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## On a special recurrent Finsler space of second order

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Geometria differenziale. - On a special recurrent Finsler space of second order. Nota di Awdhesh Kumar, presentata (*) dal Socio B. Segre.

Riassunto. - Il presente lavoro si ricollega alla ricerca [4] dello stesso Autore, dedicata allo studio dei moti affini negli spazi ricorrenti di Finsler del secondo ordine. Qui si tratta di tali spazi, sotto le condizioni (2.1), (2.2) e (2.3).

## Introduction

Let us consider an $n$-dimensional affinely connecting Finsler space $\mathrm{F}_{n}$ [I] ${ }^{(1)}$ equipped with $2 n$ line elements and a fundamental metric function $\mathrm{F}(x, \dot{x})$ which is positively homogeneous of degree one in its directional arguments. The fundamental metric tensor of the space is given by

$$
\begin{equation*}
g_{i j}(x, \dot{x}) \stackrel{\text { def }}{=} \frac{1}{2} \dot{\partial}_{i} \dot{\partial}_{j} \mathrm{~F}^{2}(x, \dot{x}), \quad \quad \dot{\partial}_{i} \equiv \partial / \dot{x}^{i} \tag{I.I}
\end{equation*}
$$

and is symmetric in its lower indices.
Let us further consider a mixed tensor field $\mathrm{T}_{j}^{i}(x, \dot{x})$ which depends both upon positional and directional arguments. The covariant derivative of $\mathrm{T}_{j}^{i}(x, \dot{x})$ with respect to $x^{k}$ in the sense of Cartan is given by

$$
\begin{equation*}
\mathrm{T}_{j \mid k}^{i}=\partial_{k} \mathrm{~T}_{j}^{i}-\dot{\partial}_{m} \mathrm{~T}_{j}^{i} \mathrm{G}_{k}^{m}+\mathrm{T}_{j}^{s} \Gamma_{s k}^{* i}-\mathrm{T}_{s}^{i} \Gamma_{j, k}^{* s}, \tag{1.2}
\end{equation*}
$$

where $\Gamma_{h k}^{* i}(x, \dot{x})$ are connection coefficients and are also symmetric in their lower indices $j$ and $k$.

The commutation formula involving the above covariant derivative for the tensor field $\mathrm{T}_{j}^{i}(x, \dot{x})$ is given by

$$
\begin{equation*}
2 \mathrm{~T}_{j[[h k]}^{i}=-\dot{\partial}_{r} \mathrm{~T}_{j}^{i} \mathrm{~K}_{h k}^{r}+\mathrm{T}_{j}^{s} \mathrm{~K}_{s h k}^{i}-\mathrm{T}_{s}^{i} \mathrm{~K}_{j h k}^{s}{ }^{(2)} \tag{I.3}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathrm{K}_{h j k}^{i}(x, \dot{x}) \stackrel{\text { def }}{=} 2\left\{\hat{a}_{[k} \Gamma_{j j h}^{* i}-\dot{\partial}_{r} \Gamma_{h[j}^{* i} \mathrm{G}_{k]}^{r}+\Gamma_{h[j}^{* r} \Gamma_{k] r r}^{* i}\right\} \tag{I.4}
\end{equation*}
$$

is called Cartan's curvature tensor field and satisfies the following identities [I]:

$$
\begin{equation*}
\mathrm{K}_{h j k}^{i}+\mathrm{K}_{j k h}^{i}+\mathrm{K}_{k h j}^{i}=\mathrm{o} \tag{I.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\text { a) } \mathrm{K}_{h j k}^{i}=-\mathrm{K}_{h k j}^{i} \quad \text { and } \quad \text { b) } \quad \mathrm{K}_{h j i}^{i}=\mathrm{K}_{h j} \tag{1.6}
\end{equation*}
$$

(*) Nella seduta del 13 novembre 1976.
(I) The numbers in square brackets refer to the References given at the end of the paper.
(2) $2 \mathrm{~A}_{[h k]}=\mathrm{A}_{h k}-\mathrm{A}_{k h}$.

## 2. Recurrent Finsler space of second order

Definition (2.1). An $n$-dimensional affinely connected $\mathrm{F}_{n}$ is called recurrent. Finsler space of the second order if its curvature tensor field $\mathrm{K}_{h j k}^{i}(x, \dot{x})$ satisfies the relation:

$$
\begin{equation*}
\mathrm{K}_{h j k \mid s m}^{i}=\alpha_{s m} \mathrm{~K}_{h j k}^{i} \tag{2.1}
\end{equation*}
$$

where $\alpha_{s, n}(x, \dot{x})$ is a non-zero symmetric tensor field.
In what follows, we shall assume the following two conditions:

$$
\begin{equation*}
v_{\mid h}^{i}=\varepsilon_{h} v^{i} \tag{2.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{K}_{h j}=\varepsilon_{h} \eta_{j} \tag{2.3}
\end{equation*}
$$

where $\eta_{j}$ is a suitable covariant vector and $\mathrm{K}_{h j}(x, \dot{x})$ is the so called Ricci tensor given by ( I .6 b ). In the following, we shall study the fundamental properties of the space under the conditions (2.1), (2.2) and (2.3).

Taking the covariant derivative of (2.2) with respect to $x^{s}$ and using the same equation (2.2), we have

$$
\begin{equation*}
v_{\mid h s}^{i}=v^{i}\left(\varepsilon_{h \mid s}+\varepsilon_{h} \varepsilon_{s}\right) \tag{2.4}
\end{equation*}
$$

By virtue of the commutation formula (I.3), commutating (2.4) with respect to the indices $h$ and $s$, we get

$$
\begin{equation*}
\mathrm{K}_{m h s}^{i} v^{m}=\dot{\Omega}_{h s} v^{i} \tag{2.5}
\end{equation*}
$$

where

$$
\begin{equation*}
\tilde{\Omega}_{h s} \stackrel{\text { dep }}{=}\left(\varepsilon_{h \mid s}-\varepsilon_{s \mid h}\right) . \tag{2.6}
\end{equation*}
$$

In view of the equation ( 1.6 b ) contracting the definition (2.1) with respect to the indices $i$ and $k$, we obtain

$$
\begin{equation*}
\mathrm{K}_{h j \mid s m}=\alpha_{s m} \mathrm{~K}_{h j} \tag{2.7}
\end{equation*}
$$

Introducing (2.3) into the left hand side of the above equation, we have

$$
\begin{equation*}
\varepsilon_{h \mid s} \eta_{j \mid m}+\eta_{j} \varepsilon_{h \mid s m}+\varepsilon_{h \mid m} \eta_{j \mid s}+\varepsilon_{h} \eta_{j \mid s m}=\alpha_{s m} \varepsilon_{h} \eta_{j} \tag{2.8}
\end{equation*}
$$

Making a commutator on the indices $s$ and $m$ from the last formula, we get

$$
\begin{equation*}
-\eta_{j} \varepsilon_{r} \mathrm{~K}_{h s m}^{r}-\varepsilon_{h} \eta_{r} \mathrm{~K}_{j s m}^{r}=Q_{s m} \varepsilon_{h} \eta_{j} \tag{2.9}
\end{equation*}
$$

where

$$
\begin{equation*}
Q_{s m} \stackrel{\text { dep }}{=} \alpha_{s m}-\alpha_{m s} . \tag{2.10}
\end{equation*}
$$

Multiplying (2.9) by $v^{h}$ and summing over $h$, we obtain

$$
\begin{equation*}
\varepsilon\left(\widetilde{\Omega}_{s m} \eta_{j}+\eta_{r} \mathrm{~K}_{j s m}^{r}+\eta_{j} \mathrm{Q}_{s m}\right)=0 \tag{2.1I}
\end{equation*}
$$

where we put

$$
\begin{equation*}
\varepsilon \stackrel{\text { def }}{=} \varepsilon_{h} v^{h} . \tag{2.12}
\end{equation*}
$$

In this way, we have to discuss here the next two cases:

$$
\begin{equation*}
\text { a) } \stackrel{\rightharpoonup}{\Omega}_{s m} \eta_{j}+\eta_{r} \mathrm{~K}_{j s m}^{r}+\eta_{j} Q_{s m}=0 \quad \text { and } \quad \text { b) } \varepsilon=0 . \tag{2.13}
\end{equation*}
$$

$$
\text { 2. The case } \eta=0
$$

By putting

$$
\begin{equation*}
\eta=\eta_{h} v^{h} \tag{3.I}
\end{equation*}
$$

condition (2.13 a) yields this case. We shall show this fact. In wiew of the above equation transvecting (2.13 a) by $v^{j}$ and summing over $j$, we get

$$
\begin{equation*}
\eta\left(2 \widetilde{\Omega}_{s m}+\mathrm{Q}_{s m}\right)=\mathrm{o} \tag{3.2}
\end{equation*}
$$

Consequently for the present case (2.13 a) we have to consider two cases:

$$
\begin{equation*}
\text { a) } \eta=0 \quad \text { and } \quad \text { b) } \quad 2 \check{\Omega}_{s m}+Q_{s m}=0 \tag{3.3}
\end{equation*}
$$

Transvecting the latter case by $v^{i}$ and making use of (2.5), we obtain

$$
\begin{equation*}
Q_{s m} v^{i}+2 \mathrm{~K}_{r s m}^{i} v^{r}=0 . \tag{3.4}
\end{equation*}
$$

By virtue of the equation ( I .6 b ) contracting the last relation with respect to the indices $i$ and $m$, we have

$$
\begin{equation*}
Q_{s m} v^{m}+2 \mathrm{~K}_{r s} v^{r}=0 . \tag{3.5}
\end{equation*}
$$

Introducing the equations (2.3) and (2.12) into the above relation we get

$$
\begin{equation*}
Q_{s m} v^{m}+2 \varepsilon \eta_{s}=0 . \tag{3.6}
\end{equation*}
$$

Multiplying this equation by $v^{s}$ and noting (2.3), (2.12) and (3.1) we obtain

$$
\begin{equation*}
Q_{s m} v^{s} v^{m}+2 \varepsilon \eta=0 \tag{3.7}
\end{equation*}
$$

or

$$
\begin{equation*}
\varepsilon \eta=0 . \tag{3.8}
\end{equation*}
$$

In this way, we have $\eta=0$, i.e. the second case ( 3.3 b ) means first case (3.3 a) and ( 2.13 a) may be replaced by $\eta=0$. By this reason there exists only one case $\eta=0$.

By virtue of the commutation formula (I.3), commutating the indices $s$ and $m$. of the definition (2.1), we get

$$
\begin{align*}
& \mathrm{Q}_{s m} \mathrm{~K}_{h j k}^{i}=-\dot{\partial}_{r} \mathrm{~K}_{h j k}^{i} \mathrm{~K}_{s m}^{r}+\mathrm{K}_{h j k}^{r} \mathrm{~K}_{r s m}^{i}-\mathrm{K}_{r j k}^{i} \mathrm{~K}_{h s m}^{r}-  \tag{3.9}\\
&-\mathrm{K}_{h r k}^{i} \mathrm{~K}_{j s m}^{r}-\mathrm{K}_{h j r}^{i} \mathrm{~K}_{k s m}^{r},
\end{align*}
$$

where we have used the equation (2.10) also.
Contracting the last relation with respect to the indices $i$ and $k$ we have

$$
\begin{equation*}
Q_{s m} \mathrm{~K}_{h j}=-\mathrm{K}_{h r} \mathrm{~K}_{i s m}^{r}-\mathrm{K}_{r j} \mathrm{~K}_{h s m}^{r} \tag{3.10}
\end{equation*}
$$

where we have made use of (r.6).
In view of the equations (2.3), (2.5) and (2.12) transvecting the last equation by $v^{h}$ and summing over $h$, we get

$$
\begin{equation*}
\eta_{r} \mathrm{~K}_{j s m}^{r}=-\eta_{j}\left(\mathrm{Q}_{s m}+\widetilde{\Omega}_{s m}\right) \tag{3.1I}
\end{equation*}
$$

where we have used the fact that $\varepsilon \neq 0$.
Introducing (3.1 I) into the left hand side of the equation (2.9), we obtain

$$
\begin{equation*}
\eta_{j}\left(\varepsilon_{r} \mathrm{~K}_{h s m}^{r}-\varepsilon_{h} \tilde{\Omega}_{s m}\right)=0 . \tag{3.12}
\end{equation*}
$$

In this way, there occur the following two cases to be discussed:
a) $\eta_{j}=0 \quad$ and $\left.\quad b\right) \quad \varepsilon_{r} \mathrm{~K}_{h s m}^{r}=\varepsilon_{h} \stackrel{\rightharpoonup}{\Omega}_{s m}$.

Introducing the Bianchi identity ( I .5 ) for the curvature tensor $\mathrm{K}_{h s m}^{r}(x, \dot{x})$ in the case of ( 3.13 b ), we have

$$
\begin{equation*}
\varepsilon_{h} \widetilde{\Omega}_{s m}+\varepsilon_{s} \stackrel{\rightharpoonup}{\Omega}_{m h}+\varepsilon_{m} \vec{\Omega}_{h s}=\mathrm{o} \tag{3.14}
\end{equation*}
$$

Transvecting the last identity by $v^{h}$ and summing over $h$, we get

$$
\begin{equation*}
\stackrel{\rightharpoonup}{\Omega}_{s m}=\eta_{s} \varepsilon_{m}-\gamma_{m} \varepsilon_{s} \tag{3.15}
\end{equation*}
$$

where we have used $\widetilde{\Omega}_{h m}=-\widetilde{\Omega}_{n h}$ and $\stackrel{\Omega}{\Omega}_{m h} v^{h}=-\mathrm{K}_{m h} v^{h}$.
In view of definitions (2.3) and (2.6), we can write down an interesting formula:

$$
\begin{equation*}
\mathrm{K}_{s m}-\mathrm{K}_{m s}=\varepsilon_{m \mid s}-\varepsilon_{s \mid m} \tag{3.16}
\end{equation*}
$$

By virtue of the equations (2.5) and (3.13 b), we have

$$
\begin{equation*}
\varepsilon_{r} \mathrm{~K}_{h m s}^{r} v^{i}=\varepsilon_{h} \mathrm{~K}_{r s m}^{i} v^{r} \tag{3.17}
\end{equation*}
$$

and this means

$$
\begin{equation*}
\mathrm{K}_{h s m}^{r} v_{\mid r}^{i}-\mathrm{K}_{r s m}^{i} v_{\mid h}^{r}=0 \tag{3.18}
\end{equation*}
$$

That is to say, we actually have

$$
\begin{equation*}
\left(v_{\mid h}^{i}\right)_{\mid s m}-\left(v_{\mid h}^{i}\right)_{\mid m s}=0 . \tag{3.19}
\end{equation*}
$$

Consequently we can imagine the existence of a gradient vector and we are able to put

$$
\begin{equation*}
v_{\mid s m}^{i}=\sigma_{m} v_{\left.\right|_{s}}^{i} \tag{3.20}
\end{equation*}
$$

In view of equations (2.2) and (3.20), we can conclude that

$$
\begin{equation*}
\varepsilon_{h \mid s}+\varepsilon_{h} \varepsilon_{s}=\sigma_{s} \varepsilon_{h} . \tag{3.2I}
\end{equation*}
$$

Transvecting the above equation by $v^{h}$ and using the relation (2.12), we have

$$
\begin{align*}
\varepsilon \sigma_{s} & =\varepsilon_{h \mid s} v^{h}+\varepsilon \varepsilon_{s}=\left(\varepsilon_{h} v^{h}\right)_{\mid s}-\varepsilon_{h} v_{\mid s}^{h}+\varepsilon \varepsilon_{s}  \tag{3.22}\\
& =\varepsilon_{\mid s}-\varepsilon_{h} \varepsilon_{s} v^{s}+\varepsilon \varepsilon_{s}=\varepsilon_{\mid s}-\varepsilon \varepsilon_{s}+\varepsilon \varepsilon_{s} \\
& =\varepsilon_{\mid s} .
\end{align*}
$$

In this way, the existence of $\sigma_{s}$ is examined and we have here a characteristic condition on $\nu_{\mid h}^{i}$ :

$$
\begin{equation*}
v_{\mid h s}^{i}=\sigma_{s} v_{\mid h}^{i} \quad, \quad \sigma_{s}=\sigma_{\mid s} / \sigma \tag{3.23}
\end{equation*}
$$

On the other hand, in case of (3.13 a), from (2.3), we have

$$
\begin{equation*}
\mathrm{K}_{h j}=\mathrm{o} . \tag{3.24}
\end{equation*}
$$

Summarizing all the considerations above, we can state the following:
Theorem (3.1). In an n-dimensional second order recurrent Finsler space admitting a contravariant vector $v^{i}$ characterized by (2.2) and having a disjoint Ricci tensor of the form (2.3) there exists a case in which $\eta_{m} v^{m}=0$. In this case, if $\eta_{m}=\mathrm{o}$, then we have the vanishing of Ricci tensor $\mathrm{K}_{h j}$ and if $\eta_{j} \neq 0$, we have (3.16). The mixed tensor $v_{\mid h}^{i}$ itself is a recurrent one characterized by (3.23) for a definite gradient vector $\sigma_{s}=\sigma_{\mid s} / \sigma$.

$$
\text { 4. The case of } \varepsilon=0
$$

Let us consider - on the case of (2.13 b). Then using an analogous method as in $\eta 3$ transvecting (3.10) by $v^{j}$, we have

$$
\begin{equation*}
\varepsilon_{r} \mathrm{~K}_{h s m}^{r}=-\varepsilon_{h}\left(Q_{s m}+\widetilde{\Omega}_{s m}\right) \tag{4.I}
\end{equation*}
$$

Substituting (4.1) into the left hand side of (2.9), we get

$$
\begin{equation*}
\varepsilon_{h}\left(\eta_{\boldsymbol{r}} \mathrm{K}_{j s m}^{r}-\eta_{j} \tilde{\Omega}_{s m}\right)=\mathrm{o} . \tag{4.2}
\end{equation*}
$$

So, we have here two cases to be discussed. They are

$$
\begin{equation*}
\text { a) } \varepsilon_{h}=0 \quad \text { and } \quad \text { b) } \eta_{r} \mathrm{~K}_{j s m}^{r}=\eta_{j} \widetilde{\Omega}_{s m} \tag{4.3}
\end{equation*}
$$

Case (4.3a) yields the one of $v_{\mid h}^{i}=0$ and $Q_{h j}=0$. Case (4.3b) may be treated as follows. We can find

$$
\begin{equation*}
\eta_{h} \stackrel{\rightharpoonup}{\Omega}_{s m}+\eta_{s} \vec{\Omega}_{m h}+\eta_{m} \widetilde{\Omega}_{h s}=0 \tag{4.4}
\end{equation*}
$$

We also have

$$
\begin{equation*}
\stackrel{\Omega}{\Omega}_{s m} v^{s}=-\mathrm{K}_{s m} v^{s}=-\varepsilon_{s} \eta_{m} v^{s}=-\varepsilon \eta_{m}=0 . \tag{4.5}
\end{equation*}
$$

In view of the above equation transvecting (4.4) by $v^{h}$ we get

$$
\begin{equation*}
\eta \check{\Omega}_{m s}=\mathrm{o}, \tag{4.6}
\end{equation*}
$$

where we have also used (3.1). From (2.5) by contraction on the indices $i$ and $h$, we obtain

$$
\begin{equation*}
\tilde{\Omega}_{m s}=0 \tag{4.7}
\end{equation*}
$$

or

$$
\begin{equation*}
\varepsilon_{m \mid s}=\varepsilon_{s \mid m} \tag{4.8}
\end{equation*}
$$

Thus, we can state here the next
ThEOREM (4.1). When $\varepsilon=0$ in our space, there exists two cases. One of them is a case in which $v_{1 h}^{i}=0$ is satisfied and $\mathrm{K}_{h j}=\mathrm{o}$ and the other is the case of (4.8).

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