \int_t^{t+\omega} x''''(s) ds \leq (A |D_2| + B |D_1| + (H + E) \sqrt{\omega})^2 := D_3^2$$

and thus [cf. (10)]

$$(17) \quad |x''(t)| \leq \sqrt{\omega} D_3 := D''$$

for every t .

The proof of this Lemma may be performed by integration of the identity resulting from (2_p) multiplied by $x'''(t)$ similarly as in Lemma 1 and Lemma 2.

THEOREM. *If the conditions (4)-(7) are satisfied and if there exist such reals $c \neq 0$, $h > 0$ that (11) holds, then the equation (1) admits an ω -periodic solution.*

Proof. By Lemmas 1, 2, 3 we have for ω -periodic solution $x(t)$ of (2_p)

$$|x(t)| + |x'(t)| + |x''(t)| \leq D + |D'| + |D''| := P$$

with P independent on $p \in \langle 0, 1 \rangle$.

Thus both conditions (i), (ii) sufficient for the existence of an ω -periodic solution of (1) are fulfilled.

Remark. Simplifying (1) into the equation

$$(18) \quad x''' + [a(t)x']' + g(x)x' + h(x) = e(t)$$

with ω -periodic $a(t)$ and $e(t)$, it is possible to require instead of (11) other conditions less restrictive with respect to $h(x)$.

The equation (18) has an ω -periodic solution, if besides (4), (5), (7) the following relations are satisfied:

$$(19) \quad h(x) \operatorname{sgn} x \geq 0 \quad \text{for every } |x| \geq h,$$

$$(20) \quad \int_0^\omega e(t) dt = 0.$$

Indeed, assuming $\operatorname{sgn} x(t) = \operatorname{const.}$ on $\langle t, t + \omega \rangle$, we obtain from (18_p) [cf. (2_p)] and (20):

$$\int_t^{t+\omega} [h(x(s)) \operatorname{sgn} x(s) + (1-p)c |x(s)|] ds = 0$$

for any c ; nevertheless (19) implies for $|x(t)| \geq h$ on $\langle t, t + \omega \rangle$ and a suitable $c \neq 0$:

$$\int_t^{t+\omega} [h(x(s)) \operatorname{sgn} x(s) + (1-p)c |x(s)|] ds \neq 0.$$

Hence there must be a point $t_1 \in \langle t, t + \omega \rangle$ such that $|x(t_1)| < h$.

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