
ATTI ACCADEMIA NAZIONALE DEI LINCEI
CLASSE SCIENZE FISICHE MATEMATICHE NATURALI
RENDICONTI

JAMES O.C. EZEILO

**On the existence of periodic solutions of certain
third order non-dissipative differential systems**

*Atti della Accademia Nazionale dei Lincei. Classe di Scienze Fisiche,
Matematiche e Naturali. Rendiconti, Serie 8, Vol. 66 (1979), n.2, p. 126–135.*

Accademia Nazionale dei Lincei

[<http://www.bdim.eu/item?id=RLINA_1979_8_66_2_126_0>](http://www.bdim.eu/item?id=RLINA_1979_8_66_2_126_0)

L'utilizzo e la stampa di questo documento digitale è consentito liberamente per motivi di ricerca e studio. Non è consentito l'utilizzo dello stesso per motivi commerciali. Tutte le copie di questo documento devono riportare questo avvertimento.

*Articolo digitalizzato nel quadro del programma
bdim (Biblioteca Digitale Italiana di Matematica)
SIMAI & UMI*

<http://www.bdim.eu/>

Equazioni differenziali ordinarie. — *On the existence of periodic solutions of certain third order non-dissipative differential systems.*
 Nota di JAMES O. C. EZEILO, presentata (*) dal Socio G. SANSONE.

RIASSUNTO. — Si danno condizioni sufficienti perché l'equazione $\ddot{X} + A\dot{X} + B\dot{X} + H(X) = P(t, X, \dot{X}, \ddot{X})$, con A, B, C matrici $n \times n$ simmetriche con elementi costanti $P(t + \omega, X, Y, Z) = P(t, X, Y, Z)$ ammetta almeno una soluzione periodica.

1. We shall be concerned here with real third order differential systems of the form:

$$(1.1) \quad \ddot{X} + A\dot{X} + B\dot{X} + H(X) = P(t, X, \dot{X}, \ddot{X})$$

in which A, B are constant symmetric $n \times n$ matrices and X, H, P are n -vectors with H, P dependent only on the arguments shown. It will be assumed as basic throughout that $P(t, X, Y, Z)$ is continuous and ω -periodic in t (that is $P(t + \omega, X, Y, Z) = P(t, X, Y, Z)$ for some $\omega > 0$), and that there are constants $\delta \geq 0, \varepsilon \geq 0$ such that

$$(1.2) \quad \|P(t, X, Y, Z)\| \leq \delta + \varepsilon(\|X\| + \|Y\| + \|Z\|)$$

for arbitrary t, X, Y and Z . Here and elsewhere the symbol $\|\cdot\|$ denotes the Euclidean norm. The derivatives $\partial h_i / \partial x_j, h_i$ and x_j ($1 \leq i, j \leq n$) here and elsewhere being the components of H and X respectively, are also assumed continuous, with the (Jacobian) matrix $J_h(X) \equiv (\partial h_i / \partial x_j)$ symmetric, for arbitrary X .

Let $\lambda_i(A), \lambda_i(B), \lambda_i(J_h(X))$ denote the eigenvalues (all real) of A, B, J_h respectively and let $\alpha \equiv \max_{1 \leq i \leq n} \lambda_i(A), \beta \equiv \max_{1 \leq i \leq n} \lambda_i(B)$. The following result, extending an earlier (scalar) result in [1], was announced, but without proof, at the International Congress Mathematicians in Helsinki (in August 1978):

THEOREM. *There exists a constant $\varepsilon_0 = \varepsilon_0(\delta, A, B, H) > 0$ such that if*

$$(1.3) \quad \gamma_0 \equiv \inf_{i, X} \lambda_i(J_h(X)) > \begin{cases} 0, & \text{if one at least of } \alpha, \beta \text{ is non} \\ & \text{positive,} \\ \alpha\beta, & \text{if } \alpha \text{ and } \beta \text{ are both positive,} \end{cases}$$

then (1.1) has at least one ω -periodic solution provided that $\varepsilon \leq \varepsilon_0$.

(*) Nella seduta del 10 febbraio 1979.

The object of the present note is to supply now a detailed proof of this theorem.

The reference here, in the title, to the system (1.1) as non-dissipative stems from the condition (1.3) which is clearly "non Routh-Hurwitz". Note that there is no loss in generality in assuming that $H(o) = o$; for the subtraction of $H(o)$ from either side of (1.1) gives an equation with H, P replaced by H_0, P_0 , where $H_0(X) = H(X) - H(o)$ which satisfies $H_0(o) = o$ and $P_0 = P - H(o)$ which satisfies

$$\|P_0(t, X, Y, Z)\| \leq \delta_0 + \varepsilon (\|X\| + \|Y\| + \|Z\|)$$

with $\delta_0 \equiv \delta + \|H(o)\|$, which is the same as (1.2).

2. NOTATION. In what follows we shall use γ 's with or without suffixes to denote positive constants whose magnitudes depend only on δ, A, B and H . The γ 's without suffixes are not necessarily the same in each place of occurrence but the numbered γ 's: $\gamma_0, \gamma_1, \gamma_2, \dots$ retain a fixed magnitude throughout.

Next, given any pair of vectors, X and Y say, with components (x_1, \dots, x_n) and (y_1, \dots, y_n) respectively, we shall use $\langle X, Y \rangle$ to denote their scalar product $\sum_{i=1}^n x_i y_i$. Thus, in particular $\langle X, X \rangle \equiv \|X\|^2$.

3. The proof is by the Leray-Schauder technique, with (1.1) embedded in the parameter-dependent equation:

$$(3.1) \quad \ddot{X} + \mu (A\ddot{X} + B\dot{X}) + (1 - \mu) \gamma_0 X + \mu H = \mu P$$

where the parameter μ is as usual restricted to the closed range $[0, 1]$. Note that, when $\mu = 0$, (3.1) reduces to the equation

$$\ddot{X} + \gamma_0 X = 0$$

which clearly has no non-trivial ω -periodic solution. Also, when $\mu = 1$, (3.1) reduces to the equation (1.1). Thus the theorem will follow from the usual Leray-Schauder fixed point considerations (see for example theorem 1.39 of [3]) if it can be shown that there are constants $\gamma_1, \gamma_2, \gamma_3$ all independent of μ , such that

$$(3.2) \quad \|X\| \leq \gamma_1, \|\dot{X}\| \leq \gamma_2 \quad \text{and} \quad \|\ddot{X}\| \leq \gamma_3 \quad (0 \leq t \leq \omega)$$

for every ω -periodic solution of (3.1) corresponding to $0 \leq \mu \leq 1$.

4. *Preliminary lemmas.* We shall make occasional use of the following lemmas:

LEMMA 1. *Let D be a symmetric $n \times n$ matrix and X any n-vector. Then*

$$(4.1) \quad d_1 \|X\|^2 \leq \langle X, DX \rangle \leq d_2 \|X\|^2$$

where d_1, d_2 are respectively the least and the greatest of the eigenvalues of D.

This is a well known result (see for example [2; p. 288]).

LEMMA 2. *If $X = X(t)$ is twice continuously differentiable in t , then*

$$(4.2) \quad \int \langle \ddot{X}, H(X) \rangle dt = \langle \dot{X}, H(X) \rangle - \int \langle J_h(X) \dot{X}, \dot{X} \rangle dt,$$

the integrals here being indefinite integrals.

Proof. Since $\langle \ddot{X}, H(X) \rangle \equiv \sum_{i=1}^n \ddot{x}_i h_i$ we have on integrating by parts, that

$$\begin{aligned} \int \langle \ddot{X}, H(X) \rangle dt &= \sum_{i=1}^n \dot{x}_i h_i - \int \sum_{i=1}^n \dot{x}_i \frac{dh_i}{dt} dt \\ &= \langle \dot{X}, H \rangle - \int \sum_{i=1}^n \sum_{k=1}^n \dot{x}_i \frac{\partial h_i}{\partial x_k} \dot{x}_k dt \\ &= \langle \dot{X}, H \rangle - \int \langle J_h(X) \dot{X}, \dot{X} \rangle dt \end{aligned}$$

which establishes (4.2).

Throughout what follows $X = X(t)$ denotes an arbitrary ω -periodic solution of (3.1) with μ restricted always to be the range $0 \leq \mu \leq 1$. The objective now will be to establish (3.2).

The main tool is the scalar function $u = u(t)$ given by

$$(5.1) \quad u = \frac{1}{2} b_1 \langle \ddot{X}, \dot{X} \rangle - b_2 \langle X, \ddot{X} \rangle + \langle \dot{X}, \ddot{X} \rangle$$

where $b_1 > 0$, $b_2 > 0$ are constants whose values are as yet undetermined but will be fixed to advantage as γ 's in the course of the proof. We have, by an elementary differentiation with respect to t that

$$\begin{aligned} \dot{u} &= -b_1 \langle \ddot{X}, \mu A \ddot{X} + \mu B \dot{X} + (1 - \mu) \gamma_0 X + \mu H - \mu P \rangle - b_2 \langle \dot{X}, \ddot{X} \rangle + \\ &\quad + b_2 \langle X, \mu A \ddot{X} + \mu B \dot{X} + (1 - \mu) \gamma_0 X + \mu H - \mu P \rangle + \langle \ddot{X}, \ddot{X} \rangle - \\ &\quad - \langle \dot{X}, \mu A \ddot{X} + \mu B \dot{X} + (1 - \mu) \gamma_0 X + \mu H - \mu P \rangle. \end{aligned}$$

Note that the terms

$$\langle \ddot{X}, B \dot{X} \rangle, \langle X, B \dot{X} \rangle, \langle \dot{X}, A \ddot{X} \rangle$$

which occur on the right hand side of (5.2) are perfect t -differentials since A , B , are symmetric. Also, since $J_h(X)$ is symmetric we have from equation 2.4 (3) of [2] that

$$\langle \dot{X}, H(X) \rangle = \frac{d}{dt} \int_0^1 \langle H(\sigma X), X \rangle d\sigma$$

so that the term $\langle \dot{X}, H(X) \rangle$ which occurs on the same right hand side of (5.2) is also a perfect t -differential. Thus we may indeed reset (5.2) in the form

$$(5.3) \quad \dot{u} \equiv u_1 + u_2 + u_3$$

where

$$\begin{aligned} u_1 &= \{ \langle \ddot{X}, \ddot{X} \rangle - \mu b_1 \langle \ddot{X}, A\ddot{X} \rangle + \\ &+ \{ \mu b_2 \langle X, A\ddot{X} \rangle - b_1 \langle \ddot{X}, (I - \mu) \gamma_0 X + \mu H \rangle - \mu \langle \dot{X}, B\dot{X} \rangle \} + \\ &+ \{ b_2 \langle X, (I - \mu) \gamma_0 X + \mu H \rangle \} \equiv \\ &\equiv u_{11} + u_{12} + u_{13}, \end{aligned}$$

say,

$$(5.4) \quad u_2 = -\mu \langle b_1 \ddot{X} + \dot{X} - b_2 X, P \rangle$$

and u_3 is a perfect t -differential. Hence, integrating both sides of (5.3) with respect to t from $t = 0$ to $t = \omega$, we have, X being ω -periodic, that

$$(5.5) \quad \int_0^\omega (u_{11} + u_{12} + u_{13}) dt + \int_0^\omega u_2 dt = 0.$$

Now, by Lemma I,

$$\langle \ddot{X}, A\ddot{X} \rangle \leq \alpha \| \ddot{X} \|^2$$

so that, since $0 \leq \mu \leq 1$,

$$(5.6) \quad u_{11} \geq (1 - \alpha b_1) \| \ddot{X} \|^2.$$

Next, since $H(0) = 0$, we have from equation 2.2 (3) of [2] that $H(X) = \int_0^1 J_h(\sigma X) X d\sigma$ so that, in particular

$$\begin{aligned} \langle X, H(X) \rangle &= \int_0^1 \langle X, J_h(\sigma X) X \rangle d\sigma \\ &\geq \gamma_0 \| X \|^2 \end{aligned}$$

by (1.3) and (4.1); and hence

$$(5.7) \quad u_{13} \geq b_2 \gamma_0 \| X \|^2$$

Finally, we have by Lemma 2 that

$$(5.8) \quad \int_0^{\omega} \langle X, A\ddot{X} \rangle dt = - \int_0^{\omega} \langle A\dot{X}, \dot{X} \rangle dt \\ \geq -\alpha \int_0^{\omega} \|\dot{X}\|^2 dt$$

in view of (4.1), and then analogously for the terms $\int_0^{\omega} \langle \ddot{X}, X \rangle dt$ and $\int_0^{\omega} \langle \ddot{X}, H(X) \rangle dt$ appearing in $\int_0^{\omega} u_{12} dt$ that

$$(5.9) \quad \int_0^{\omega} \langle \ddot{X}, X \rangle dt = - \int_0^{\omega} \|\dot{X}\|^2 dt$$

$$(5.10) \quad \int_0^{\omega} \langle \ddot{X}, H(X) \rangle dt = - \int_0^{\omega} \langle J_h(X) \dot{X}, \dot{X} \rangle dt \\ \leq -\gamma_0 \int_0^{\omega} \|\dot{X}\|^2 dt,$$

the latter inequality deriving immediately from the use of (1.3) and (4.1). Since

$$(5.11) \quad \langle \dot{X}, B\dot{X} \rangle \leq \beta \|\dot{X}\|^2$$

it is clear from (5.8), (5.9), (5.10) and (5.11) that

$$(5.12) \quad \int_0^{\omega} u_{12} dt \geq (b_1 \gamma_0 - b_2 \alpha - \beta) \int_0^{\omega} \|\dot{X}\|^2 dt$$

Thus we have from (5.6), (5.7) and (5.12) that

$$(5.13) \quad \int_0^{\omega} (u_{11} - u_{12} + u_{13}) dt \geq \int_0^{\omega} u_4 dt$$

where

$$u_4 \equiv (1 - \alpha b_1) \|\ddot{X}\|^2 + (b_1 \gamma_0 - b_2 \alpha - \beta) \|\dot{X}\|^2 + b_2 \gamma_0 \|X\|^2.$$

A most crucial part of our proof is to show now that the so far undefined positive constants b_1, b_2 in (5.1) can in fact be fixed such that

$$(5.14) \quad u_4 \geq \gamma_4 (\| \ddot{X} \|^2 + \| \dot{X} \|^2 + \| X \|^2)$$

for some γ_4 . We shall distinguish here two cases (already highlighted in (1.3)) namely: (I) *one at least of α, β is non positive*, (II) *α and β are both positive*.

We start with the case (I). Suppose for example that $\alpha \leq 0$. Then clearly $(1 - \alpha b_1) \geq 1$ for arbitrary $b_1 > 0$.

Also

$$b_1 \gamma_0 - b_2 \alpha - \beta \geq \gamma_5$$

if

$$(5.15) \quad b_1 \geq (\gamma_5 + |\beta|) \gamma_0^{-1},$$

for arbitrary $b_2 > 0$. Thus when $\alpha \leq 0$ we have that

$$u_4 \geq (\| \ddot{X} \|^2 + \gamma_5 \| \dot{X} \|^2 + \gamma_0 \gamma_6 \| X \|^2)$$

if b_1 is fixed by (5.15) and $b_2 = \gamma_6$, which establishes (5.14) with $\gamma_4 = \min(1, \gamma_5, \gamma_0 \gamma_6)$. Suppose on the other hand that $\beta \leq 0$. Then, if $\alpha \leq 0$, $b_1 = \gamma_7 = b_2$ clearly secures the estimate:

$$u_4 \geq \| \ddot{X} \|^2 + \gamma_0 \gamma_7 (\| \dot{X} \|^2 + \| X \|^2)$$

which implies (5.14) (with $\gamma_4 = \min(1, \gamma_0 \gamma_7)$) while the choice

$$b_1 = \frac{1}{2} \alpha^{-1}, b_2 = \frac{1}{4} \gamma_0 \alpha^{-2}$$

when $\alpha > 0$ secures the estimate:

$$u_4 \geq \frac{1}{2} (\| \ddot{X} \|^2 + \frac{1}{2} \gamma_0 \alpha^{-1} \| \dot{X} \|^2 + \frac{1}{2} \gamma_0^2 \alpha^{-2} \| X \|^2)$$

which again implies (5.14) but with $\gamma_4 = \frac{1}{2} \min(1, \frac{1}{2} \gamma_0 \alpha^{-1}, \frac{1}{2} \gamma_0^2 \alpha^{-2})$. Thus whether $\alpha \leq 0$ or $\beta \leq 0$ it is possible to fix $b_1 = \gamma, b_2 = \gamma$ so that (5.14) holds.

We turn next to the case (II): $\alpha > 0$ and $\beta > 0$. Note that, since $\gamma_0 > \alpha\beta$, by (1.3), it is possible to choose γ_8 such that

$$(5.16) \quad \beta \gamma_0^{-1} < \gamma_8 < \alpha^{-1}$$

Now fix $b_1 = \gamma_8$ and $b_2 \equiv \frac{1}{2} \alpha^{-1} (\gamma_0 \gamma_8 - \beta) > 0$, by (5.16). Then

$$\begin{aligned} u_4 &\geq (1 - \alpha \gamma_8) \| \ddot{X} \|^2 + \frac{1}{2} (\gamma_8 \delta_0 - \beta) \| \dot{X} \|^2 + \frac{1}{2} \alpha^{-1} (\gamma_0 \gamma_8 - \beta) \| X \|^2 \\ &\geq \gamma (\| \ddot{X} \|^2 + \| \dot{X} \|^2 + \| X \|^2) \end{aligned}$$

for some γ , since $(1 - \alpha \gamma_8)$ and $(\gamma_8 \gamma_0 - \beta)$ are both positive, by (5.16). Thus in the case (II), (5.14) holds for some appropriate choice of b_1, b_2 as γ 's. We

have thus conclusively verified that, subject to (1.3), there exist γ_9, γ_{10} such that if

$$(5.17) \quad b_1 = \gamma_9, b_2 = \gamma_{10}$$

then (5.14) holds.

We assume henceforth that b_1 and b_2 are fixed by (5.17) and define $\rho = \rho(t) \geq 0$ by

$$\rho^2 = \|\ddot{X}\|^2 + \|\dot{X}\|^2 + \|X\|^2.$$

It is clear then from (5.14), (5.13), (5.5), (5.3) and (1.2) that

$$\gamma_4 \int_0^\omega \rho^2 dt \leq \gamma_{11} \int_0^\omega \rho dt + \varepsilon \gamma_{12} \int_0^\omega \rho^2 dt$$

for some γ_{11} and γ_{12} ; so that if, for example,

$$(5.18) \quad \varepsilon \leq \frac{1}{2} \gamma_4 \gamma_{12}^{-1}$$

as we assume henceforth and $\gamma_{13} \equiv 2 \gamma_4^{-1} \gamma_{11}$ then

$$\begin{aligned} \int_0^\omega \rho^2 dt &\leq \gamma_{13} \int_0^\omega \rho dt \\ &\leq \gamma_{13} \omega^{\frac{1}{2}} \left(\int_0^\omega \rho^2 dt \right)^{1/2} \end{aligned}$$

by Schwarz's inequality. Hence

$$\left(\int_0^\omega \rho^2 dt \right)^{1/2} \leq \gamma_{13} \omega^{\frac{1}{2}}$$

that is

$$(5.19) \quad \int_0^\omega \rho^2 dt \leq \gamma_{14} \equiv \gamma_{13}^2 \omega.$$

The result (3.2) is a consequence of (5.19) as will now be shown. We begin by noting that (5.19) implies that

$$(5.20) \quad \int_0^\omega x_i^2 dt \leq \gamma_{14}, \int_0^\omega \dot{x}_i^2 dt \leq \gamma_{14}, \int_0^\omega \ddot{x}_i^2 dt \leq \gamma_{14} \quad (i = 1, 2, \dots, n).$$

The inequality: $\int_0^\omega \dot{x}_i^2 dt \leq \gamma_{14}$ here implies that $|x_i(\tau)| \leq \gamma_{15} \equiv \gamma_{14}^{\frac{1}{2}} \omega^{-\frac{1}{2}}$ for some $\tau \in [0, \omega]$. Thus, since $x_i(t) = x_i(\tau) + \int_\tau^t \dot{x}_i(s) ds$, we have at once that

$$\begin{aligned} \text{Sup}_{0 \leq t \leq \omega} |x_i(t)| &\leq \gamma_{15} + \int_\tau^{\tau+\omega} |\dot{x}_i(s)| ds \\ &\leq \gamma_{15} + \omega^{\frac{1}{2}} \left(\int_\tau^{\tau+\omega} \dot{x}_i^2(s) ds \right)^{1/2}, \end{aligned}$$

by Schwarz's inequality, which, in view of the second inequality in (5.20), leads in turn to the estimate

$$\text{Sup}_{0 \leq t \leq \omega} |\dot{x}_i(t)| \leq \gamma_{15} + \omega^{\frac{1}{2}} \gamma_{14}^{\frac{1}{2}}.$$

This is true for each $i = 1, 2, \dots, n$ and hence

$$(5.21) \quad \|X\| \leq \gamma_{16} \quad (0 \leq t \leq \omega)$$

for each ω -periodic solution $X(t)$ of (3.1) corresponding to $0 \leq \mu \leq 1$. Analogously the second and third inequalities in (5.20) also lead to the estimate

$$\text{Sup}_{0 \leq t \leq \omega} |\dot{x}_i(t)| \leq \gamma_{15} + \omega^{\frac{1}{2}} \gamma_{14}^{\frac{1}{2}} \quad (i = 1, 2, \dots, n)$$

which in turn implies that

$$(5.22) \quad \|\dot{X}\| \leq \gamma_{17} \quad (0 \leq t \leq \omega)$$

for each ω -periodic solution $X(t)$ of (3.1) corresponding to $0 \leq \mu \leq 1$. It should be pointed out, however, that the middle inequality in (5.20), whose only role, as far as the verification of (5.22) is concerned, is to secure the existence of a $\tau \in [0, \omega]$ such that $|\dot{x}_i(\tau)| \leq \gamma_{15}$ is not actually crucial to the proof of (5.22) once the last inequality in (5.20) is available. This is because the existence of a $\tau \in [0, \omega]$ such that $|\dot{x}_i(\tau)| \leq \gamma$ for some γ is already a consequence of the ω -periodicity condition: $x_i(0) = x_i(\omega)$ which in fact implies that $\dot{x}_i(\tau_0) = 0$ for some $\tau_0 \in [0, \omega]$, so that because of the identity:

$$\dot{x}_i(t) = \dot{x}_i(\tau_0) + \int_{\tau_0}^t \ddot{x}_i(s) ds$$

we have that

$$(5.23) \quad \begin{aligned} \text{Sup}_{0 \leq t \leq \omega} |x_1(t)| &\leq \omega^{\frac{1}{2}} \left(\int_{\tau_0}^{\tau_0 + \omega} \ddot{x}_i^2(s) ds \right)^{1/2} \\ &\leq \omega^{\frac{1}{2}} \gamma_{14}^{\frac{1}{2}}, \end{aligned} \quad (i = 1, 2, \dots, n),$$

thus leading to: $\|\dot{X}\| \leq \gamma$ ($0 \leq t \leq \omega$) as before.

To establish the last of the inequalities (3.2) it will suffice now to verify that

$$(5.24) \quad \int_0^{\omega} \|\ddot{X}\|^2 dt \leq \gamma_{18}$$

for any ω -periodic solution of (3.1) with $0 \leq \mu \leq 1$. For if indeed (5.24) holds, so that

$$(5.25) \quad \int_0^{\omega} \ddot{x}_i^2 dt \leq \gamma_{18} \quad (i = 1, 2, \dots, n),$$

then, since $\ddot{x}_i(\tau_1) = 0$ for some $\tau_1 \in [0, \omega]$ so that

$$\ddot{x}_i(t) = \ddot{x}_i(\tau_1) + \int_{\tau_1}^t \ddot{x}(s) ds = \int_{\tau_1}^t \ddot{x}(s) ds,$$

we shall have that

$$(5.26) \quad \begin{aligned} \text{Sup}_{0 \leq t \leq \omega} |\ddot{x}_i(t)| &\leq \omega^{\frac{1}{2}} \left(\int_{\tau_1}^{\tau_1 + \omega} \ddot{x}^2(s) ds \right)^{1/2} \\ &\leq \omega^{\frac{1}{2}} \gamma_{18}^{\frac{1}{2}} \end{aligned} \quad (i = 1, 2, \dots, n),$$

by (5.25), which leads to the remaining estimate:

$$(5.27) \quad \|\ddot{X}\| \leq \gamma_{19} \quad (0 \leq t \leq \omega)$$

in (3.2). As for the actual verification of (5.24) it is convenient to take a scalar product of either side of (3.1) with \ddot{X} and integrate with respect to t from $t = 0$ to $t = \omega$. Since X, \dot{X} are already subject to (5.21) and (5.22) and $\langle A\ddot{X}, \ddot{X} \rangle$ is a perfect t -differential, this integration shows readily in view of (1.2), that

$$\begin{aligned} \int_0^{\omega} \|\ddot{X}\|^2 dt &\leq \gamma_{20} \int_0^{\omega} \|\ddot{X}\| dt + \varepsilon \int_0^{\omega} \|\ddot{X}\| \cdot \|\ddot{X}\| dt \\ &\leq \left\{ \gamma_{10} \omega^{\frac{1}{2}} + \varepsilon \left(\int_0^{\omega} \|\ddot{X}\|^2 dt \right)^{1/2} \right\} \left(\int_0^{\omega} \|\ddot{X}\|^2 dt \right)^{1/2}. \end{aligned}$$

Thus, since $\int_0^{\omega} \|\ddot{\mathbf{X}}\|^2 dt \leq \gamma$, we must have that

$$(5.28) \quad \int_0^{\omega} \|\ddot{\mathbf{X}}\|^2 dt \leq \gamma_{21} \left(\int_0^{\omega} \|\ddot{\mathbf{X}}\|^2 dt \right)^{1/2}.$$

Hence

$$\int_0^{\omega} \|\ddot{\mathbf{X}}\|^2 dt \leq \gamma_{21}^2$$

which is (5.24). Thus (5.27) holds if $\varepsilon \leq \frac{1}{2} \gamma_4 \gamma_{12}^{-1}$. This completely verifies the theorem with $\varepsilon_0 = \frac{1}{2} \gamma_4 \gamma_{12}^{-1}$.

REFERENCES

- [1] J. O. C. EZEILO (1975) - «Math. Proc. Cambridge Philos. Soc.», 77, 547-551.
- [2] J. O. C. EZEILO and H. O. TEJUMOLA (1966) - «Ann. Mat. Pura Appl.», IV, Vol. LXXIV, 283-316.
- [3] R. REISSIG, R. CONTI and G. SANSONE (1969) - *Nichtlineare Differentialgleichungen höherer Ordnung*, Edizioni Cremonese, Rome.