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**Curvature invariants associated with fundamental
forms of a Finsler hypersurface**

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Geometria differenziale. — *Curvature invariants associated with fundamental forms of a Finsler hypersurface.* Nota di U. P. SINGH, presentata (*) dal Socio E. BOMPIANI.

RIASSUNTO. — H Rund ha definito [1] n forme fondamentali per una ipersuperficie in uno spazio di Riemann n -dimensionale.

Analoga ricerca è qui fatta per le ipersuperficie di uno spazio di Finsler n -dimensionale nel caso siano soddisfatte ulteriori condizioni semplificatrici.

I. FUNDAMENTAL FORMULAE

Let a hypersurface F_{n-1} given by the equations

$$x^i = x^i(u^\alpha); \quad i = 1, \dots, n; \quad \alpha = 1, \dots, n-1$$

be immersed in a n -dimensional Finsler space F_n . The components $\dot{x}^i$ and $\dot{u}^\alpha$ of a vector [of magnitude $F(x, \dot{x})$] tangent to F_{n-1} are related by $\dot{x}^i = B_\alpha^i \dot{u}^\alpha$, where $B_\alpha^i = \frac{\partial x^i}{\partial u^\alpha}$. The unit vector l^i along $\dot{x}^i$ (or $\dot{x}^i$ itself) is taken to be the element of support. The metric tensors $g_{ij}(x, \dot{x})$ and $g_{\alpha\beta}(u, \dot{u})$ of F_n and F_{n-1} are related by

$$(1.1) \quad g_{\alpha\beta}(u, \dot{u}) = g_{ij}(x, \dot{x}) B_\alpha^i B_\beta^j.$$

The normal vector $N^i(x, \dot{x})$ of the hypersurface satisfies the conditions

$$(1.2) \quad g_{ij}(x, \dot{x}) N^j B_\alpha^i = N_i B_\alpha^i = 0 \quad \text{and} \quad g_{ij}(x, \dot{x}) N^i N^j = 1.$$

For a displacement vector dx^i and a vector-field X^i , both tangent to F_{n-1} , we may write

$$(1.3) \quad dx^i = B_\alpha^i du^\alpha, \quad X^i = B_\alpha^i X^\alpha.$$

If DX^i stands for the Cartan's covariant differential, then the 'induced' differential is defined by

$$(1.4) \quad \bar{D}X^\alpha = B_i^\alpha DX^i$$

where

$$B_i^\alpha = g^{\alpha\epsilon}(u, \dot{u}) g_{ij}(x, \dot{x}) B_\epsilon^j.$$

The connection parameters Γ_{hk}^{*i} , $\Gamma_{\alpha\beta}^{*\epsilon}$ and tensor $I_{\alpha\beta}^i$ (defined in [2] pages 68, 160 and 193 respectively) may be used in defining

$$(1.5) \quad \Omega_{\alpha\beta} = N_i I_{\alpha\beta}^i.$$

The tensor $\Omega_{\alpha\beta}$ is called the second fundamental tensor of the hypersurface.

(*) Nella seduta del 15 gennaio 1972.

We quote the following results (Rund [2] pages 158-165) for reference in the later sections of this paper.

Defining

$$(1.6) \quad N_j^i = N^i N_j \quad , \quad A_{kh}^i = F C_{kh}^i \quad , \quad C_{\alpha\beta\gamma} = \mathcal{G}_{\beta\varepsilon} C_{\alpha\gamma}^\varepsilon = C_{ijk} B_\alpha^i B_\beta^j B_\gamma^k$$

and

$$(1.7) \quad \bar{\Gamma}_{\beta\gamma}^{*\alpha} = \Gamma_{\beta\gamma}^\alpha - C_{\beta\delta}^\alpha \Gamma_{\varepsilon\gamma}^\delta u^\varepsilon$$

it has been shown

$$(1.8) \quad D\ell^i - B_\alpha^i \bar{D}\ell^\alpha = H_\gamma^i du^\gamma$$

and

$$(1.9) \quad DX^i - B_\alpha^i \bar{D}X^\alpha = H_{\gamma\beta}^i X^\beta du^\gamma + N_j^i A_{hk}^j B_\beta^h B_\gamma^k X^\beta \bar{D}\ell^\gamma$$

where we have taken

$$(1.10) \quad FH_\gamma^i = (B_{\beta\gamma}^i - B_\varepsilon^i \Gamma_{\beta\gamma}^\varepsilon) u^\beta + B_\gamma^k \Gamma_{hk}^i x^h$$

and

$$(1.11) \quad H_{\gamma\beta}^i = (B_{\beta\gamma}^i - B_\alpha^i \bar{\Gamma}_{\beta\gamma}^{*\alpha} + \Gamma_{hk}^{*i} B_\beta^h B_\gamma^k + A_{hk}^i B_\beta^h H_\gamma^k)$$

(the symbols C_{ijk} , $\Gamma_{\beta\gamma}^\alpha$ have their usual meanings [2]).

Rund [3] has introduced certain tensors which vanish identically in any locally Euclidean (or Riemannian) theory of hypersurface. These tensors and associated quantities are defined by

$$(1.12) \quad M_{ij}(x, \dot{x}) = C_{ijk}(x, \dot{x}) N^k \quad , \quad M_i(x, \dot{x}) = C_{ijk} N^j N^k$$

and

$$(1.13) \quad M_{\alpha\beta} = C_{ijk}(x, \dot{x}) B_\alpha^i B_\beta^j N^k \quad , \quad M_\alpha = C_{ijk}(x, \dot{x}) B_\alpha^i N^j N^k.$$

The following relations are obvious.

$$(1.14) \quad M_{ij}(x, \dot{x}) \dot{x}^j = 0 \quad , \quad M_i(x, \dot{x}) \dot{x}^i = 0 , \\ M_{\alpha\beta}(u, \dot{u}) \dot{u}^\beta = 0 \quad \text{and} \quad M_\alpha(u, \dot{u}) \dot{u}^\alpha = 0 .$$

2. n FUNDAMENTAL FORMS OF F_{n-1}

Let $C: u^\alpha = u^\alpha(\lambda)$ be a curve (not in the direction of u^α) of the hypersurface. The components in x^i and u^α of the tangent vector of this curve are such that

$$(2.1) \quad \frac{dx^i}{d\lambda} = B_\alpha^i \frac{du^\alpha}{d\lambda} .$$

This vector is normalised by the condition

$$(2.2) \quad g_{ij}(x, \dot{x}) \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda} = g_{\alpha\beta}(u, \dot{u}) \frac{du^\alpha}{d\lambda} \frac{du^\beta}{d\lambda} = 1.$$

Considering the displacement du^γ in (1.9) along C and putting $X^\alpha = \frac{du^\alpha}{d\lambda}$ we obtain

$$(2.3) \quad \frac{D\xi^i}{D\lambda} - B_\alpha^i \frac{\bar{D}\xi^\alpha}{D\lambda} = H_{\gamma\beta}^i(u, \dot{u}) \xi^\beta \xi^\gamma + FN^i M_{\beta\gamma}(u, \dot{u}) \xi^\beta \frac{\bar{D}\lambda^\gamma}{D\lambda},$$

where for brevity we have taken $\xi^i = \frac{dx^i}{d\lambda}$, $\xi^\alpha = \frac{du^\alpha}{d\lambda}$. It may be verified with the help of equation (1.4) that $\left(\frac{D\xi^i}{D\lambda} - B_\alpha^i \frac{\bar{D}\xi^\alpha}{D\lambda}\right)$ is along the normal to the hypersurface and this may be used in obtaining

$$(2.4) \quad K\left(u, \dot{u}, \frac{du}{d\lambda}\right) = N_i \left(\frac{D\xi^i}{D\lambda} - B_\alpha^i \frac{\bar{D}\xi^\alpha}{D\lambda}\right) = \\ = H_{\gamma\beta}^i(u, \dot{u}) \xi^\beta \xi^\gamma + FM_{\beta\gamma}(u, \dot{u}) \xi^\beta \frac{\bar{D}\lambda^\gamma}{D\lambda},$$

where in view of (1.11), (1.10) and relations $\Gamma_{\beta\gamma}^\varepsilon \dot{u}^\beta = \Gamma_{\beta\gamma}^{*\varepsilon} \dot{u}^\beta$; $\Gamma_{hk}^i \dot{x}^h = \Gamma_{hk}^{*i} \dot{x}^h$,

$$(2.5) \quad H_{\beta\gamma} = N_i H_{\beta\gamma}^i = \Omega_{\beta\gamma} + M_\gamma \Omega_{\beta\delta} \dot{u}^\delta.$$

An expression for induced mixed covariant derivative $N_{||\gamma}^i$ has been obtained by Rund [3]. This expression when substituted in $DN^i = N_{||\gamma}^i du^\gamma$ will reduce it to

$$(2.6) \quad DN^i = -\Omega_{\sigma\gamma} g^{\sigma\varepsilon} B_\varepsilon^i du^\gamma + (N^i M_i - 2 M_i^i) I_{\sigma\gamma}^i \dot{u}^\sigma du^\gamma.$$

In view of this relation and equations (1.5), (1.13) we obtain

$$(2.7) \quad -g_{ij} DN^i dx^j = [\Omega_{\gamma\delta} + 2 M_\delta \Omega_{\sigma\gamma} \dot{u}^\sigma] du^\gamma du^\delta.$$

Using the equation (2.6) once again and simplifying with the help of (1.2), (1.12) and (1.13) we obtain

$$(2.8) \quad g_{ij} DN^i DN^j = [\Omega_{\theta\gamma} \Omega_{\rho\delta} g^{\rho\theta} + 4 M_\varepsilon \Omega_{\sigma\gamma} g^{\sigma\varepsilon} \Omega_{\rho\delta} \dot{u}^\rho + \\ + (4 M_i M^i - 3 M^2) \Omega_{\sigma\gamma} \Omega_{\rho\delta} \dot{u}^\sigma \dot{u}^\rho] du^\gamma du^\delta$$

where

$$M = M_i N^i.$$

The expressions (2.7) and (2.8) are too complicated to make further study possible. We, therefore, assume that at x^i the element of support $\dot{x}^i$ satisfies the condition

$$(2.9) \quad M_{ij}(x, \dot{x}) = C_{ijk}(x, \dot{x}) N^k(x, \dot{x}) = 0.$$

This implies $M_i = 0$, $M_{\alpha\beta} = 0$, $M_\alpha = 0$. The equations (2.7) and (2.8) will now reduce to

$$(2.10) \quad -g_{ij}(x, \dot{x}) DN^i dx^j = C_{(2)\alpha\beta}(u, \dot{u}) du^\alpha du^\beta$$

and

$$(2.11) \quad g_{ij}(x, \dot{x}) DN^i DN^j = C_{(3)\alpha\beta}(u, \dot{u}) du^\alpha du^\beta$$

where

$$(2.12) \quad C_{(2)\alpha\beta} = \Omega_{\alpha\beta} \quad \text{and} \quad C_{(3)\alpha\beta} = C_{(2)\alpha\varepsilon} \Omega_\beta^\varepsilon.$$

Using the above equation and following the method of Rund [1] we write down the recurrence formula

$$(2.13) \quad C_{(p)\alpha\beta}(u, \dot{u}) = C_{(p-1)\alpha\varepsilon}(u, \dot{u}) \Omega_\beta^\varepsilon(u, \dot{u})$$

$C_{(p)\alpha\beta}(u, \dot{u})$ will be called the coefficients of the p -th fundamental form of F_{n-1} .

3. PRINCIPAL DIRECTIONS

The conditions $M_{ij}(x, \dot{x}) = 0$ and equations (2.4), (2.5) yield

$$(3.1) \quad K(u, \dot{u}, \xi) = \Omega_{\beta\gamma}(u, \dot{u}) \xi^\beta \xi^\gamma.$$

Since $K(u, \dot{u}, \xi)$ is a quadratic expression in ξ , the equations determining principal directions will be linear. They will therefore yield $(n-1)$ principal directions of the hypersurface. It is now possible to extend the theory developed in [1] to the hypersurface of a Finsler space.

In particular, if $\lambda_1, \lambda_2, \dots, \lambda_{n-1}$ are principal curvatures of the hypersurface then we define p -th curvature H_p of F_{n-1} in the form

$$(3.2) \quad H_p = \Sigma \lambda_1 \lambda_2 \cdots \lambda_p = \text{Sum of products (taken } p \text{ at a time) of the principal curvatures} \\ (p = 1, \dots, n-1).$$

The p -th associated mean curvature of F_{n-1} is defined as

$$(3.3) \quad M_p = g^{\alpha\beta} C_{(p+1)\alpha\beta}.$$

Using the method of [1] we deduce

$$(3.4) \quad H_1 = g^{\alpha\beta} \Omega_{\alpha\beta} = M_1,$$

$$(3.5) \quad H_2 = \sum_{r=1}^{n-1} \sum_{s=1}^{n-1} (\Omega_r^r \Omega_{r+s}^{r+s} - \Omega_{r+s}^r \Omega_r^{r+s})$$

(no summation with respect to r and s) and

$$(3.6) \quad 2 H_2 = M_1^2 - M_2.$$

4. INDUCED AND INTRINSIC CURVATURE TENSORS

In view of the results obtained above, it is possible to express the curvature invariant H_2 defined by (3.2) in terms of the curvature tensors of F_n and F_{n-1} . In the existing literature two types of curvature tensors have been defined for a Finsler hypersurface (Rund [3]). These are called induced and intrinsic curvature tensors. The generalised equations of Gauss have been obtained for both the curvature tensors. The generalised Gauss equation for the induced curvature tensor $K_{\alpha\beta\gamma\varepsilon}$ is given by (Rund [4])

$$(4.1) \quad K_{\alpha\beta\gamma\varepsilon} = K_{ijhk} B_\alpha^i B_\beta^j B_\gamma^h B_\varepsilon^k + (\Omega_{\alpha\gamma} \Omega_{\beta\varepsilon} - \Omega_{\alpha\varepsilon} \Omega_{\beta\gamma}) + \\ + 2 C_{jhl} B_\beta^j N^h (I_{\sigma\varepsilon}^l \Omega_{\alpha\gamma} - I_{\sigma\gamma}^l \Omega_{\alpha\varepsilon}) u^\sigma + \\ + g_{ij} B_\beta^j B_\alpha^h (\partial \Gamma_{hk}^{*i} / \partial x^l) (I_{\sigma\varepsilon}^l B_\gamma^k - I_{\sigma\gamma}^l B_\varepsilon^k) u^\sigma,$$

where K_{ijhk} is the curvature tensor of F_n .

The intrinsic curvature tensor $\hat{K}_{\alpha\beta\gamma\varepsilon}$ has been obtained with the help of intrinsic connection parameter $\hat{\Gamma}_{\beta\gamma}^\alpha$ which is defined with respect to the metric of F_{n-1} in a manner formally identical with the mode of definition of Γ_{hk}^{*i} in terms of the metric of F_n . Writing

$$2 G^\varepsilon = \gamma_{\alpha\beta}^\varepsilon u^\alpha u^\beta, \quad G_v^\varepsilon = \frac{\partial G^\varepsilon}{\partial u^v},$$

we have the definition (Rund [3])

$$\hat{K}_{\alpha\gamma\varepsilon}^\delta = \frac{\partial \hat{\Gamma}_{\alpha\gamma}^\delta}{\partial u^\varepsilon} - \frac{\partial \hat{\Gamma}_{\alpha\varepsilon}^\delta}{\partial u^\gamma} - \left(\frac{\partial \hat{\Gamma}_{\alpha\gamma}^\delta}{\partial u^\lambda} G_\varepsilon^\lambda - \frac{\partial \hat{\Gamma}_{\alpha\varepsilon}^\delta}{\partial u^\lambda} G_\gamma^\lambda \right) + \hat{\Gamma}_{\lambda\varepsilon}^\delta \hat{\Gamma}_{\alpha\gamma}^\lambda - \hat{\Gamma}_{\lambda\gamma}^\delta \hat{\Gamma}_{\alpha\varepsilon}^\lambda,$$

($\gamma_{\alpha\beta}^\varepsilon$ being the Christoffel's symbols of the second kind for F_{n-1}). The Gauss equation for intrinsic curvature tensor is now given by ([3])

$$(4.2) \quad \hat{K}_{\alpha\beta\gamma\varepsilon} = K_{ijhk} B_\alpha^i B_\beta^j B_\gamma^h B_\varepsilon^k + (\Omega_{\alpha\gamma} \Omega_{\beta\varepsilon} - \Omega_{\alpha\varepsilon} \Omega_{\beta\gamma}) + \\ + 2 M_{jl} B_\beta^j (\Omega_{\alpha\gamma}^l J_{\sigma\varepsilon}^l - \Omega_{\alpha\varepsilon}^l J_{\sigma\gamma}^l) u^\sigma + (\Delta_{\sigma\beta\gamma} \Delta_{\alpha\varepsilon}^\sigma - \Delta_{\sigma\varepsilon} \Delta_{\alpha\gamma}^\sigma) + \\ + (\Delta_{\alpha\beta\gamma/\varepsilon} - \Delta_{\alpha\varepsilon/\gamma}) + g_{ij} B_\beta^j \frac{\partial \Gamma_{hk}^{*i}}{\partial x^l} B_\alpha^h (J_{\sigma\varepsilon}^l B_\gamma^k - J_{\sigma\gamma}^l B_\varepsilon^k) u^\sigma,$$

where

$$(4.3) \quad \Delta_{\alpha\beta}^\varepsilon = \hat{\Gamma}_{\alpha\beta}^\varepsilon - \Gamma_{\alpha\beta}^{*\varepsilon}, \quad \Delta_{\alpha\gamma\beta} = g_{\gamma\varepsilon} \Delta_{\alpha\beta}^\varepsilon,$$

$$(4.4) \quad J_{\alpha\beta}^i = B_{\alpha\beta}^i - B_\varepsilon^i \hat{\Gamma}_{\alpha\beta}^\varepsilon + \Gamma_{hk}^{*i} B_\alpha^h B_\beta^k = I_{\alpha\beta}^i - B_\varepsilon^i \Delta_{\alpha\beta}^\varepsilon,$$

and $\Delta_{\alpha\gamma\beta/\varepsilon}$ is the intrinsic covariant derivative of $\Delta_{\alpha\beta\gamma}$. In view of the fact ([3])

$$(4.5) \quad \Delta_{\alpha\gamma\beta} = C_{hjk} N^j [(B_\beta^h B_\gamma^k \Omega_{\alpha\sigma} + B_\alpha^h B_\gamma^k \Omega_{\beta\sigma} - B_\beta^h B_\alpha^k \Omega_{\gamma\sigma}) u^\sigma - \\ - (C_{\beta\gamma}^\lambda B_\alpha^h B_\lambda^k + C_{\alpha\gamma}^\lambda B_\beta^h B_\lambda^k - C_{\beta\alpha}^\lambda B_\gamma^h B_\lambda^k) \Omega_{\sigma\lambda} u^\sigma u^\lambda]$$

and the condition (2.9) we deduce

$$\Delta_{\alpha\gamma\beta} = 0, \quad \hat{\Gamma}_{\alpha\beta}^{\varepsilon} = \Gamma_{\alpha\beta}^{*\varepsilon} \quad \text{and} \quad J_{\alpha\beta}^i = I_{\alpha\beta}^i.$$

A comparison of equations (4.1) and (4.2) now reveals that if the element of support satisfies the condition (2.9) then each induced and intrinsic generalised Gauss equation will reduce to

$$(4.6) \quad K_{\alpha\beta\gamma\varepsilon} = \hat{K}_{\alpha\beta\gamma\varepsilon} = K_{ijhk} B_{\alpha}^i B_{\beta}^j B_{\gamma}^h B_{\varepsilon}^k + (\Omega_{\alpha\gamma} \Omega_{\beta\varepsilon} - \Omega_{\alpha\varepsilon} \Omega_{\beta\gamma}) + \\ + g_{ij} B_{\beta}^j B_{\alpha}^h \frac{\partial \Gamma_{hk}^{*i}}{\partial x^l} N^l (\Omega_{\sigma\varepsilon} B_{\gamma}^k - \Omega_{\sigma\gamma} B_{\varepsilon}^k) u^{\sigma}.$$

This shows that (2.9) is a sufficient condition for the equality of induced and intrinsic curvature tensors.

The equation (4.6) will be used in finding a relation between the curvature invariant H_2 defined by (3.2) and the curvature tensors of F_n and F_{n-1} . In order to obtain this relation we multiply (4.6) by $g^{\beta\varepsilon} g^{\alpha\gamma}$. The left hand side reduces to

$$(4.7) \quad g^{\beta\varepsilon} g^{\alpha\gamma} K_{\alpha\beta\gamma\varepsilon} = g^{\alpha\gamma} K_{\alpha\gamma} = \bar{K},$$

where $\bar{K}$ and $K_{\alpha\gamma}$ are respectively the scalar curvature and Ricci tensor (both intrinsic and induced) of the hypersurface. For simplifying the right hand side we shall use the relations

$$(4.8) \quad g^{\beta\varepsilon} B_{\beta}^j B_{\varepsilon}^k = g^{jk} - N^j N^k,$$

and

$$(4.9) \quad K_{jihk} = -K_{ijhk} - 2 C_{ijl} K_{\gamma hk}^l x^{\gamma}$$

(Rund [2] page 105).

The above two equations the condition (2.9) and the skew-symmetry of K_{jihk} in h and k yield

$$(4.10) \quad (K_{jihk} B_{\alpha}^i B_{\beta}^j B_{\gamma}^h B_{\varepsilon}^k) g^{\alpha\gamma} g^{\beta\varepsilon} = (K_{ih} - K_{jihk} N^j N^k) (g^{ih} - N^i N^h) = \\ = K - 2 K_{ih} N^i N^h,$$

where K and K_{ih} stand for the scalar curvature and Ricci tensor of F_n . Further, in view of equations (2.12), (3.3) and (3.6) we have

$$(4.11) \quad g^{\beta\varepsilon} g^{\alpha\gamma} (\Omega_{\alpha\gamma} \Omega_{\beta\varepsilon} - \Omega_{\alpha\varepsilon} \Omega_{\beta\gamma}) = M_1^2 - M_2 = 2 H_2.$$

The last term in (4.6) on being multiplied by $g^{\beta\varepsilon} g^{\alpha\gamma}$ gives

$$(4.12) \quad g_{ij} B_{\beta}^j B_{\alpha}^h \frac{\partial \Gamma_{hk}^{*i}}{\partial x^l} N^l (\Omega_{\sigma\varepsilon} B_{\gamma}^k - \Omega_{\sigma\gamma} B_{\varepsilon}^k) g^{\beta\varepsilon} g^{\alpha\gamma} u^{\sigma} = \\ = \left(\frac{\partial \Gamma_{jkh}^{*i}}{\partial x^l} g^{jk} - \frac{\partial \Gamma_{jkh}^{*i}}{\partial x^l} N^j N^k - \frac{\partial \Gamma_{hi}^{*i}}{\partial x^l} + \frac{\partial \Gamma_{hjk}^{*i}}{\partial x^l} N^j N^k \right) N^l \Omega_{\sigma}^{\alpha} B_{\alpha}^h u^{\sigma},$$

where we have used the condition (2.9), the equation (4.8) and the relations

$$B_i^\varepsilon = g_{ij} g^{\beta\varepsilon} B_\beta^j, \quad B_i^\varepsilon B_\varepsilon^k = (\delta_i^k - N^k N_i)$$

in simplification. After differentiating $g^{jk} g_{ij} = \delta_i^k$ and using (2.9) we deduce

$$\frac{\partial g^{jk}}{\partial x^l} N^l = 0.$$

This relation helps in getting

$$(4.13) \quad \frac{\partial \Gamma_{hi}^{*i}}{\partial x^l} N^l = g^{jk} \frac{\partial \Gamma_{hjk}^{*}}{\partial x^l} N^l.$$

Combining (4.7), (4.10), (4.11), (4.12) and (4.13) we obtain

$$(4.14) \quad \bar{K} = K - 2 K_{ij} N^i N^j + 2 H_2 + 2 \frac{\partial \Gamma_{[jk]k}^{*}}{\partial x^l} (g^{jk} - N^j N^k) N^l \Omega_\sigma^\alpha B_\alpha^h u^\sigma$$

where

$$\frac{\partial \Gamma_{[jk]k}^{*}}{\partial x^l} = \frac{1}{2} \left(\frac{\partial \Gamma_{jkk}^{*}}{\partial x^l} - \frac{\partial \Gamma_{hjk}^{*}}{\partial x^l} \right).$$

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