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# Some Simple Nonlinear PDE's Without Solutions. 

Haïm Brezis - Xavier Cabré


#### Abstract

Sunto. - In questo articolo consideriamo alcune semplici equazioni a derivate parziali elittiche nonlineari, per le quali il Teorema della Funzione Inversa, se applicato in modo formale, suggerisce l'esistenza di soluzioni. Nonostante ciò, proviamo che non esistono soluzioni neppure in vari sensi deboli. Un problema modello è dato da $-\Delta u=\left(u^{2} /|x|^{2}\right)+\operatorname{cin} \Omega, u=0$ su $\partial \Omega$, dove $\Omega \subset \mathbb{R}^{N}, N \geqslant 2$, è un dominio limitato contenente 0. Per qualunque costante $c>0$, arbitrariamente piccola, proviamo che questo problema non ammette soluzioni distribuzionali in $\mathscr{O}^{\prime}(\Omega \backslash\{0\})$. Mostriamo anche come, cercando di approssimare il problema mediante un certo procedimento naturale, accada che le soluzioni dei problemi approssimati esplodano dappertutto in $\Omega$. Infine proviamo gli analoghi parabolici di questi risultati e, in particolare, alcuni fenomeni di «blow-up» istantaneo e completo..


## 0. - Introduction.

The original motivation of this work is the following. Consider the simple problem

$$
\begin{cases}-\Delta u=a(x) u^{2}+f(x) & \text { in } \Omega  \tag{0.1}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

where $\Omega$ is a smooth bounded domain in $\mathbb{R}^{N}, N \geqslant 3$. If $a(x) \in L^{p}(\Omega)$ and $p>$ $N / 2$, then for any $f \in L^{p}(\Omega)$ with $\|f\|_{p}$ small, problem (0.1) has a unique small solution $u$ in $W^{2, p}(\Omega)$. This is an easy consequence of the Inverse Function Theorem applied to $F(u)=-\Delta u-a(x) u^{2}$ which maps $X=W^{2, p}(\Omega) \cap$ $W_{0}^{1, p}(\Omega)$ into $Y=L^{p}(\Omega)$ (recall that $W^{2, p}(\Omega) \subset L^{\infty}(\Omega)$ by the Sobolev imbedding theorem), since its differential at $0, D F(0)=-\Delta$ is bijective.

As a special case suppose $a(x)=|x|^{-\alpha}$ in a domain $\Omega$ containing 0 , with $0<\alpha<2$. Then for any small constant $c$ the problem

$$
\begin{cases}-\Delta u=\frac{u^{2}}{|x|^{\alpha}}+c & \text { in } \Omega  \tag{0.2}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

has a unique small solution.

The case $\alpha=2$ is interesting since $a(x)=|x|^{-2}$ does not belong to $L^{p}(\Omega)$ for $p>N / 2$. One may then wonder what happens to the problem

$$
\begin{cases}-\Delta u=\frac{u^{2}}{|x|^{2}}+c & \text { in } \Omega,  \tag{0.3}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

On the one hand, a formal computation suggests that since the linearized operator at 0 is $-\Delta$, which is bijective, problem ( 0.3 ) has a solution for small $c$. On the other hand, the $F$ above does not map $X=W^{2, p}(\Omega) \cap W_{0}^{1, p}(\Omega)$ into $Y=L^{p}(\Omega)$ for any $1<p<\infty$. One may then try to construct other function spaces, for example weighted spaces, where the Inverse Function Theorem might apply. This is doomed to fail. In fact, our main results show that for any constant $c>0$ (no matter how small) problem (0.3) has no solution, even in a very weak sense. When $c<0$, problem (0.3) does, however, have a solution (see Remark 1.4).

In Sections 1 and 4 we propose several notions of weak solutions and establish nonexistence. A basic ingredient in Section 1 is the following:

Theorem 0.1. - Assume $0 \in \Omega$. If $u \in L_{\text {loc }}^{1}(\Omega), u \geqslant 0$ a.e. with $u^{2} /|x|^{2} \in$ $L_{\text {loc }}^{1}(\Omega)$ is such that

$$
\begin{equation*}
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{D}^{\prime}(\Omega) \tag{0.4}
\end{equation*}
$$

then $u \equiv 0$.

The proof of Theorem 0.1 uses an adaptation of a method introduced in [4]. In Section 4 we prove a stronger result, namely:

Theorem 0.2. - Assume $0 \in \Omega$. If $u \in L_{\text {loc }}^{2}(\Omega \backslash\{0\}), u \geqslant 0$ a.e. is such that

$$
\begin{equation*}
-|x|^{2} \Delta u \geqslant u^{2} \quad \text { in } \mathscr{O}^{\prime}(\Omega \backslash\{0\}) \tag{0.5}
\end{equation*}
$$

then $u \equiv 0$.

Theorem 0.2 is proved using appropriate powers of testing functions-an idea due to Baras-Pierre [2]. As a consequence we obtain the nonexistence of local solutions (i.e., in any neighborhood of 0 , without prescribing any boundary condition) for a very simple nonlinear equation:

Theorem 0.3. - Assume $0 \in \Omega$ and $c>0$. There is no function $u \in$ $L_{\mathrm{loc}}^{2}(\Omega \backslash\{0\})$ satisfying

$$
\begin{equation*}
-|x|^{2} \Delta u=u^{2}+c \quad \text { in } \mathcal{O}^{\prime}(\Omega \backslash\{0\}) . \tag{0.6}
\end{equation*}
$$

In Section 3 we examine what happens to a natural approximation procedure of (0.3). Consider for example the equation

$$
\begin{cases}-\Delta u=\frac{\min \left\{u^{2}, n\right\}}{|x|^{2}+(1 / n)}+c & \text { in } \Omega, c>0  \tag{0.7}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

For any $n$ there is a minimal solution $u_{n}$. We prove that $u_{n}(x) \rightarrow+\infty$ for every $x \in \Omega$ as $n \rightarrow \infty$, i.e., there is complete blow-up in the sense of Baras-Cohen [1]. Again, this rules out any reasonable notion of weak solution for (0.3).

In Section 2 we extend the previous results to more general problems such as

$$
-\Delta u=a(x) g(u)+b(x)
$$

assuming only that $g \geqslant 0$ on $\mathbb{R}, g$ is continuous and nondecreasing on $[0, \infty)$ and

$$
\int^{\infty} \frac{d s}{g(s)}<\infty
$$

with $a \in L_{\text {loc }}^{1}(\Omega), a \geqslant 0$ and

$$
\int \frac{a(x)}{|x|^{N-2}}=\infty
$$

The original motivation of our research came from observations made in [6] and [4].

For any $N \geqslant 3$ the problem

$$
\begin{cases}-\Delta u=2(N-2) e^{u} & \text { in } B_{1}=\left\{x \in \mathbb{R}^{N} ;|x|<1\right\}, \\ u=0 & \text { on } \partial B_{1},\end{cases}
$$

admits the weak solution $\bar{u}(x)=\log \left(1 /|x|^{2}\right)$. It was observed in [6] that when
$N \geqslant 11$ the linearized operator at $\bar{u}$ namely

$$
\begin{aligned}
L v & =-\Delta v-2(N-2) e^{\bar{u}} v, \\
& =-\Delta v-\frac{2(N-2)}{|x|^{2}} v,
\end{aligned}
$$

is coercive and thus formally bijective; this is a simple consequence of Hardy's inequality:

$$
\int|\nabla v|^{2} \geqslant \frac{(N-2)^{2}}{4} \int \frac{v^{2}}{|x|^{2}}, \quad \forall v \in H_{0}^{1}\left(B_{1}\right)
$$

(note that $(N-2)^{2} / 4>2(N-2)$ when $N \geqslant 11$ ). On the other hand the results of [4] show that when $N \geqslant 10$ the perturbed problem

$$
\begin{cases}-\Delta u=2(N-2) e^{u}+c & \text { in } B_{1} \\ u=0 & \text { on } \partial B_{1}\end{cases}
$$

has no solution even in a weak sense and no matter how small $c$ is, provided $c>0$.

This strange «failure» of the Inverse Function Theorem is only apparent. As was pointed out in [6] this just means that there is no functional setting in which it can be correctly applied. We have tried here, in the spirit of Open Problem 6 in [6], to analyze the same phenomenon for simple examples in low dimensions.

After our investigation was completed we learned about an interesting work of N. J. Kalton and I. E. Verbitsky [8] (which was carried out independently of ours). Consider for example the problem

$$
\begin{cases}-\Delta u=a(x) u^{2}+c & \text { in } \Omega  \tag{0.8}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

with $a \in L_{\mathrm{loc}}^{1}(\Omega), a \geqslant 0$ and $c$ a positive constant.
Their result says that if (0.8) has a weak solution, then necessarily

$$
\begin{equation*}
G\left(a \delta^{2}\right) \leqslant C \delta \quad \text { in } \Omega \tag{0.9}
\end{equation*}
$$

for some constant $C$, where $G=(-\Delta)^{-1}$ (with zero boundary condition) and $\delta(x)=\operatorname{dist}(x, \partial \Omega)$. In particular, if $0 \in \Omega$ and $a(x)=1 /|x|^{2}$, then $G\left(a \delta^{2}\right) \simeq$ $|\log | x|\mid$ as $x \rightarrow 0$ and thus (0.9) fails; hence (0.8) has no weak solution.

We present in Section 5 a very simple proof of the main result of [8] using a variant of the method developed in Section 1.

Finally, in Section 6 we present a parabolic analogue of Theorem 0.2. It ex-
tends, in particular, a result of I. Peral and J. L. Vázquez [12]. Namely, the problem

$$
\begin{cases}u_{t}-\Delta u=2(N-2) e^{u} & \text { in } B_{1} \times(0, T), \\ u=0 & \text { on } \partial B_{1} \times(0, T), \\ u(x, 0)=u_{0}, & \end{cases}
$$

with $u_{0} \geqslant \bar{u}=\log \left(1 /|x|^{2}\right), u_{0} \neq \bar{u}$, has no solution $u \geqslant \bar{u}$ even for small time: instantaneous and complete blow-up occurs.

The plan of the paper is the following:
(1) Proof of Theorem 0.1.
(2) General nonlinearities.
(3) Complete blow-up.
(4) Very weak solutions. Proofs of Theorems 0.2 and 0.3 .
(5) Connection with a result of Kalton-Verbitsky.
(6) Evolution equations.

Notation. - Throughout this paper, $\Omega$ is a bounded smooth domain of $\mathbb{R}^{N}$, $N \geqslant 1$, such that $0 \in \Omega$. We write

$$
\delta(x)=\operatorname{dist}(x, \partial \Omega)
$$

and $L_{\delta}^{1}(\Omega)=L^{1}(\Omega, \delta(x) d x)$. We denote by $C_{0}^{\infty}(\Omega)$ the space of $C^{\infty}$ functions with compact support in $\Omega$, and by $\mathscr{\omega}^{\prime}(\Omega)$ the space of distributions in $\Omega$. By $C$ we denote a positive constant which may be different in each inequality.

## 1. - Proof of Theorem 0.1.

In this section we prove Theorem 0.1 and its consequences. We first introduce some terminology about weak solutions.

Definition 1.1. - Let $h(x, u)$ be a Carathéodory function in $\Omega \times \mathbb{R}$, that is, $h(x, u)$ is measurable in $x$ and continuous in $u$ for a.e. $x$.
(a) We say that

$$
-\Delta u=h(x, u) \quad \text { in } \mathscr{o}^{\prime}(\Omega)
$$

if $u \in L_{\text {loc }}^{1}(\Omega), h(x, u) \in L_{\text {loc }}^{1}(\Omega)$ and $-\int u \Delta \varphi=\int h(x, u) \varphi$ for any $\varphi \in$ $C_{0}^{\infty}(\Omega)$.
(b) We say that

$$
-\Delta u \geqslant h(x, u) \quad \text { in } \mathscr{O}^{\prime}(\Omega)
$$

if $u \in L_{\text {loc }}^{1}(\Omega), h(x, u) \in L_{\text {loc }}^{1}(\Omega)$ and $-\int u \Delta \varphi \geqslant \int h(x, u) \varphi$ for any $\varphi \in$ $C_{0}^{\infty}(\Omega)$ with $\varphi \geqslant 0$.
(c) We say that $u$ is a weak solution of

$$
\begin{cases}-\Delta u=h(x, u) & \text { in } \Omega, \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

if $u \in L^{1}(\Omega), h(x, u) \in L_{\delta}^{1}(\Omega)$ and $-\int_{\Omega} u \Delta \zeta=\int_{\Omega} h(x, u) \zeta$ for any $\zeta \in C^{2}(\bar{\Omega})$
with $\zeta=0$ on $\partial \Omega$.
The following is the main result of this section; it is Theorem 0.1 of the Introduction.

Theorem 1.2. - Let $N \geqslant 1$ and $u \in L_{\text {loc }}^{1}(\Omega)$ satisfy $u \geqslant 0$ a.e. in $\Omega, u^{2} /|x|^{2} \in$ $L_{\text {loc }}^{1}(\Omega)$ and

$$
\begin{equation*}
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } ब^{\prime}(\Omega) \tag{1.1}
\end{equation*}
$$

Then $u \equiv 0$.

This theorem easily implies two nonexistence results. The first one deals with the following boundary value problem.

Corollary 1.3. - Let $N \geqslant 1$ and $f \in L_{\delta}^{1}(\Omega)$ satisfy $f \geqslant 0$ a.e. and $f \not \equiv 0$. Then there is no weak solution of

$$
\begin{cases}-\Delta u=\frac{u^{2}}{|x|^{2}}+f(x) & \text { in } \Omega  \tag{1.2}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

in the sense of Definition 1.1.

Remark 1.4. - When $f \equiv 0$, problem (1.2) has $u \equiv 0$ as the only weak solution; this follows immediately from Theorem 1.2 . When $N \geqslant 3$ and $f \in L^{1}(\Omega)$ satisfies $f \leqslant 0$ and $f \not \equiv 0$, (1.2) has a weak solution $u \leqslant 0$, which is the unique solution among nonpositive functions. This fact is a consequence of a more general result of Gallouët and Morel [7], which extends work of Brezis and Strauss [5].

As a consequence of Theorem 1.2, we may write down a simple PDE without local solutions, i.e., no solution exists in any neighborhood of 0 . Here, we do not impose any boundary condition.

Corollary 1.5. - Let $N \geqslant 3$ and $c>0$ be any positive constant. Then there is no function $u$ such that $u^{2} /|x|^{2} \in L_{\mathrm{loc}}^{1}(\Omega)$ and

$$
\begin{equation*}
-\Delta u=\frac{u^{2}+c}{|x|^{2}} \quad \text { in } \mathscr{O}^{\prime}(\Omega) \tag{1.3}
\end{equation*}
$$

Remark 1.6. - In contrast with (1.3), the equation

$$
\begin{equation*}
-\Delta u=\frac{u^{2}}{|x|^{2}}+c \tag{1.4}
\end{equation*}
$$

has a weak solution in some neighborhood of 0 , if $N \geqslant 3$ and $c>0$ is a constant. This solution is nonpositive and can be obtained from the results of [7] as follows.

We introduce the new unknown $v=-u-c|x|^{2}$ so that (1.4) becomes

$$
\begin{equation*}
-\Delta v+\frac{v^{2}}{|x|^{2}}+2 c v=c(2 N-1)-c^{2}|x|^{2} \equiv h(x) \tag{1.5}
\end{equation*}
$$

We solve (1.5) on $B_{R}$ with the boundary condition $v=0$ on $\partial B_{R}$. Note that $h \geqslant 0$ on $B_{R}$ provided $R$ is sufficiently small $\left(R^{2} \leqslant(2 N-1) / c\right)$. The results of [7] give a unique solution $v \geqslant 0$.

In Section 4 we will prove stronger nonexistence results for problems (1.2) and (1.3).

The proof of Theorem 1.2 is based on the following variant of Kato's inequality [9].

LEmMA 1.7. - Let $u \in L_{\mathrm{loc}}^{1}(\Omega)$ and $f \in L_{\mathrm{loc}}^{1}(\Omega)$ satisfy

$$
-\Delta u \geqslant f \quad \text { in } \mathscr{D}^{\prime}(\Omega)
$$

Let $\phi: \mathbb{R} \rightarrow \mathbb{R}$ be a $C^{1}$, concave function such that

$$
0 \leqslant \phi^{\prime} \leqslant C \quad \text { in } \mathbb{R}
$$

for some constant $C$. Then $\phi(u) \in L_{\mathrm{loc}}^{1}(\Omega)$ and

$$
-\Delta \phi(u) \geqslant \phi^{\prime}(u) f \quad \text { in } \sigma^{\prime}(\Omega) .
$$

The proof is standard, smoothing $u, f$ and $\phi$ by convolution; see also Lemma 2 in [4].

The proof of Theorem 1.2 is a variant of a method introduced in [4]. Consid-
er the function $\phi(s)=(1 / \varepsilon)-(1 / s)$. It is nonnegative, bounded, increasing and concave in the interval $[\varepsilon, \infty), \varepsilon>0$. Note that if $u \geqslant \varepsilon$ satisfies

$$
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}}
$$

then

$$
-\Delta \phi(u) \geqslant \phi^{\prime}(u) \frac{u^{2}}{|x|^{2}}=\frac{1}{|x|^{2}} .
$$

This will lead to a contradiction with the fact that $\phi(u)$ is bounded. The details go as follows.

Proof of Theorem 1.2. - Suppose that $u$ is as in the theorem, and that $u \not \equiv 0$. Since $u \geqslant 0, u \not \equiv 0,-\Delta u \geqslant 0$ in $\mathscr{\sigma}^{\prime}(\Omega)$ and $\Omega$ is connected, we have that

$$
u \geqslant \varepsilon \quad \text { a.e. in } \mathrm{B}_{\eta},
$$

for some $\varepsilon>0$ and $B_{\eta}=B_{\eta}(0)$ with closure in $\Omega$. If $N \leqslant 2$, this is already a contradiction with $u^{2} /|x|^{2} \in L_{\text {loc }}^{1}(\Omega)$.

Suppose $N \geqslant 3$. Let

$$
\phi(s)=\frac{1}{\varepsilon}-\frac{1}{s} \quad \text { for } s \geqslant \varepsilon
$$

and extend it by $\phi(s)=(s-\varepsilon) / \varepsilon^{2}$ for $s \leqslant \varepsilon$. Note that $\phi: \mathbb{R} \rightarrow \mathbb{R}$ is $C^{1}$, concave and $0 \leqslant \phi^{\prime} \leqslant 1 / \varepsilon^{2}$, so that $\phi$ satisfies all the conditions of Lemma 1.7. Recall that $u \geqslant \varepsilon$ in $B_{\eta}$, and consider

$$
v=\phi(u)=\frac{1}{\varepsilon}-\frac{1}{u} \quad \text { in } B_{\eta} .
$$

It satisfies $0 \leqslant v \leqslant 1 / \varepsilon$ in $B_{\eta}$ and, by Lemma 1.7,

$$
-\Delta v \geqslant \phi^{\prime}(u) \frac{u^{2}}{|x|^{2}}=\frac{1}{|x|^{2}} \quad \text { in } \mathscr{O}^{\prime}\left(B_{\eta}\right)
$$

Hence $v-(1 /(N-2)) \log 1 /|x|) \in L^{1}\left(B_{\eta}\right)$ and

$$
-\Delta\left(v-\frac{1}{N-2} \log \frac{1}{|x|}\right) \geqslant 0 \quad \text { in } \mathscr{D}^{\prime}\left(B_{\eta}\right)
$$

which implies

$$
v-\frac{1}{N-2} \log \frac{1}{|x|} \geqslant-C \quad \text { in } B_{\eta / 2}
$$

for some constant $C>0$. In particular, $v(x) \rightarrow+\infty$ as $x \rightarrow 0$, which is a contradiction with the fact that $v \leqslant 1 / \varepsilon$.

Remark 1.8. - In the previous proof we could have used (in the spirit of [4]) the function

$$
w=\frac{u}{\varepsilon u+1}
$$

for any $\varepsilon>0$, instead of $v=1 / \varepsilon-1 / u$. Note that

$$
\psi(s)=\frac{s}{\varepsilon s+1}
$$

satisfies $\psi^{\prime}(s)=(\varepsilon s+1)^{-2}$ and hence

$$
\psi^{\prime}(s) s^{2}=\psi(s)^{2}
$$

Moreover, $\psi$ is bounded in [ $0, \infty$ ) and satisfies all the conditions (for $\phi$ ) of Lemma 1.7 in $[0, \infty)$. In particular, if $u \geqslant 0, u \not \equiv 0$ and $-\Delta u \geqslant u^{2} /|x|^{2}$ then

$$
-\Delta w \geqslant \psi^{\prime}(u) \frac{u^{2}}{|x|^{2}}=\frac{w^{2}}{|x|^{2}}
$$

and we can conclude as before, since $w^{2} /|x|^{2} \geqslant v /|x|^{2}$ in a subdomain of $\Omega$, for some constant $v>0$.

Analogous versions of the functions $\phi$ and $\psi$ will appear, in Sections 2 and 5 , when the nonlinearity $u^{2}$ is replaced by more general nonlinearities $g(u)$.

Finally, we use Theorem 1.2 to prove the nonexistence results of this section.

Proof of Corollary 1.3. - Suppose that $u$ is a weak solution of (1.2). Since
$\int_{\Omega} u(-\Delta \zeta) \geqslant 0$ for any $\zeta \in C^{2}(\bar{\Omega})$ with $\zeta \geqslant 0$ in $\Omega$ and $\zeta=0$ on $\partial \Omega$, we easily deduce $u \geqslant 0$ in $\Omega$. Moreover, $-\Delta u \geqslant u^{2} /|x|^{2}$ in $\mathscr{D}^{\prime}(\Omega)$. We obtain, by Theorem $1.2, u \equiv 0$. This is a contradiction with (1.2), since $f \not \equiv 0$.

PRoof of Corollary 1.5. - Suppose that $u$ is a solution of (1.3) in $\mathscr{\sigma}^{\prime}(\Omega)$. Then $-\Delta u \geqslant c /|x|^{2}$ in $\mathscr{D}^{\prime}\left(B_{\eta}\right)$, for some ball $B_{\eta}=B_{\eta}(0)$ with closure in $\Omega$. As
in the proof of Theorem 1.2, we deduce that

$$
u-\frac{c}{N-2} \log \frac{1}{|x|} \geqslant-C \quad \text { in } B_{\eta / 2}
$$

for some constant $C$. In particular, $u \geqslant 0$ in $B_{v}$ for some small $v>0$.
We therefore have

$$
\begin{gathered}
u \geqslant 0 \quad \text { in } B_{v} \\
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{D}^{\prime}\left(B_{v}\right) .
\end{gathered}
$$

By Theorem 1.2, $u \equiv 0$ in $B_{v}$, which is a contradiction with equation (1.3).

## 2. - General nonlinearities.

In this section we extend the previous nonexistence results to more general problems of the form

$$
-\Delta u=a(x) g(u)+b(x)
$$

We assume (here and throughout the rest of the paper) that $g: \mathbb{R} \rightarrow[0, \infty$ ) is continuous on $\mathbb{R}$, nondecreasing on $[0, \infty), g(s)>0$ if $s>0$, and

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d s}{g(s)}<\infty \tag{2.1}
\end{equation*}
$$

Power functions $g(u)=u^{p}$ for $u \geqslant 0$, with $p>1$ are examples of such nonlinearities. We suppose that $N \geqslant 3$.

For the function $a(x)$ we assume in this section that $a \in L_{\mathrm{loc}}^{1}(\Omega), a \geqslant 0$ in $\Omega$, and

$$
\begin{equation*}
\int_{B_{\eta}(0)} \frac{a(x)}{|x|^{N-2}}=\infty \tag{2.2}
\end{equation*}
$$

for some $\eta>0$ small enough (or, equivalently, for any $\eta>0$ small).
Theorem 2.1. - Let $g$ and a satisfy the assumptions above, with $N \geqslant 3$.
(a) Let $u \geqslant 0$ a.e. in $\Omega$ satisfy

$$
-\Delta u \geqslant a(x) g(u) \quad \text { in } \mathscr{\sigma}^{\prime}(\Omega)
$$

Then $u \equiv 0$.
(b) Let $f \in L_{\delta}^{1}(\Omega)$ satisfy $f \geqslant 0$ a.e. and $f \not \equiv 0$. Then there is no weak solution of

$$
\begin{cases}-\Delta u=a(x) g(u)+f(x) & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

(c) If $b(x)$ satisfies the same conditions as $a(x)$, then there is no weak solution of

$$
-\Delta u=a(x) g(u)+b(x) \quad \text { in } \mathscr{D}^{\prime}(\Omega)
$$

This theorem is proved with the same method as in the previous section. We only need to adapt two points. First, $\phi(s)=(1 / \varepsilon)-(1 / s)$ has to be replaced by a solution of

$$
\phi^{\prime}(s)=\frac{1}{g(s)} \quad \text { if } s \geqslant \varepsilon
$$

where $\varepsilon>0$ is a constant.
We therefore define

$$
\phi(s)=\int_{\varepsilon}^{s} \frac{d t}{g(t)} \quad \text { if } s \geqslant \varepsilon
$$

which satisfies $0 \leqslant \phi \leqslant \int_{\varepsilon}^{\infty} d t / g(t)<\infty$ in [ $\varepsilon, \infty$ ) (by assumption (2.1)), $\phi(\varepsilon)=$ $0, \phi^{\prime}(\varepsilon)=1 / g(\varepsilon), \phi$ is $C^{1}$, concave (since $\phi^{\prime}(s)=1 / g(s)$ is nonincreasing) and $0 \leqslant \phi^{\prime} \leqslant 1 / g(\varepsilon)$ in $[\varepsilon, \infty)$. Extending $\phi$ by $\phi(s)=(s-\varepsilon) / g(\varepsilon)$ for $s \leqslant \varepsilon$, we obtain a function $\phi$ on all of $\mathbb{R}$, satisfying the conditions of Lemma 1.7 and with $\phi$ bounded from above.

To complete the proof of Theorem 2.1, we only need to consider a solution $w \in L_{\text {loc }}^{1}\left(B_{\eta}\right)$, where $B_{\eta}=B_{\eta}(0)$ with closure in $\Omega$, of

$$
-\Delta w=a(x) \quad \text { in } \mathscr{O}^{\prime}\left(B_{\eta}\right)
$$

(a solution always exists since $a \in L^{1}\left(B_{\eta}\right)$ ) and show that

$$
\begin{equation*}
\underset{B_{1 / n}(0)}{\operatorname{ess} \inf } w \rightarrow \infty+\infty \quad \text { as } n \rightarrow+\infty \tag{2.3}
\end{equation*}
$$

For this purpose, we consider the convolution in $B_{\eta}, \tilde{w}=a * C_{N}|x|^{2-N}$, where $C_{N}$ is chosen such that $-\Delta\left(C_{N}|x|^{2-N}\right)=\delta_{0}$. Then $w-\tilde{w}$ is harmonic in $B_{\eta}$ and hence bounded in $B_{\eta / 2}$. In particular, it suffices to show (2.3) for $\widetilde{w}$. But
this is true since, for $|x| \leqslant 1 / n$,

$$
\tilde{w}(x)=C_{N} \int_{B_{\eta}} \frac{a(y)}{|y-x|^{N-2}} d y \geqslant C \int_{B_{\eta}} \frac{a(y)}{|y|^{N-2}+(1 / n)^{N-2}} d y \rightarrow+\infty
$$

as $n \rightarrow \infty$, by (2.2).

## 3. - Complete blow-up.

In Corollary 1.3 and Theorem 2.1 we have proved the nonexistence of weak solutions of some problems of the form

$$
\begin{cases}-\Delta u=a(x) g(u)+f(x) & \text { in } \Omega  \tag{3.1}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

In this section we prove that, under the same assumptions on $a(x), g(u)$ and $f(x)$ made in Section 2, approximate solutions of (3.1) blow up everywhere in $\Omega$, that is, there is complete blow-up. More precisely, we have the following.

Let $g(u)$ and $a(x)$ be as in Section 2. Suppose that $f \in L_{\delta}^{1}(\Omega), f \geqslant 0$ a.e. and $f \equiv 0$. Let ( $g_{n}$ ) be a sequence of nonnegative, bounded, nondecreasing and continuous functions in $[0, \infty)$ such that $\left(g_{n}\right)$ increases pointwise to $g$. Let $a_{n}$ and $f_{n}$ be two sequences of nonnegative bounded functions in $\Omega$, increasing pointwise to $a$ and $f$, respectively.

Theorem 3.1. - Under the above assumptions, let $\underline{u}_{n}$ be the minimal nonnegative solution of the approximate problem

$$
\begin{cases}-\Delta u=a_{n}(x) g_{n}(u)+f_{n}(x) & \text { in } \Omega,  \tag{3.2}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

Then, as $n \rightarrow+\infty$,

$$
\frac{\underline{u}_{n}(x)}{\delta(x)} \rightarrow+\infty \quad \text { uniformly in } \Omega
$$

In the proof of Theorem 3.1 we use two ingredients. First, the nonexistence result of Theorem $2.1(b)$ (see also Corollary 1.3) and, second, the following estimate for the linear Laplace equation. It asserts that, for some positive constant $c$,

$$
\begin{equation*}
G(x, y) \geqslant c \delta(x) \delta(y) \quad \text { in } \bar{\Omega} \times \bar{\Omega}, \tag{3.3}
\end{equation*}
$$

where $G$ is the Green's function of the Laplacian in $\Omega$ with zero Dirichlet con-
dition. In an equivalent way, we can state this lower bound on $G$ as follows.

Lemma 3.2. - Suppose that $h \geqslant 0$ belongs to $L^{\infty}(\Omega)$. Let $v$ be the solution of

$$
\begin{cases}-\Delta v=h & \text { in } \Omega, \\ v=0 & \text { on } \partial \Omega .\end{cases}
$$

Then

$$
\begin{equation*}
\frac{v(x)}{\delta(x)} \geqslant c \int_{\Omega} h \delta, \quad \forall x \in \Omega, \tag{3.4}
\end{equation*}
$$

where $c>0$ is a constant depending only on $\Omega$.
Estimate (3.3) was proved by Morel and Oswald [11] (in unpublished work), and by Zhao [13] (in a stronger form). For the convenience of the reader we give a simple proof of (3.3).

Proof of Lemma 3.2. - We proceed in two steps.
Step 1. For any compact set $K \subset \Omega$, we show

$$
\begin{equation*}
v(x) \geqslant c \int_{\Omega} h \delta, \quad \forall x \in K \tag{3.5}
\end{equation*}
$$

where $c$ is a positive constant depending only on $K$ and $\Omega$. To prove (3.5), let $\varrho=\operatorname{dist}(K, \partial \Omega) / 2$, and take $m$ balls of radius $\varrho$ such that

$$
K \subset B_{\varrho}\left(x_{1}\right) \cup \ldots \cup B_{\varrho}\left(x_{m}\right) \subset \Omega .
$$

Let $\zeta_{1}, \ldots, \zeta_{m}$ be the solutions of

$$
\begin{cases}-\Delta \zeta_{i}=\chi_{B_{e}\left(x_{i}\right)} & \text { in } \Omega, \\ \zeta_{i}=0 & \text { on } \partial \Omega,\end{cases}
$$

where $\chi_{A}$ denotes the characteristic function of $A$. The Hopf boundary lemma implies that there is a constant $c>0$ such that

$$
\zeta_{i}(x) \geqslant c \delta(x), \quad \forall x \in \Omega, \quad \forall 1 \leqslant i \leqslant m .
$$

Here and in the rest of the proof, $c$ denotes various constants depending only on $K$ and $\Omega$. Let now $x \in K$, and take a ball $B_{\varrho}\left(x_{i}\right)$ containing $x$. Then $B_{\varrho}\left(x_{i}\right) \subset$
$B_{2 \varrho}(x) \subset \Omega$, and since $-\Delta v \geqslant 0$ in $\Omega$, we conclude

$$
v(x) \geqslant \int_{B_{2_{Q}}(x)} v=c \int_{B_{2_{e}}(x)} v \geqslant c \int_{B_{e}\left(x_{i}\right)} v=c \int_{\Omega} v\left(-\Delta \zeta_{i}\right)=c \int_{\Omega} h \zeta_{i} \geqslant c \int_{\Omega} h \delta .
$$

Step 2. Fix a smooth compact set $K \subset \Omega$. By (3.5), $v \geqslant c \int_{\Omega} h \delta$ in $K$, so that it
suffices to prove (3.4) for $x \in \Omega \backslash K$.
Let $w$ be the solution of

$$
\begin{cases}-\Delta w=0 & \text { in } \Omega \backslash K \\ w=0 & \text { on } \partial \Omega \\ w=1 & \text { on } \partial K\end{cases}
$$

The Hopf boundary lemma gives again

$$
w(x) \geqslant c \delta(x), \quad \forall x \in \Omega \backslash K
$$

Since $v$ is superharmonic and $v \geqslant c \int_{\Omega} h \delta$ on $\partial K$, the maximum principle
plies implies

$$
v(x) \geqslant c\left(\int_{\Omega} h \delta\right) w(x) \geqslant c\left(\int_{\Omega} h \delta\right) \delta(x), \quad \forall x \in \Omega \backslash K
$$

This completes the proof.
Proof of Theorem 3.1. - Consider the approximate problem

$$
\begin{cases}-\Delta u=a_{n}(x) g_{n}(u)+f_{n}(x) & \text { in } \Omega  \tag{3.6}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

Since $0 \leqslant a_{n}(x) g_{n}(s)+f_{n}(x) \leqslant C_{n}$ in $\Omega \times[0, \infty)$ for some constant $C_{n}$, we have that $C_{n} z$ is a supersolution of (3.6) $)_{n}$ where

$$
\begin{cases}-\Delta z=1 & \text { in } \Omega  \tag{3.7}\\ z=0 & \text { on } \partial \Omega\end{cases}
$$

On the other hand, 0 is a subsolution of $(3.6)_{n}$. We therefore obtain a minimal solution $\underline{u}_{n}$ of (3.6) ${ }_{n}$ by monotone iteration:

$$
\begin{cases}-\Delta u_{m+1}=a_{n}(x) g_{n}\left(u_{m}\right)+f_{n}(x) & \text { in } \Omega \\ u_{m+1}=0 & \text { on } \partial \Omega\end{cases}
$$

starting with $u_{0} \equiv 0$. In particular, since $a_{n}(x) g_{n}(s)+f_{n}(x)$ increases with $n$,
$\underline{u}_{n+1}$ is a supersolution of $(3.6)_{n}$, and hence

$$
\underline{u}_{n} \leqslant \underline{u}_{n+1} .
$$

We claim that

$$
\int_{\Omega} a_{n}(x) g_{n}\left(\underline{u}_{n}\right) \delta \nearrow+\infty \quad \text { as } n \nearrow+\infty .
$$

Lemma 3.2 then gives

$$
\frac{\underline{u}_{n}(x)}{\delta(x)} \nearrow+\infty \quad \text { uniformly in } \Omega
$$

as $n \rightarrow+\infty$; this proves Theorem 3.1. Thus, it only remains to show the claim. Suppose not, that

$$
\begin{equation*}
\int_{\Omega} a_{n}(x) g_{n}\left(\underline{u}_{n}\right) \delta \leqslant C, \quad \forall n \tag{3.8}
\end{equation*}
$$

Then, multiplying (3.6) ${ }_{n}$ by the solution $z$ of (3.7), we see that

$$
\int_{\Omega} \underline{u}_{n} \leqslant C, \quad \forall n
$$

(we have used that $0 \leqslant f_{n} \leqslant f \in L_{\delta}^{1}(\Omega)$ ). Hence, $\underline{u}_{n} \nearrow u$ in $L^{1}(\Omega)$, for some $u$, by monotone convergence.

Since $g_{n}$ is a nondecreasing function, $a_{n} g_{n}\left(\underline{u}_{n}\right)+f_{n}$ increases to $a g(u)+f$ a.e. in $\Omega$; (3.8) also gives

$$
a_{n} g_{n}\left(\underline{u}_{n}\right)+f_{n} \nearrow a g(u)+f \quad \text { in } L_{\delta}^{1}(\Omega),
$$

again by monotone convergence. We can now pass to the limit in the weak formulation of $(3.6)_{n}$ (recall Definition 1.1(c)), and obtain that $u$ is a weak solution of

$$
\begin{cases}-\Delta u=a(x) g(u)+f(x) & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

This is impossible by Theorem 2.1(b).

## 4. - Very weak solutions. Proofs of Theorems 0.2 and 0.3 .

In this section we return to the study of equation

$$
-\Delta u=\frac{u^{2}}{|x|^{2}}+f(x)
$$

and its corresponding boundary value problem. We prove stronger versions of the nonexistence results of Section 1 by considering a more general notion of solutions, that we call very weak solutions. More precisely, we prove the following results.

Theorem 4.1. - Let $N \geqslant 2$ and $u \in L_{\mathrm{loc}}^{2}(\Omega \backslash\{0\})$ satisfy $u \geqslant 0$ a.e. in $\Omega$ and

$$
\begin{equation*}
-|x|^{2} \Delta u \geqslant u^{2} \quad \text { in } \mathscr{O}^{\prime}(\Omega \backslash\{0\}), \tag{4.1}
\end{equation*}
$$

in the sense that $-\int u \Delta\left(|x|^{2} \varphi\right) \geqslant \int u^{2} \varphi$ for any $\varphi \geqslant 0, \quad \varphi \in$ $C_{0}{ }^{\infty}(\Omega \backslash\{0\})$.
Then $u \equiv 0$.

Note that now we are testing (4.1) only against functions with compact support in $\Omega$ and vanishing in a neighborhood of 0 . As a consequence of the previous theorem, we will prove the following stronger version of Corollary 1.3.

Corollary 4.2. - Let $N \geqslant 2$ and $f \in L_{\mathrm{loc}}^{1}(\Omega \backslash\{0\}), f \delta$ integrable near $\partial \Omega$, $f \geqslant 0$ a.e. in $\Omega, f \not \equiv 0$. Then there is no very weak solution of

$$
\begin{cases}-\Delta u=\frac{u^{2}}{|x|^{2}}+f(x) & \text { in } \Omega  \tag{4.2}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

in the sense that $u \in L_{\text {loc }}^{2}(\Omega \backslash\{0\}), u$ and $u^{2} \delta$ are integrable near $\partial \Omega$ and

$$
-\int_{\Omega} u \Delta\left(|x|^{2} \zeta\right)=\int_{\Omega} u^{2} \zeta+\int_{\Omega} f|x|^{2} \zeta
$$

for any $\zeta \in C^{2}(\bar{\Omega}), \zeta=0$ on $\partial \Omega$ and $\zeta \equiv 0$ in a neighborhood of 0.

Remark 4.3. - Theorem 4.1 does not hold in dimension $N=1$; a direct computation shows that $u(x)=|x|^{\alpha}$, for any $0<\alpha<1$, satisfies (4.1) when $\Omega$ is a small interval containing 0 . Corollary 4.2 is also false in dimension $N=1$; in fact, for any $p>1$ and $f \in L^{p}(-1,1)$, small in $L^{p}$, there exists a very weak solution $u$ of (4.2) in $\Omega=(-1,1)$ which is continuous in $[-1,1]$ and satisfies
$u(0)=0$. This solution $u$ is defined in $(0,1)$ to be the solution of

$$
\left\{\begin{array}{l}
-u^{\prime \prime}=\frac{u^{2}}{|x|^{2}}+f(x) \quad \text { in }(0,1)  \tag{4.3}\\
u(0)=u(1)=0
\end{array}\right.
$$

and similarly in $(-1,0)$; we obtain in this way a very weak solution $u$ of (4.2) in $(-1,1)$.

Note that if $\|f\|_{L^{p}(0,1)}$ is small then there is a solution $u \in C^{2}(0,1) \cap$ $C^{1}([0,1])$ of (4.3). It can be obtained through the Inverse Function Theorem applied to the operator $-u^{\prime \prime}-u^{2} /|x|^{2}$, which maps $X=W^{2, p}(0,1) \cap$ $W_{0}^{1, p}(0,1)$ into $L^{p}(0,1)$; note that if $u \in X$ then $u /|x| \in L^{\infty}(0,1)$ and $u \in C^{1}([0,1])$.

We also extend the local nonexistence result of Corollary 1.5, now for all $N \geqslant 1$.

Corollary 4.4. - Let $N \geqslant 1$ and $c>0$ be any positive constant. Then there is no very weak solution of

$$
\begin{equation*}
-|x|^{2} \Delta u=u^{2}+c \quad \text { in } \mathscr{D}^{\prime}(\Omega \backslash\{0\}), \tag{4.4}
\end{equation*}
$$

in the sense that $u \in L_{\text {loc }}^{2}(\Omega \backslash\{0\})$, and $-\int u \Delta\left(|x|^{2} \varphi\right)=\int\left(u^{2}+c\right) \varphi$ for any function $\varphi \in C_{0}^{\infty}(\Omega \backslash\{0\})$.

The proofs of the results of this section consist of using appropriate powers of testing functions; this is an idea due to Baras and Pierre [2] and employed in [2] for the study of removable singularities of solutions of semilinear equations. In fact, the first step in our proof of Theorem 4.1 will be to use equation (4.1) to show that $u /|x| \in L_{\text {loc }}^{2}(\Omega)$; this can be interpreted as a «removable singularity» result. The second step of the proof is to show that $-\Delta u \geqslant u^{2} /|x|^{2}$ is satisfied in $\mathscr{\sigma}^{\prime}(\Omega)$. We may then conclude that $u \equiv 0$ with the help of Theorem 1.2. We present here an alternative proof of Theorem 1.2 based on multiplication by a sequence of appropriate testing functions.

Proof of Theorem 4.1.
Step 1. We prove that $u /|x| \in L_{\text {loc }}^{2}(\Omega)$. For this purpose, let $\zeta_{n} \in$ $C_{0}^{\infty}(\Omega \backslash\{0\})$ be such that $0 \leqslant \zeta_{n} \leqslant 1$,

$$
\zeta_{n}= \begin{cases}0 & \text { if }|x| \leqslant 1 / n \\ 1 & \text { if }|x| \geqslant 2 / n, \quad x \in \omega \\ \zeta & \text { if } x \in \Omega \backslash \omega\end{cases}
$$

where $\omega$ is open, $0 \in \omega, \bar{\omega} \subset \Omega$ and $\xi$ is a fixed «tail» for all $\zeta_{n}$. We take $\zeta_{n}$ such
that if $1 / n<|x|<2 / n$ then

$$
\begin{gathered}
\nabla \zeta_{n}^{4}=4 \zeta_{n}^{3} \nabla \zeta_{n} \\
\Delta \zeta_{n}^{4}=4 \zeta_{n}^{3} \Delta \zeta_{n}+12 \zeta_{n}^{2}\left|\nabla \zeta_{n}\right|^{2}
\end{gathered}
$$

and

$$
\left|\Delta \zeta_{n}^{4}\right| \leqslant C n^{2} \zeta_{n}^{2}
$$

for some constant $C$ independent of $n$. Multiplying (4.1) by $\zeta_{n}^{4} /|x|^{2}$ yields

$$
\int \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4} \leqslant-\int u \Delta \zeta_{n}^{4} \leqslant C\left(\int_{\{1 / n<|x|<2 / n\}} n^{2} u \zeta_{n}^{2}\right)+C
$$

Using the Cauchy-Schwarz inequality, we obtain

$$
\int \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4} \leqslant C n\left(\int_{\{1 / n<|x|<2 / n\}} \frac{u}{|x|} \zeta_{n}^{2}\right)+C \leqslant \frac{C}{n^{N / 2-1}}\left(\int \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4}\right)^{1 / 2}+C .
$$

Since $N \geqslant 2$ we deduce that

$$
\int \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4} \leqslant C
$$

and hence $u /|x| \in L_{\text {loc }}^{2}(\Omega)$.
For later purposes, let us retain two more consequences of the previous proof. First, we did not use $u \geqslant 0$. Second, if carried out for $N=1$, the proof gives

$$
\begin{equation*}
\int_{2 / n}^{3 / n} \frac{u^{2}}{|x|^{2}} \leqslant C n+C, \quad(N=1) \tag{4.5}
\end{equation*}
$$

Step 2. We show that

$$
\begin{equation*}
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{O}^{\prime}(\Omega) \tag{4.6}
\end{equation*}
$$

Indeed, let $\varphi \in C_{0}^{\infty}(\Omega), \varphi \geqslant 0$, and $\eta_{n}(x)=\eta_{1}(n|x|)$ be such that $0 \leqslant \eta_{1} \leqslant 1$ and

$$
\eta_{1}(x)= \begin{cases}0 & \text { if }|x| \leqslant 1 \\ 1 & \text { if }|x| \geqslant 2\end{cases}
$$

Multiplying (4.1) by $\left(\varphi \eta_{n}\right) /|x|^{2}$ yields

$$
\int \frac{u^{2}}{|x|^{2}} \varphi \eta_{n} \leqslant-\int u \Delta\left(\varphi \eta_{n}\right)
$$

If we show that, as $n \rightarrow \infty$,

$$
\begin{equation*}
\int u|\nabla \varphi|\left|\nabla \eta_{n}\right| \rightarrow 0 \tag{4.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\int u \varphi\left|\Delta \eta_{n}\right| \rightarrow 0 \tag{4.8}
\end{equation*}
$$

then we obtain $-\int u \Delta\left(\varphi \eta_{n}\right) \rightarrow-\int u \Delta \varphi$ and hence the statement of Step 2. To prove (4.7) and (4.8), we use that $u /|x| \in L_{\text {loc }}^{2}(\Omega)$-which we established in Step 1. Hence

$$
\int u|\nabla \varphi|\left|\nabla \eta_{n}\right| \leqslant C n \int_{\{1 / n<|x|<2 / n\}} u \leqslant C \int_{\{1 / n<|x|<2 / n\}} \frac{u}{|x|} \rightarrow 0
$$

and

$$
\begin{aligned}
\int u \varphi\left|\Delta \eta_{n}\right| \leqslant C n^{2} \int_{\{1 / n<|x|<2 / n\}} u \leqslant C n & \int_{\{1 / n<|x|<2 / n\}} \frac{u}{|x|} \leqslant \\
& \frac{C}{n^{N / 2-1}}\left(\int_{\{1 / n<|x|<2 / n\}} \frac{u^{2}}{|x|^{2}}\right)^{1 / 2} \rightarrow 0 .
\end{aligned}
$$

Note that, as in Step 1, we have not used $u \geqslant 0$.
Step 3. We show that $u \equiv 0$ (it is only here where we use $u \geqslant 0$ ). Let us suppose that $u \not \equiv 0$. Then, since $u \geqslant 0,-\Delta u \geqslant 0$ in $\mathscr{\sigma}^{\prime}(\Omega)$ and $\Omega$ is connected, we have that

$$
u \geqslant \varepsilon \quad \text { a.e. in } B_{\eta}
$$

for some $\varepsilon>0$ and $B_{\eta}=B_{\eta}(0)$ with closure in $\Omega$. When $N=2$ this is impossible since $u /|x| \geqslant \varepsilon /|x|$ near 0 and hence $u /|x| \notin L_{\text {loc }}^{2}(\Omega)$-a contradiction with Step 1.

When $N \geqslant 3$, we use that

$$
-\Delta u \geqslant \frac{\varepsilon^{2}}{|x|^{2}}=-\Delta\left(\frac{\varepsilon^{2}}{N-2} \log \frac{1}{|x|}\right) \quad \text { in } \mathscr{D}^{\prime}\left(B_{\eta}\right)
$$

and we conclude (as in the proof of Theorem 1.2) that

$$
\begin{equation*}
u \geqslant \frac{\varepsilon^{2}}{N-2}|\log | x| |-C \quad \text { near } 0 \tag{4.9}
\end{equation*}
$$

for some $C>0$.
Let us now choose a sequence of functions $\chi_{n}(x)=\chi_{1}(n|x|)$ such that $0 \leqslant \chi_{n} \leqslant 1$ and

$$
\chi_{n}(x)= \begin{cases}1 & \text { if }|x| \leqslant 1 / n \\ 0 & \text { if }|x| \geqslant 2 / n\end{cases}
$$

Mutiplying (4.6) by $\chi_{n}^{4}$ yields

$$
\int \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \leqslant \int u\left|\Delta \chi_{n}^{4}\right| \leqslant C n^{2} \int_{\{1 / n<|x|<2 / n\}} u \chi_{n}^{2} \leqslant
$$

$$
\operatorname{Cn} \int \frac{u}{|x|} \chi_{n}^{2} \leqslant \frac{C}{n^{N / 2-1}}\left(\int \frac{u^{2}}{|x|^{2}} \chi_{n}^{4}\right)^{1 / 2},
$$

and therefore

$$
\int \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \leqslant \frac{C}{n^{N-2}}
$$

But, using (4.9), we have

$$
\int \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \geqslant c \int_{\{1 /(2 n)<|x|<1 / n\}} n^{2}|\log n|^{2} \sim \frac{|\log n|^{2}}{n^{N-2}}
$$

which contradicts the previous statement.
Finally we give the proofs of Corollaries 4.2 and 4.4.
Proof of Corollary 4.2. - Recall that Steps 1 and 2 of the previous proof hold for any $u$ satisfying (4.1)-without the assumption $u \geqslant 0$. Therefore, since $f \geqslant 0$, Step 1 of the proof of Theorem 4.1 gives

$$
\frac{u}{|x|} \in L_{\mathrm{loc}}^{2}(\Omega) .
$$

Moreover, proceeding as in Step 2 of the same proof, we see that

$$
-\int_{\Omega} u \Delta \zeta=\int_{\Omega} \frac{u^{2}}{|x|^{2}} \zeta+\int_{\Omega} f \zeta
$$

for any $\zeta \in C^{2}(\bar{\Omega}), \zeta=0$ on $\partial \Omega$. In particular, $-\int u \Delta \zeta \geqslant 0$ if, in addition, $\zeta \geqslant 0$ in $\Omega$. We conclude that $u \geqslant 0$ in $\Omega$. Theorem 4.1 implies that $u \equiv 0$, which contradicts (4.2) and $f \not \equiv 0$.

Proof of Corollary 4.4. - Suppose that $u$ satisfies

$$
-|x|^{2} \Delta u=u^{2}+c \quad \text { in } \mathscr{O}^{\prime}(\Omega \backslash\{0\}) .
$$

The proof of Step 1 in Theorem 4.1 gives

$$
\frac{u^{2}}{|x|^{2}}+\frac{c}{|x|^{2}} \in L_{\text {loc }}^{1}(\Omega) \quad \text { when } N \geqslant 2 .
$$

This is impossible when $N=2$, since $1 /|x|^{2}$ is not integrable near 0 .
When $N \geqslant 3$ we get (as in the proof of Theorem 4.1, Step 2) that

$$
-\Delta u=\frac{u^{2}}{|x|^{2}}+\frac{c}{|x|^{2}} \quad \text { in } \mathscr{D}^{\prime}(\Omega) .
$$

We now proceed as in the proof of Corollary 1.5, i.e., we compare $u$ with $\log (1 /|x|)$. We obtain that $u \geqslant 0$ in a neighborhood of 0 and hence, by Theorem 4.1, $u \equiv 0$. This is a contradiction with equation (4.4).

We finally treat the case $N=1$. We would have $u \in L_{\text {loc }}^{2}(\Omega \backslash\{0\})$ and

$$
-x^{2} u^{\prime \prime}=u^{2}+c \quad \text { in } \mathscr{O}^{\prime}(\Omega \backslash\{0\}) .
$$

In particular, $u$ belongs to $C^{2}(0, a)$ for some $a>0$. Integrating the inequality $-u^{\prime \prime} \geqslant c / x^{2}$ in ( $s, a / 2$ ), we obtain

$$
u^{\prime}(s) \geqslant \frac{c}{s}-C
$$

for some constant $C$. Integrating again yields

$$
-u(s) \geqslant-c \log s-C .
$$

Thus

$$
|u(s)| \geqslant c|\log s|-C \quad \text { near } 0 .
$$

On the other hand, we recall (4.5):

$$
\int_{2 / n}^{3 / n} \frac{u^{2}}{s^{2}} d s \leqslant C n+C
$$

which was proved in Step 1 of the proof of Theorem 4.1. Using $|u(s)| \geqslant$ $c|\log s|-C$, this inequality yields

$$
\frac{c}{6}|\log n|^{2} n=c|\log n|^{2} \int_{2 / n}^{3 / n} \frac{d s}{s^{2}} \leqslant \int_{2 / n}^{3 / n} \frac{u^{2}}{s^{2}} d s \leqslant C n+C,
$$

which is a contradiction.

## 5. - Connection with a result of Kalton-Verbitsky.

In this section we consider the problem (for $u \geqslant 0$ )

$$
\begin{cases}-\Delta u=a(x) u^{p}+f(x) & \text { in } \Omega,  \tag{5.1}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

where $p>1, a \geqslant 0$ and $f \geqslant 0$ in $\Omega$. Recently, Kalton and Verbitsky [8] have found an interesting necessary condition for the existence of a weak solution of (5.1). Their result states that if (5.1) has a weak solution, then necessarily

$$
\begin{equation*}
G\left(a G(f)^{p}\right) \leqslant C G(f) \quad \text { in } \Omega \tag{5.2}
\end{equation*}
$$

for some constant $C$, where $G=(-\Delta)^{-1}$ (with zero Dirichlet boundary condition). In [8] the authors also prove (5.2) for more general second-order elliptic operators.

In this section we give a simple proof of the necessary condition (5.2) (for the Laplacian) using a refinement of the method that we have developed in Section 1. Our proof gives (5.2) with constant $C=1 /(p-1)$. Next, we replace $f(x)$ by $\lambda f(x)$ in (5.1) (where $\lambda>0$ is a parameter) and we study the problem of existence of solution depending on the value of $\lambda$.

As pointed out in the Introduction, condition (5.2) easily implies some of our nonexistence results. For instance, it gives the result of Theorem 2.1(b) when $g(u)=u^{p}$, for some $p>1$, and $f \in L^{\infty}$, since in this case $G(f) \sim \delta$.

We recall that there is another necessary condition-due to Baras and Pierre [3]-for the existence of a weak solution of (5.1). Its proof consists of multiplying (5.1) by test functions and using Young's inequality.

Throughout this section we assume that

$$
\begin{equation*}
a \in L_{\mathrm{loc}}^{1}(\Omega), \quad a \geqslant 0 \text { a.e. }, \quad \mathrm{a} \not \equiv 0, \tag{5.3}
\end{equation*}
$$

and

$$
\begin{equation*}
f \in L_{\delta}^{1}(\Omega), \quad f \geqslant 0 \text { a.e. , } f \not \equiv 0 . \tag{5.4}
\end{equation*}
$$

A function $u \in L^{1}(\Omega), u \geqslant 0$ a.e. is a weak solution of (5.1) if $a u^{p} \in L_{\delta}^{1}(\Omega)$ and (5.1) is satisfied in the sense of Definition 1.1(c). Finally, for $h \in L_{\delta}^{1}(\Omega)$ we denote by $G(h)$ the unique function in $L^{1}(\Omega)$ satisfying

$$
\begin{cases}-\Delta(G(h))=h & \text { in } \Omega \\ G(h)=0 & \text { on } \partial \Omega\end{cases}
$$

again in the sense of Definition 1.1(c)—see e.g. Lemma 1 of [4] for such result
about the linear Laplace equation. We now give a necessary condition and a sufficient condition for the existence of a weak solution of (5.1).

Theorem 5.1. - Assume (5.3) and (5.4).
(a) If

$$
\begin{cases}-\Delta u=a(x) u^{p}+f(x) & \text { in } \Omega  \tag{5.5}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

has a weak solution $u \geqslant 0$, then $a G(f)^{p} \in L_{\delta}^{1}(\Omega)$ and

$$
\frac{G\left(a G(f)^{p}\right)}{G(f)} \leqslant \frac{1}{p-1} \quad \text { in } \Omega
$$

(b) If $a G(f)^{p} \in L_{\delta}^{1}(\Omega)$ and

$$
\frac{G\left(a G(f)^{p}\right)}{G(f)} \leqslant\left(\frac{p-1}{p}\right)^{p} \frac{1}{p-1} \quad \text { in } \Omega
$$

then (5.5) has a weak solution $u$ satisfying $G(f) \leqslant u \leqslant C G(f)$ in $\Omega$ for some constant $C$.

The second result of this section is the following.
Theorem 5.2. - Assume (5.3), (5.4) and

$$
\frac{G\left(a G(f)^{p}\right)}{G(f)} \in L^{\infty}(\Omega)
$$

For $\lambda$ a positive parameter, consider the problem

$$
\begin{cases}-\Delta u=a(x) u^{p}+\lambda f(x) & \text { in } \Omega  \tag{5.6}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

Then there exists $\lambda^{\star} \in(0, \infty)$ such that
(i) if $0<\lambda<\lambda^{\star}$, then (5.6) has a weak solution $u_{\lambda}$ satisfying

$$
\lambda \leqslant \frac{u_{\lambda}}{G(f)} \leqslant C(\lambda) \quad \text { in } \Omega
$$

for some constant $C(\lambda)$ depending on $\lambda$.
(ii) if $\lambda=\lambda^{\star}$, then $(5.6)_{\lambda}$ has a weak solution.
(iii) if $\lambda>\lambda^{\star}$, then (5.6) has no weak solution.

Moreover,

$$
\begin{equation*}
\left(\frac{p-1}{p}\right)^{p} \frac{1}{p-1} \leqslant\left(\lambda^{\star}\right)^{p-1}\left\|\frac{G\left(a G(f)^{p}\right)}{G(f)}\right\|_{L^{\infty}(\Omega)} \leqslant \frac{1}{p-1} \tag{5.7}
\end{equation*}
$$

The main ingredient in the proof of the above theorems is the following.

Lemma 5.3. - Suppose $u$ and $v$ are $C^{2}$ functions in $\Omega$, and that $v>0$. Let $\phi: \mathbb{R} \rightarrow \mathbb{R}$ be a $C^{2}$, concave function. Then

$$
-\Delta\left[v \phi\left(\frac{u}{v}\right)\right] \geqslant \phi^{\prime}\left(\frac{u}{v}\right)(-\Delta u)+\left[\phi\left(\frac{u}{v}\right)-\frac{u}{v} \phi^{\prime}\left(\frac{u}{v}\right)\right](-\Delta v) .
$$

If, in addition, $-\Delta v \geqslant 0$ in $\Omega$, then

$$
\begin{equation*}
-\Delta\left[v \phi\left(\frac{u}{v}\right)\right] \geqslant \phi^{\prime}\left(\frac{u}{v}\right)[-\Delta u+\Delta v]+\phi(1)(-\Delta v) \tag{5.8}
\end{equation*}
$$

Proof. - We simply compute and use that $\phi^{\prime \prime} \leqslant 0$ and $v>0$. We have (using the notation $\partial_{i} v=v_{i}$ )

$$
\begin{aligned}
& -\Delta\left[v \phi\left(\frac{u}{v}\right)\right]=-\sum_{i=1}^{N}\left[\phi^{\prime}\left(\frac{u}{v}\right) v\left(\frac{u}{v}\right)_{i}+\phi\left(\frac{u}{v}\right) v_{i}\right]_{i}= \\
& \quad-\phi^{\prime \prime}\left(\frac{u}{v}\right) v\left|\nabla\left(\frac{u}{v}\right)\right|^{2}-\sum_{i=1}^{N}\left\{\phi^{\prime}\left(\frac{u}{v}\right)\left[v\left(\frac{u}{v}\right)_{i}\right]_{i}+\left[\phi\left(\frac{u}{v}\right) v_{i}\right]_{i}\right\} \geqslant \\
& \quad-\sum_{i=1}^{N}\left\{\phi^{\prime}\left(\frac{u}{v}\right)\left[u_{i}-\frac{u}{v} v_{i}\right]_{i}+\phi\left(\frac{u}{v}\right) v_{i i}+\phi^{\prime}\left(\frac{u}{v}\right)\left(\frac{u}{v}\right)_{i} v_{i}\right\}= \\
& \phi^{\prime}\left(\frac{u}{v}\right)(-\Delta u)+\phi^{\prime}\left(\frac{u}{v}\right) \nabla\left(\frac{u}{v}\right) \nabla v-\frac{u}{v} \phi^{\prime}\left(\frac{u}{v}\right)(-\Delta v)+ \\
& \quad \phi\left(\frac{u}{v}\right)(-\Delta v)-\phi^{\prime}\left(\frac{u}{v}\right) \nabla\left(\frac{u}{v}\right) \nabla v= \\
& \phi^{\prime}\left(\frac{u}{v}\right)(-\Delta u)+\left[\phi\left(\frac{u}{v}\right)-\frac{u}{v} \phi^{\prime}\left(\frac{u}{v}\right)\right](-\Delta v),
\end{aligned}
$$

which is the first inequality of the lemma. From this, we deduce the second inequality, as follows. Since $\phi$ is concave, we have

$$
\phi(s)+(1-s) \phi^{\prime}(s) \geqslant \phi(1), \quad \forall s \in \mathbb{R} .
$$

Thus

$$
\phi\left(\frac{u}{v}\right)-\frac{u}{v} \phi^{\prime}\left(\frac{u}{v}\right) \geqslant-\phi^{\prime}\left(\frac{u}{v}\right)+\phi(1)
$$

multiplying this inequality by $-\Delta v$ (which is nonnegative by assumption), we easily deduce (5.8).

To use the previous lemma, we will need (5.8) in its weak version for $L^{1}$ functions.

LEMMA 5.4. - Let $\phi: \mathbb{R} \rightarrow \mathbb{R}$ be a $C^{1}$, concave function with $\phi^{\prime}$ bounded. Let $h$ and $k$ belong to $L_{\delta}^{1}(\Omega)$, with $k \geqslant 0, k \not \equiv 0$, and let $u$ and $v$ be the $L^{1}(\Omega)$ solutions of

$$
\begin{cases}-\Delta u=h & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

and

$$
\begin{cases}-\Delta v=k & \text { in } \Omega \\ v=0 & \text { on } \partial \Omega\end{cases}
$$

Then

$$
\begin{equation*}
-\Delta\left[v \phi\left(\frac{u}{v}\right)\right] \geqslant \phi^{\prime}\left(\frac{u}{v}\right)(h-k)+\phi(1) k, \tag{5.9}
\end{equation*}
$$

in the sense that $v \phi(u / v) \in L^{1}(\Omega), \phi^{\prime}(u / v)(h-k)+\phi(1) k \in L_{\delta}^{1}(\Omega)$ and

$$
-\int_{\Omega} v \phi\left(\frac{u}{v}\right) \Delta \zeta \geqslant \int_{\Omega}\left\{\phi^{\prime}\left(\frac{u}{v}\right)(h-k)+\phi(1) k\right\} \zeta
$$

for all $\zeta \in C^{2}(\bar{\Omega}), \zeta \geqslant 0$, with $\zeta=0$ on $\partial \Omega$.
Proof. - We first point out that (5.8) holds when $\phi$ is $C^{1}$ and concave-not necessarily $C^{2}$. This follows immediately from Lemma 5.3 convoluting $\phi$ with mollifiers.

We approximate $h$ and $k$ in $L_{\delta}^{1}(\Omega)$ by sequences $\left(h_{n}\right)$ and $\left(k_{n}\right)$, respectively, of $C_{0}^{\infty}(\Omega)$ functions and with $k_{n} \geqslant 0, k_{n} \not \equiv 0$. Let $u_{n}, v_{n}$ be the solutions of

$$
\begin{cases}-\Delta u_{n}=h_{n} & \text { in } \Omega, \\ u_{n}=0 & \text { on } \partial \Omega\end{cases}
$$

and

$$
\begin{cases}-\Delta v_{n}=k_{n} & \text { in } \Omega \\ v_{n}=0 & \text { on } \partial \Omega\end{cases}
$$

It follows that $u_{n} \rightarrow u$ and $v_{n} \rightarrow v$ in $L^{1}(\Omega)$ (for this, subtract the equations for $u_{n}$ and $u$, multiply by $G(1)$ and integrate). Moreover, $u_{n}, v_{n} \in C^{2}(\Omega)$ and $v_{n}>0$, $-\Delta v_{n} \geqslant 0$ in $\Omega$.

By (5.8) we have

$$
-\Delta\left[v_{n} \phi\left(\frac{u_{n}}{v_{n}}\right)\right] \geqslant \phi^{\prime}\left(\frac{u_{n}}{v_{n}}\right)\left(h_{n}-k_{n}\right)+\phi(1) k_{n} .
$$

Moreover, using that $\phi^{\prime}$ is bounded, we see that

$$
\begin{equation*}
\left|v_{n} \phi\left(\frac{u_{n}}{v_{n}}\right)\right|=\left|v_{n}\left(\phi\left(\frac{u_{n}}{v_{n}}\right)-\phi(0)\right)+\phi(0) v_{n}\right| \leqslant C\left(\left|u_{n}\right|+\left|v_{n}\right|\right) \tag{5.10}
\end{equation*}
$$

for some constant $C$. Hence, $v_{n} \phi\left(u_{n} / v_{n}\right)$ vanishes on $\partial \Omega$ and thus

$$
\begin{equation*}
-\int_{\Omega} v_{n} \phi\left(\frac{u_{n}}{v_{n}}\right) \Delta \zeta \geqslant \int_{\Omega}\left\{\phi^{\prime}\left(\frac{u_{n}}{v_{n}}\right)\left(h_{n}-k_{n}\right)+\phi(1) k_{n}\right\} \zeta \tag{5.11}
\end{equation*}
$$

for all $\zeta \in C^{2}(\bar{\Omega}), \zeta \geqslant 0$ in $\Omega$ and $\zeta=0$ on $\partial \Omega$.
Note that $v>0$ a.e. in $\Omega$, so that $v \phi(u / v)$ is well defined a.e. Moreover, $v_{n} \phi\left(u_{n} / v_{n}\right)$ converges a.e. to $v \phi(u / v)$ - up to a subsequence. Since $u_{n}$ and $v_{n}$ converge in $L^{1}(\Omega)$, they are dominated (also up to a subsequence) by an $L^{1}(\Omega)$ function. Thus, by (5.10), $\left|v_{n} \phi\left(u_{n} / v_{n}\right)\right|$ is also dominated by an $L^{1}$ function (for a subsequence). We conclude that

$$
v_{n} \phi\left(\frac{u_{n}}{v_{n}}\right) \rightarrow v \phi\left(\frac{u}{v}\right) \quad \text { in } L^{1}(\Omega) .
$$

Passing to the limit in (5.11), we finally obtain (5.9) and the lemma.
We write explicitly the concave functions $\phi$ that we use in this section. For Theorem 5.1 we will use

$$
\begin{equation*}
\phi(s)=\int_{1}^{s} \frac{d t}{t^{p}}=\frac{1}{p-1}\left(1-\frac{1}{s^{p-1}}\right) \quad \text { for } s \geqslant 1 \tag{5.12}
\end{equation*}
$$

It satisfies

$$
\phi^{\prime}(s) s^{p}=1 \quad \text { for } s \geqslant 1,
$$

and hence $\phi$ is concave and $\phi^{\prime}$ is bounded in [1, $\infty$ ). Moreover,

$$
\phi(1)=0 \quad \text { and } \quad 0 \leqslant \phi(s) \leqslant \frac{1}{p-1} \quad \text { for } \mathrm{s} \geqslant 1
$$

Since $\phi^{\prime}(1)=1$, we can extend $\phi$ by $\phi(s)=s-1$ in $(-\infty, 1]$ obtaining a function $\phi$ that satisfies the conditions of Lemma 5.4.

Note that the functions $\phi$ above are analogous versions of the ones that we employed in Sections 1 and 2, in the sense that they all satisfy $\phi^{\prime}(s) g(s)=1$ (in an appropriate interval) where $g$ is the nonlinearity.

In Theorem 5.2 we will be led to take $\phi$ (that we denote now by $\psi$ ) satisfying another differential equation, namely: $\psi^{\prime}(s) s^{p}=\psi(s)^{p}$. Precisely, we will take

$$
\begin{equation*}
\psi(s)=\frac{s}{\left(\varepsilon s^{p-1}+1\right)^{1 /(p-1)}} \quad \text { for } s \geqslant 0 \tag{5.13}
\end{equation*}
$$

It satisfies

$$
\psi^{\prime}(s)=\frac{1}{\left(\varepsilon s^{p-1}+1\right)^{p /(p-1)}} \quad \text { for } s \geqslant 0
$$

and hence

$$
\psi^{\prime}(s) s^{p}=\psi(s)^{p} \quad \text { for } s \geqslant 0 .
$$

Note that $\psi$ is concave and $\psi^{\prime}$ is bounded in [0, $\infty$ );

$$
\psi(1)=\left(\frac{1}{1+\varepsilon}\right)^{1 /(p-1)} \quad \text { and } \quad 0 \leqslant \psi(s) \leqslant\left(\frac{1}{\varepsilon}\right)^{1 /(p-1)} \quad \text { for } s \geqslant 0
$$

Extending $\psi$ by $\psi(s)=s$ in $(-\infty, 0]$, we obtain a function $\psi$ which satisfies the conditions (for $\phi$ ) of Lemma 5.4. Note that, for $p=2, \psi$ was already considered in Remark 1.8.

## Proof of Theorem 5.1.

Part (a). Let $u \geqslant 0$ be a weak solution of

$$
\begin{cases}-\Delta u=a(x) u^{p}+f(x) & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

We consider

$$
v=G(f)
$$

note that, since $a(x) u^{p} \geqslant 0$,

$$
\frac{u}{v}=\frac{u}{G(f)} \geqslant 1 \quad \text { in } \Omega
$$

by the weak maximum principle (which is an easy consequence of the weak formulation of Definition 1.1(c) used here). In particular $\alpha G(f)^{p} \leqslant a u^{p}$ and hence $a G(f)^{p} \in L_{\delta}^{1}(\Omega)$.

We take $\phi$, defined by (5.12), and apply Lemma 5.4 to obtain (we use $u / v \geqslant$ 1 and the properties of $\phi(s)$ for $s \geqslant 1$ )

$$
-\Delta\left[v \phi\left(\frac{u}{v}\right)\right] \geqslant \phi^{\prime}\left(\frac{u}{v}\right)\left(a u^{p}+f-f\right)+\phi(1) f=\phi^{\prime}\left(\frac{u}{v}\right) a u^{p}=a v^{p}=\alpha G(f)^{p}
$$

in the weak sense of the lemma. Hence (by the weak maximum principle)

$$
G\left(a G(f)^{p}\right) \leqslant v \phi\left(\frac{u}{v}\right) \leqslant \frac{1}{p-1} v=\frac{1}{p-1} G(f)
$$

Part ( $a$ ) is now proved.

Part (b). We assume that $a G(f)^{p} \in L_{\delta}^{1}(\Omega)$ and

$$
G\left(a G(f)^{p}\right) \leqslant\left(\frac{p-1}{p}\right)^{p} \frac{1}{p-1} G(f) .
$$

It follows that the $L^{1}(\Omega)$ function

$$
\tilde{u}:=\left(\frac{p}{p-1}\right)^{p} G\left(a G(f)^{p}\right)+G(f)
$$

satisfies

$$
\tilde{u} \leqslant \frac{p}{p-1} G(f)
$$

Therefore

$$
-\Delta \tilde{u}=\left(\frac{p}{p-1}\right)^{p} a G(f)^{p}+f \geqslant a \tilde{u}^{p}+f
$$

in the weak sense. That is, $\tilde{u}$ is a weak supersolution of (5.5). On the other hand, 0 is a subsolution of the problem. It is then easy to obtain a weak sol-
ution $u$ of (5.5) by monotone iteration (see e.g. Lemma 3 of [4]). Moreover

$$
G(f) \leqslant u \leqslant \tilde{u} \leqslant \frac{p}{p-1} G(f)
$$

This completes the proof of Theorem 5.1.
Proof of Theorem 5.2. - We assume that

$$
\frac{G\left(a G(f)^{p}\right)}{G(f)} \in L^{\infty}(\Omega)
$$

moreover, $a \not \equiv 0$ and $f \not \equiv 0$ and hence

$$
0<M_{\infty}:=\|\left.\frac{G\left(a G(f)^{p}\right)}{G(f)}\right|_{L^{\infty}(\Omega)}<\infty
$$

Theorem 5.1 (applied with $f$ replaced by $\lambda f$ ) implies that if (5.6) $)_{\lambda}$ has a weak solution then

$$
\lambda^{p-1} M_{\infty} \leqslant \frac{1}{p-1} .
$$

The theorem also gives that if

$$
\lambda^{p-1} M_{\infty} \leqslant\left(\frac{p-1}{p}\right)^{p} \frac{1}{p-1}
$$

then (5.6) ${ }_{\lambda}$ has a weak solution. Hence, defining

$$
\lambda^{\star}=\sup \left\{\lambda>0 ; \quad(5.6)_{\lambda} \text { has a weak solution }\right\}
$$

we have $0<\lambda^{\star}<\infty$, and also estimate (5.7) of Theorem 5.2. Note that part (iii) of the theorem is now obvious.

To prove part (i), we have to show that if $0<\lambda<\mu$ and $(5.6)_{\mu}$ has a weak solution $u$ then (5.6) $)_{\lambda}$ has a weak solution $u_{\lambda}$ satisfying

$$
\begin{equation*}
\lambda \leqslant \frac{u_{\lambda}}{G(f)} \leqslant C(\lambda) \quad \text { in } \Omega \tag{5.14}
\end{equation*}
$$

For this purpose, we consider

$$
v=G(\mu f)=\mu G(f)
$$

and the function $\psi$ defined by (5.13) with $\varepsilon>0$ chosen small enough such that

$$
\lambda \leqslant\left(\frac{1}{1+\varepsilon}\right)^{1 /(p-1)} \mu=\psi(1) \mu
$$

We apply Lemma 5.4 (with $\phi$ replaced by $\psi$ ) to obtain

$$
\begin{aligned}
& -\Delta\left[v \psi\left(\frac{u}{v}\right)\right] \geqslant \psi^{\prime}\left(\frac{u}{v}\right)\left(a u^{p}+\mu f-\mu f\right)+\psi(1) \mu f= \\
& \quad a v^{p}\left[\psi\left(\frac{u}{v}\right)\right]^{p}+\psi(1) \mu f \geqslant a\left[v \psi\left(\frac{u}{v}\right)\right]^{p}+\lambda f
\end{aligned}
$$

in the weak sense of the lemma. Hence $v \psi(u / v)$ is a weak supersolution of (5.6) $\lambda_{\lambda}$. Again by monotone iteration we deduce that (5.6) $)_{\lambda}$ has a weak solution $u_{\lambda}$ such that

$$
u_{\lambda} \leqslant v \psi\left(\frac{u}{v}\right) \leqslant C v=C \mu G(f) .
$$

This, together with the trivial bound $u_{\lambda} \geqslant G(\lambda f)$, gives (5.14) and proves (i).

It remains to show part (ii). For $\lambda<\lambda^{\star}$ we can take $u_{\lambda}$ to be the minimal solution of (5.6) , i.e., the solution obtained by monotone iteration starting from the function 0 . In this manner, $u_{\lambda} \leqslant u_{\mu}$ if $0<\lambda<\mu<\lambda^{\star}$. Hence, in order to obtain a weak solution of $(5.6)_{\lambda^{\star}}$, it suffices to show

$$
\begin{equation*}
\int_{\Omega}\left[a(x) u_{\lambda}^{p}+\lambda f(x)\right] \delta \leqslant C \tag{5.15}
\end{equation*}
$$

for some constant $C$ independent of $\lambda$.
To prove (5.15) we proceed as follows. Since $a \geqslant 0$ and $\alpha \not \equiv 0$, there exists a constant $M \in(0, \infty)$ such that

$$
a_{M}:=\chi_{\{a \leqslant M\}} a
$$

satisfies $0 \leqslant a_{M} \leqslant M$ and $a_{M} \not \equiv 0$ (here $\chi_{\{a \leqslant M\}}$ denotes the characteristic function of $\{a \leqslant M\}$ ).

It is well-known that the problem

$$
\begin{cases}-\Delta w=a_{M}^{1 / p} w^{1 / p} & \text { in } \Omega  \tag{5.16}\\ w>0 & \text { in } \Omega \\ w=0 & \text { on } \partial \Omega\end{cases}
$$

has a unique solution $w \in W^{2, r}(\Omega)$ (for any $1<r<\infty$ ) with $w \not \equiv 0$. This sol-
ution can be obtained for example by minimizing in $H_{0}^{1}(\Omega)$ the energy associated to (5.16):

$$
E(v)=\frac{1}{2} \int_{\Omega}|\nabla v|^{2}-\frac{p}{p+1} \int_{\Omega} a_{M}^{1 / p}\left(v^{+}\right)^{(p+1) / p}
$$

which is a coercive functional, bounded from below, in $H_{0}^{1}(\Omega)$ (note that $a_{M}^{1 / p} \in$ $L^{\infty}(\Omega)$ and $\left.1<(p+1) / p<2\right)$. We have that $E\left(t \varphi_{1}\right)<0$ for $t$ small, where $\varphi_{1}$ denotes the first eigenfunction of $-\Delta$. In particular, the minimizer $w$ of $E$ satisfies $E(w)<0$, and thus $w \not \equiv 0$. The strong maximum principle then gives

$$
\begin{equation*}
w \geqslant c \delta \tag{5.17}
\end{equation*}
$$

for some positive constant $c$.
Since $\Delta w \in L^{\infty}(\Omega)$ we can multiply (5.6) by $w$ and integrate by parts. We also use Young's inequality-in the spirit of the ideas of Baras and Pierre [3]. We obtain

$$
\begin{aligned}
\int_{\Omega} a u_{\lambda}^{p} w+\int_{\Omega} \lambda f w= & \int_{\Omega} u_{\lambda}(-\Delta w)= \\
& \int_{\Omega} a_{M}^{1 / p} u_{\lambda} w^{1 / p} \leqslant \int_{\Omega} a^{1 / p} u_{\lambda} w^{1 / p} \leqslant \frac{1}{p} \int_{\Omega} a u_{\lambda}^{p} w+\frac{p-1}{p} \int_{\Omega} 1,
\end{aligned}
$$

and therefore

$$
\int_{\Omega} a u_{\lambda}^{p} w+\int_{\Omega} \lambda f w \leqslant C
$$

for some constant $C$ independent of $\lambda$. Using (5.17), we conclude (5.15), and hence the proof of part (ii).

Remark 5.5. - An analogous version of Theorem 5.2 also holds for the problem

$$
\begin{cases}-\Delta u=\lambda\left(a(x) u^{p}+f(x)\right) & \text { in } \Omega, \\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

This follows from Theorem 5.2 rescaling the solution $u$, i.e., considering the problem for $\alpha u$, for appropriate $\alpha$.

Remark 5.6. - In Theorem 5.2 we have shown that, for $\lambda<\lambda^{\star},(5.6)_{\lambda}$ has a weak solution $u_{\lambda}$ satisfying

$$
\begin{equation*}
\frac{u_{\lambda}}{G(f)} \in L^{\infty}(\Omega) \tag{5.18}
\end{equation*}
$$

We point out that this property may not be true for the minimal solution obtained for $\lambda=\lambda^{\star}$. To give an example of this, consider first the problem

$$
\begin{cases}-\Delta v=\lambda(v+1)^{p} & \text { in } B_{1}=\{|x|<1\} \subset \mathbb{R}^{N}  \tag{5.19}\\ v=0 & \text { on } \partial B_{1}\end{cases}
$$

For $N$ and $p$ sufficiently large, (5.19) $)_{\lambda}$ has

$$
\bar{v}(x)=|x|^{-2 /(p-1)}-1
$$

as weak solution for a certain parameter $\bar{\lambda}>0$; moreover $\bar{v}$ is the pointwise, increasing limit (as $\lambda \nearrow \bar{\lambda}$ ) of the classical minimal solutions $v_{\lambda}$ of (5.19) for $\lambda<\bar{\lambda}$ (see e.g. [6]). Let us define

$$
f=(\bar{v}+1)^{p}-\bar{v}^{p}
$$

so that the problem

$$
\begin{cases}-\Delta u=\bar{\lambda} u^{p}+\lambda f(x) & \text { in } B_{1}  \tag{5.20}\\ u=0 & \text { on } \partial B_{1}\end{cases}
$$

has $u=\bar{v}$ as weak solution for $\lambda=\bar{\lambda}$. Note that $(5.20)_{\lambda}$ is of the form (5.6) $)_{\lambda}$ considered in Theorem 5.2. We claim that $\bar{v}$ is the minimal weak solution of $(5.20)_{\bar{\lambda}}$. This is shown as follows: for $\lambda<\bar{\lambda}$, let $u_{\lambda}$ be the minimal weak solution of (5.20) $)_{\lambda}$. Since $u_{\lambda} \leqslant \bar{v}$, we have $\left(u_{\lambda}+1\right)^{p}-u_{\lambda}^{p} \leqslant f$, and hence $u_{\lambda}$ is a weak supersolution of (5.19) . Thus $u_{\lambda} \geqslant v_{\lambda}$, and since $v_{\lambda} \nearrow \bar{v}$ as $\lambda \nearrow \bar{\lambda}$, we conclude that $\bar{v}$ is the minimal weak solution of $(5.20)_{\bar{\lambda}}$.

Finally, we have that

$$
\begin{equation*}
\frac{\bar{v}}{G(f)} \notin L^{\infty}\left(B_{1}\right) \tag{5.21}
\end{equation*}
$$

since $\bar{v} \sim|x|^{-2 /(p-1)}$ near 0 and $G(f) \leqslant C \log (1 /|x|)$ (note that $f \leqslant p(\bar{v}+$ $1)^{p-1}=p|x|^{-2}$ in $B_{1}$ ). This proves that (5.18) does not hold for $\lambda=\bar{\lambda}$ and, in particular, that $\bar{\lambda}$ is the extremal parameter $\lambda^{\star}$ for problem $(5.20)_{\lambda}$.

## 6. - Evolution equations.

In Section 4 we proved that $u \equiv 0$ is the only nonnegative very weak supersolution of

$$
-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{O}^{\prime}(\Omega \backslash\{0\}) .
$$

Here we prove an analogous result for the equation

$$
u_{t}-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{O}^{\prime}((\Omega \backslash\{0\}) \times(0, T))
$$

More precisely, we have the following parabolic analogue of Theorem 4.1.

Theorem 6.1. - Let $N \geqslant 2, T>0$ and $u \in L_{\text {loc }}^{2}((\Omega \backslash\{0\}) \times(0, T))$ satisfy $u \geqslant 0$ a.e. in $\Omega \times(0, T)$ and

$$
\begin{equation*}
|x|^{2}\left(u_{t}-\Delta u\right) \geqslant u^{2} \quad \text { in } \mathscr{O}^{\prime}((\Omega \backslash\{0\}) \times(0, T)), \tag{6.1}
\end{equation*}
$$

in the sense that

$$
-\iint u\left\{\left(|x|^{2} \varphi\right)_{t}+\Delta\left(|x|^{2} \varphi\right)\right\} \geqslant \iint u^{2} \varphi
$$

for any $\varphi \geqslant 0, \varphi \in C_{0}^{\infty}((\Omega \backslash\{0\}) \times(0, T))$.
Then $u \equiv 0$.

As an immediate consequence of the theorem we obtain an extension of a result of Peral and Vázquez (Theorem 7.1 of [12]). Here $\Omega=B_{1}$, the unit ball of $\mathbb{R}^{N}$, and $N \geqslant 3$. We consider the function

$$
\bar{u}(x)=\log \frac{1}{|x|^{2}},
$$

which is a weak solution of

$$
\begin{cases}-\Delta \bar{u}=2(N-2) e^{\bar{u}} & \text { in } B_{1}, \\ \bar{u}=0 & \text { on } \partial B_{1} .\end{cases}
$$

Corollary 6.2. - Let $N \geqslant 3, T>0$ and $u \in L_{\text {loc }}^{1}\left(\left(B_{1} \backslash\{0\}\right) \times(0, T)\right)$ be a function such that $e^{u} \in L_{\text {loc }}^{1}\left(\left(B_{1} \backslash\{0\}\right) \times(0, T)\right)$,

$$
u(x, t) \geqslant \bar{u}(x) \quad \text { a.e. in } B_{1} \times(0, T),
$$

and

$$
\begin{equation*}
u_{t}-\Delta u=2(N-2) e^{u} \quad \text { in } \propto^{\prime}\left(\left(B_{1} \backslash\{0\}\right) \times(0, T)\right), \tag{6.2}
\end{equation*}
$$

in the sense that

$$
-\iint u\left(\varphi_{t}+\Delta \varphi\right)=2(N-2) \iint e^{u} \varphi
$$

for any $\varphi \in C_{0}^{\infty}\left(\left(B_{1} \backslash\{0\}\right) \times(0, T)\right)$.
Then $u(x, t) \equiv \bar{u}(x)$.
Note again that we only assume equation (6.2) to be satisfied in the distributional sense and away from $\{x=0\} \times[0, T]$. In particular, given any $u_{0}(x) \geqslant \bar{u}(x)$ a.e. in $B_{1}, u_{0} \neq \bar{u}$, and any $T>0$, there is no weak solution $u$, with

$$
u(x, t) \geqslant \bar{u}(x) \quad \text { in } B_{1} \times(0, T),
$$

of the problem

$$
\begin{cases}u_{t}-\Delta u=2(N-2) e^{u} & \text { in } B_{1} \times(0, T),  \tag{6.3}\\ u=0 & \text { on } \partial B_{1} \times(0, T), \\ u(x, 0)=u_{0} & \text { on } B_{1},\end{cases}
$$

(as stated in [12]).
A second consequence of Theorem 6.1 is the following nonexistence and complete blow-up result. For any $u_{0} \geqslant 0, u_{0} \not \equiv 0$ (say $u_{0} \in C_{0}^{\infty}(\Omega)$ ) and for any $T>0$, the problem

$$
\begin{cases}u_{t}-\Delta u=\frac{u^{2}}{|x|^{2}} & \text { in } \Omega \times(0, T),  \tag{6.4}\\ u=0 & \text { on } \partial \Omega \times(0, T), \\ u(x, 0)=u_{0} & \text { on } \Omega\end{cases}
$$

has no weak solution (by Theorem 6.1). Using similar ideas as in the elliptic case (see Section 3), we prove that approximate solutions of (6.4) blow up completely. More precisely, let $g_{n}$ be a sequence of nonnegative, nondecreasing and globally Lipschitz functions in $[0, \infty)$ such that $g_{n}(u)$ increases pointwise to $u^{2}$. Let $a_{n}$ be a sequence of nonnegative bounded functions in $\Omega$, increasing pointwise to $1 /|x|^{2}$.

Theorem 6.3. - Under the above assumptions, let $u_{n}$ be the solution of
(6.5) $)_{n} \begin{cases}\frac{\partial u_{n}}{\partial t}-\Delta u_{n}=a_{n}(x) g_{n}\left(u_{n}\right) & \text { in } \Omega \times(0,+\infty), \\ u_{n}=0 & \text { on } \partial \Omega \times(0,+\infty), \\ u_{n}(x, 0)=u_{0} & \text { on } \Omega .\end{cases}$

Then, for any $0<\varepsilon<T$,

$$
\begin{equation*}
\frac{u_{n}(x, t)}{\delta(x)} \rightarrow+\infty \quad \text { uniformly in }(x, t) \in \Omega \times[\varepsilon, T] \tag{6.6}
\end{equation*}
$$

as $n \rightarrow \infty$.
To prove Theorem 6.1 we adapt the method given in Section 4 for the elliptic case; it consists of using appropriate powers of testing functions.

Proof of Theorem 6.1. - We proceed as in the proof of Theorem 4.1; we use the same notation as there. We also fix a cut-off function in time:

$$
\psi \in C_{0}^{\infty}((0, T)), \quad 0 \leqslant \psi \leqslant 1
$$

with $\psi \equiv 1$ in a given compact sub-interval of $(0, T)$.
Step 1. We prove that $u /|x| \in L_{\mathrm{loc}}^{2}(\Omega \times(0, T))$. For this purpose, we multiply (6.1) by $\zeta_{n}^{4}(x) \psi^{2}(t) /|x|^{2}$; it yields

$$
\iint \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4} \psi^{2} \leqslant-\iint u\left(\Delta \zeta_{n}^{4}\right) \psi^{2}-\iint u \zeta_{n}^{4}\left(\psi^{2}\right)_{t}
$$

Let us denote by $C$ different constants independent of $n$, but that may depend on $\Omega, T$ and the cut-off $\psi$. We have

$$
-\iint u\left(\Delta \zeta_{n}^{4}\right) \psi^{2} \leqslant C n\left(\int_{\{1 / n<|x|<2 / n\} \times(0, T)} \int_{|x|} \frac{u}{n} \zeta_{n}^{2} \psi\right)+C
$$

and

$$
-\iint u \zeta_{n}^{4}\left(\psi^{2}\right)_{t} \leqslant C \iint \frac{u}{|x|} \zeta_{n}^{2} \psi
$$

Hence we can conclude, as in Section 4, that

$$
\iint \frac{u^{2}}{|x|^{2}} \zeta_{n}^{4} \psi^{2} \leqslant C
$$

Thus $u /|x| \in L_{\text {loc }}^{2}(\Omega \times(0, T))$.
Step 2. We show that

$$
\begin{equation*}
u_{t}-\Delta u \geqslant \frac{u^{2}}{|x|^{2}} \quad \text { in } \mathscr{\sigma}^{\prime}(\Omega \times(0, T)) \tag{6.7}
\end{equation*}
$$

This is done exactly as in the proof of Theorem 4.1, where now $\varphi=\varphi(x, t)$ belongs to $C_{0}^{\infty}(\Omega \times(0, T))$.

Step 3 . We finally prove $u \equiv 0$. We suppose $u \not \equiv 0$. Since $u \geqslant 0, u_{t}-\Delta u \geqslant 0$ in $\Omega^{\prime}(\Omega \times(0, T))$ and $\Omega$ is connected, we have that

$$
u \geqslant \varepsilon \quad \text { a.e. in } B_{\eta} \times(\tau, T)
$$

for some $0<\tau<T, \varepsilon>0$ and $B_{\eta}=B_{\eta}(0)$ with closure in $\Omega$. When $N=2$ this is a contradiction with $u /|x| \in L_{\text {loc }}^{2}(\Omega \times(0, T))$.

When $N \geqslant 3$, (6.7) gives

$$
u_{t}-\Delta u \geqslant \frac{\varepsilon^{2}}{|x|^{2}}=\left(\partial_{t}-\Delta\right)\left(\frac{\varepsilon^{2}}{N-2} \log \frac{1}{|x|}\right) \quad \text { in } \sigma^{\prime}\left(B_{\eta} \times(\tau, T)\right)
$$

We deduce that

$$
\begin{equation*}
u \geqslant \frac{\varepsilon^{2}}{N-2} \log \frac{1}{|x|}-C \quad \text { in } B_{\eta / 2} \times\left(\tau^{\prime}, T\right) \tag{6.8}
\end{equation*}
$$

for some $C>0$ and $\tau<\tau^{\prime}<T$.
Following the proof of Theorem 4.1, we now multiply (6.7) by $\chi_{n}^{4}(x) \psi^{2}(t)$, with $\psi$ a cut-off as in the beginning of this proof, and with $\psi \equiv 1$ in $\left(\tau^{\prime}, T^{\prime}\right)$ for some $T^{\prime}$ with $\tau^{\prime}<T^{\prime}<T$. We have

$$
\iint \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \psi^{2} \leqslant \iint u\left|\Delta \chi_{n}^{4}\right| \psi^{2}+\iint u \chi_{n}^{4}\left|\left(\psi^{2}\right)_{t}\right|
$$

(where we are integrating on $\{|x|<2 / n\} \times(0, T)$, since it contains the support of $\chi_{n}^{4} \psi^{2}$ ). Hence

$$
\iint \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \psi^{2} \leqslant C n \iint \frac{u}{|x|} \chi_{n}^{2} \psi
$$

where $C$ is independent of $n$. From the Cauchy-Schwarz inequality, we deduce

$$
\iint \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \psi^{2} \leqslant \frac{C}{n^{N-2}}
$$

But using (6.8), we have

$$
\iint \frac{u^{2}}{|x|^{2}} \chi_{n}^{4} \psi^{2} \geqslant c \frac{|\log n|^{2}}{n^{N-2}}
$$

which contradicts the previous statement. This proves the theorem.
Remark 6.4. - Step 3 of the previous proof, (i.e., to show $u \equiv 0$ from $u \geqslant 0$ and (6.7)) could have been done using a parabolic analogue of the method of Section 1. That is, one considers

$$
v=\frac{1}{\varepsilon}-\frac{1}{u}
$$

in a subcylinder where $u \geqslant \varepsilon$. Then (with the aid of the parabolic Kato's inequality) $v$ satisfies

$$
v_{t}-\Delta v \geqslant \frac{1}{|x|^{2}}
$$

which leads to contradiction since $v$ is bounded $(v \leqslant 1 / \varepsilon)$.
Corollary 6.2 follows immediately from Theorem 6.1:
Proof of Corollary 6.2. - Let $u$ be as in the corollary. Consider $v(x, t)=$ $u(x, t)-\bar{u}(x)$. It satisfies, from our assumptions,

$$
v \geqslant 0 \quad \text { a.e. in } B_{1} \times(0, T)
$$

Moreover, in the distributional sense $\mathscr{\sigma}^{\prime}\left(\left(B_{1} \backslash\{0\}\right) \times(0, T)\right)$, $v_{t}-\Delta v=u_{t}-\Delta u+\Delta \bar{u}=2(N-2)\left(e^{u}-e^{\bar{u}}\right)=2(N-2) e^{\bar{u}}\left(e^{v}-1\right)=$

$$
\frac{2(N-2)}{|x|^{2}}\left(e^{v}-1\right) \geqslant \frac{(N-2)}{|x|^{2}} v^{2}
$$

since $v \geqslant 0$. Hence $(N-2) v \geqslant 0$ satisfies (6.1). By Theorem 6.1 we deduce $v \equiv$ 0 , that is $u \equiv \bar{u}$.

We finally give the proof of the complete blow-up result.

Proof of Theorem 6.3. - We proceed in three steps. Recall that $0 \leqslant u_{n} \leqslant$ $u_{n+1}$, by the maximum principle.

Step 1. We prove that, for any $\tau>0$,

$$
\int_{0}^{\tau} \int_{\Omega} a_{n} g_{n}\left(u_{n}\right) \delta \rightarrow+\infty
$$

Suppose not, that $\int_{0}^{\tau} \int_{\Omega} a_{n} g_{n}\left(u_{n}\right) \delta \leqslant C$ for some $\tau>0$. We multiply (6.5) $n_{n}$ by
the solution of

$$
\begin{cases}-\Delta z=1 & \text { in } \Omega \\ z=0 & \text { on } \partial \Omega\end{cases}
$$

We obtain

$$
\int_{0}^{\tau} \int_{\Omega} u_{n}+\int_{\Omega} u_{n}(x, \tau) z-\int_{\Omega} u_{0} z=\int_{0}^{\tau} \int_{\Omega} a_{n} g_{n}\left(u_{n}\right) z \leqslant C
$$

and, in particular

$$
\int_{0}^{\tau} \int_{\Omega} u_{n} \leqslant C
$$

Hence $u_{n}$ and $a_{n} g_{n}\left(u_{n}\right) \delta$ are bounded in $L^{1}(\Omega \times(0, \tau))$. By monotone convergence, we obtain that $u_{n} \rightarrow u$ in $L^{1}(\Omega \times(0, \tau))$ with $u$ satisfying the assumptions of Theorem 6.1. By this theorem, $u \equiv 0$. Thus $u_{1} \equiv 0$ and $u_{0}(x)=$ $u_{1}(x, 0) \equiv 0$, a contradiction.

Step 2. We show that

$$
\int_{\Omega} u_{n}(x, t) \delta(x) d x \rightarrow+\infty \text { uniformly in } t \in[\varepsilon / 2, T] .
$$

Indeed, let us multiply (6.5) ${ }_{n}$ by $e^{\lambda_{1} t} \varphi_{1}$, where $\varphi_{1}$ is the first eigenfunction of $-\Delta$ in $\Omega$ with zero Dirichlet condition and $\lambda_{1}$ its corresponding eigenvalue. We then integrate in space and time, to obtain

$$
e^{\lambda_{1} \tau} \int_{\Omega} u_{n}(\cdot, \tau) \varphi_{1}-\int_{\Omega} u_{0} \varphi_{1}=\int_{0}^{\tau} \int_{\Omega} a_{n} g_{n}\left(u_{n}\right) \varphi_{1} e^{\lambda_{1} t}
$$

Hence, if $\tau \in[\varepsilon / 2, T]$,

$$
\int_{\Omega} u_{n}(x, \tau) \delta(x) d x \geqslant c e^{-\lambda_{1} T} \int_{0}^{\varepsilon / 2} \int_{\Omega} a_{n} g_{n}\left(u_{n}\right) \delta \rightarrow+\infty
$$

by Step 1.
Step 3. We finally prove (6.6). For this purpose, we use a parabolic analogue of Lemma 3.2 due to Martel (see Lemma 2 of [10]). It asserts that if $\varphi \geqslant 0$ in $\Omega$ (say, $\varphi \in L^{\infty}(\Omega)$ ) then

$$
\frac{T(\tau) \cdot \varphi(x)}{\delta(x)} \geqslant c(\tau) \int_{\Omega} \varphi \delta, \quad \forall x \in \Omega
$$

for any $\tau>0$, where $c(\tau)>0$ is a constant depending on $\tau$, and where $T(\tau)$ is the heat semigroup at time $\tau$. We apply this estimate with $\tau=\varepsilon / 2$, and obtain

$$
\frac{u_{n}(x, t)}{\delta(x)} \geqslant c(\varepsilon / 2) \int_{\Omega} u_{n}(x, t-\varepsilon / 2) \delta(x) d x \rightarrow+\infty
$$

uniformly in $t \in[\varepsilon, T]$ by Step 2 , since $t-\varepsilon / 2 \geqslant \varepsilon / 2$.
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